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# THE ST. LAWRENCE RIVER

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#### ABSTRACT

Use has been made of the fundamental concepts of hydrodynamics to study the steady and tidal flows in the St. Lawrence River.

The Equations of Hydrodynamics take on a very simple form for steady flow as it is the case between Montréal and Trois Rivières. The definition and the meaning of the quantities they relate such as the mean depth D, the inclination of the river bed I, of the surface  $S_X$ , have been carefully reviewed. The mean inclination of the surface of the river  $S_X$ , which is always positive, is found to be remarkably constant from year to year in spite of the variations occurring in the discharge. Using this fact it is possible to deduce a value of the Chézy coefficient between Montréal and Sorel.

A plot of the mean level for all the stations located between King Edward Pier in Montréal and Pointe au Père shows the mean downwards slope of the surface of the river when undisturbed by the tide. The monthly fluctuations of this mean level decrease downstream since the capacity of the river bed increases in this direction. The mean flow is found to be subcritical and therefore cannot impede the propagation of the tide upstream.

The propagation of the tide may be understood using the time dependent form of the Equations of Hydrodynamics. An elementary solution of these equations shows that under the influence of linearized friction, the tide is damped exponentially and is retarded in its progress upstream. The actual occurrence of quadratic friction in the equations implies the creation of fast and slow shallow water constituents along the river; the speed of progress of the tide upstream is ascertained from observations. The use of the current observations in conjunction with the time dependent equations of motion allows the deduction of the value of the Chézy coefficient which prevails in the part of the river which extends between Trois-Rivières and Québec.

The spectral analysis of the water levels recorded by the water level gauges reveals that the low frequency band contains oscillation of 15, 9, 7 and 5 days which may be traced to the interaction of the tidal constituents. The change in phase and amplitude of the major tidal constituents deduced from the same analysis has been plotted and supplies the material which has to be reproduced in a mathematical model of the river. The current observations carried out by G.C. Dohler and the Ship Channel Division have been analyzed as well, whenever possible. The suggestion by W.D. Forrester that the average one dimensional currents in a river could be indirectly deduced from the readings on the tide gauges and the equation of continuity rather than by direct observations has been confirmed by our analyses.

A one dimensional numerical model of the St. Lawrence River has been elaborated. In it, the contribution of the tributaries, especially the Saguenay and the bifurcation of the river bed at Isle d'Orléans have been neglected. Still the model reproduces adequately the main features of the major tidal constituents; it reproduces qualitatively but not quantitatively, the shallow water constituents. It indicates that the fluctuations in the discharge are felt only upstream of the Ouebec constriction. The amplitude of the tide is little affected but its progress may be retarded by up to 3 hours. The slow shallow water constituents are increased when the discharge is increased but in the range of the value of the friction considered (C=58m½/sec to C=52m½/sec), the amplitude of the fast and slow shallow water constituents is decreased when friction is increased. Finally the model indicates that blockage anywhere in the river automatically involves an enhancement of the tide downstream.

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#### THE FUNDAMENTAL PARAMETERS OF A RIVER

The regime of a river may be characterized by the following quantities:

1) Q: the discharge. The amount of water which flows through a vertical cross section of the channel per unit time. This quantity is not easy to measure directly but it may be estimated by integrating the values of the velocity of the current measured at various points of the cross section. A change in discharge will cause a change in the depth of water, larger discharges normally implying greater depths at a given point (as seen in paragraph 4).

Units: meters3/second.

2) D: the depth of water. D varies at every point of the river depending on the bottom configuration and it normally varies with the discharge. If a portion of the river channel is schematized in such a way that a vertical cross section is a rectangle of width B and depth D, the mean depth D acquires a precise meaning.

Units: meters.

Units: dimensionless.

3) I: the inclination of the bed of the river. The water can flow in a river only if the bed of the river slopes downward and if it is remote from a large body of water with which it would be in hydrostatic equilibrium. The inclination of a river bed is not at all easy to visualize from ordinary chart soundings and it is necessary to schematize the river into channels of rectangular cross section in order to get some idea of it.

- 4) B: the width of the river. B happens to be an easy quantity to visualize but if marshes, reeds and islands are scattered near the shores, the river may "feel" a width which is quite different from the width which the eye can see.

  Units: meters.
- 5) The friction. Friction is an invisible quantity but it is the major factor controlling the flow of a river. Without it a river would always be dry as the water would run very rapidly towards the ocean. With friction present, there develops an equilibrium between the urge of the water to tumble downstream and the force of friction holding it back. Under steady conditions, a constant velocity of outflow develops and the river surface will be inclined at an angle with respect to the horizontal.

Units of the Chézy friction coefficient: meters 2/second.

#### STEADY STATE

When nothing changes in time, although the conditions may change from point to point, we wish to study the relations between the various quantities which we have just described and which are descriptive of a river. To measure them we have water level gauges and current meters. Gauges are not difficult to install and they can operate unattended; they give us the elevation of the water surface above a reference level. This quantity seems to satisfy the hydrographers but it does not suffice for our purpose as we are really in need of the quantity D, the actual depth of water. Current meters are more difficult to handle than gauges and cannot be left at one place

indefinitely. Used judiciously they can give us an idea of the value of the discharge Q at a section of the river, a quantity which is fundamental in our investigations.

Gauges strung along the river, if they are all calibrated to the same reference level, will show the inclination of the surface of the river as it flows downstream. In the portion of the river where the tide makes itself felt they will show the distortion of the water surface as the tidal wave moves in and out. In conjunction with hydrographic charts which give mean soundings undisturbed by the tide or by the irregularities in the discharge, we may compute mean values of the depth D over the river.

In the upstream portion of a river, steady state conditions do not quite prevail. The discharge varies from day to day, ice jams form in the spring which distort the profile drastically for short intervals of time. As a consequence the depth D and the inclination D vary from day to day. There arises the need to average out the fluctuations. The mean values of Q, D, I and B may be considered as characteristic of the particular river under study.

#### 1. The Equations of Hydrodynamics for steady state

For time independent motion, the quantities which we have just discussed are related through the equations (Stoker: Water Waves, p. 456; Dronkers: Tidal Computations, p. 158):

$$I - D_{x} - Q_{0}^{2}/C^{2}D^{3}B^{2} = 0$$
 (2)

where

Q =the discharge in  $m^3/$ sec

q = the acceleration due to gravity, 9.80 m/sec<sup>2</sup>

I = the inclination of the schematized bed of the river

D = the mean depth of the schematized portion of the channel,
in meters

 $D_{\rm X}$  = the rate of change looking downstream of the mean depth of water

C = the Chézy coefficient of friction in m2/sec

B = the effective width of the river in meters

x = the distance along the river measured positive downstream,
in meters

(2) contains an extra convective term  $(Q_0^2/gBD)$  (1/BD)<sub>x</sub> which is usually small compared to the other quantities. x used as an index indicates differentiation with respect to x.

We assume the absence of tributaries in (1) and (2).

Equation (1) simply states that the discharge throughout any
vertical section of the river is the same everywhere if steady
state prevails. This implies in turn that D cannot change
locally although it does change along the river as (2) indicates.

Essentially, all the information which (2) supplies is that:

$$D_{\rm X} = I - Q_0^2/C^2B^2D^3$$
 (3)

which tells us that the local rate of change of the depth of water depends on  $Q_0$ , on the friction coefficient C and on the width B. Strictly speaking, B and D vary for each value of  $Q_0$  and in the tidal portions of the river, vary at each instant of time. However, these quantities oscillate around mean values which are shown in Figs. 24 and 25 and which may be considered as representative of the St. Lawrence River. We had to plot the mean width B on a vertical logarithmic scale because it varies so immensely: from 800 meters near Deschambault to 52,000 meters at Matane. The depth D decreases more or less steadily from 250 meters at Pointe des Monts to 3 meters in Lake St. Peter.

The quantity which we know the least about in (3) is really not  $D_X$  which we may calculate indirectly from the readings on tide gauges and hydrographic charts but C the Chézy coefficient of friction. We may therefore actually use (3) to obtain estimates of the Chézy coefficient for the upstream portion of the river where steady state conditions prevail on the average. However, before we succeed in this task, we have to look closer on how we may go about measuring D and  $D_X$  from schematized sections of the river, gauge readings and hydrographic soundings and how we can succeed in obtaining some estimates of the inclination of the river bed I.

## 2. The profile of the river

Gauges strung along a river actually measure  $h_{\rm X}$ , the inclination of the water surface with reference to a horizontal level, the International Great Lakes Datum (IGLD) in the case of

the St. Lawrence River, rather than the rate of change of the mean depth with distance  $\text{D}_{\text{X}}$ .

Fig. 1 may help to visualize the various quantities involved. This figure refers to a hypothetical small portion of a river made up of many elementary sections of length  $\Delta x$ . The depth has been averaged across the river and over the elements of length  $\Delta x$ . This is what we call the depth D. The dark solid line indicates the actual surface of the river, the hatched line, the bottom of the river, assumedly sloping at a constant angle I downward. D is the depth of water at a given point. The gauge itself measures h, the depth of water above IGLD or its complement S. The dotted line indicates a constant depth  $D_0$  which is a possible regime of flow but which does not actually occur in the St. Lawrence River.

Using the diagram we see that we can write

$$D = z + h = z + h_O + \Delta$$

 $h_{\text{O}}$  would be the height read on the gauge if the depth were everywhere constant,  $\Delta$  is the deviation from this constant depth. Differentiation with respect to x gives

$$D_X = z_X + h_{OX} + \Delta_X = I - I + \Delta_X = \Delta_X$$

 $\Delta_{\mathrm{X}}$  simply measures the deviation of D from a surface of constant depth. Therefore,

$$D_{X} = \Delta_{X} \tag{4}$$

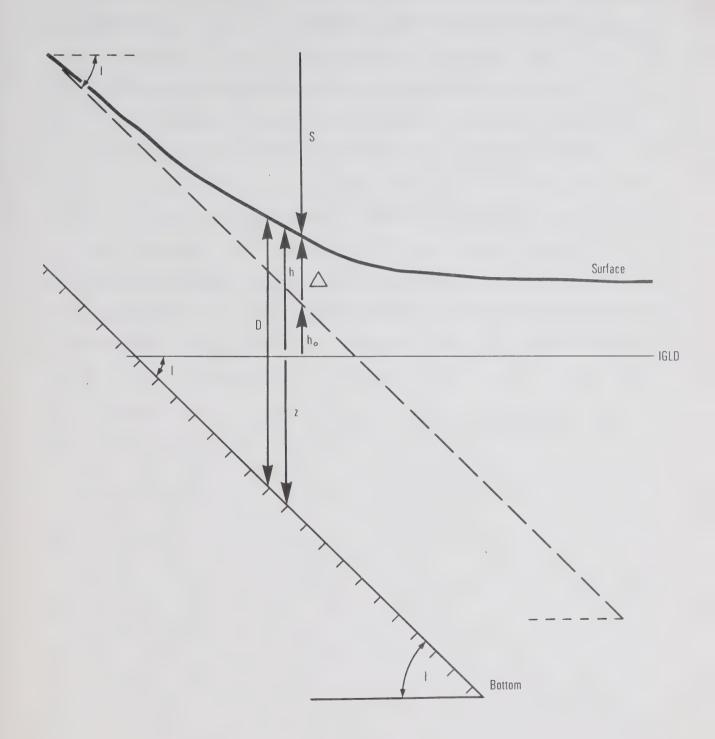


FIG.1 RIVER PROFILE

From a physical point of view we are more interested in h as read by the tide gauge or its complement S. Fig. 1 indicates

$$h_{X} = h_{O_{X}} + \Delta_{X} = -I + \Delta_{X} = -I + D_{X} = -I + I - Q_{O}^{2}/C^{2}B_{2}D_{3}$$

or

$$h_x = -S_x = -Q_0^2/C^2B^2D^3$$
 (5)

 $Q_0^2/C^2B^2D^3$  does not change sign and is always a positive quantity; therefore,  $h_X$  is always negative downstream, i.e. the water surface must always slope downward in steady state conditions or equivalently,  $S_X$  is always positive.

## 3. The significance of the quantities D and I

Although we have used D and I in the previous paragraph, we have not yet given them a precise meaning: this is now our task before we proceed to use formula (5).

D is meant to be the depth of water in the river at a given section corresponding to a value of x; in fact it varies at every point across and along the river. Hydrographic charts show a fair collection of depth samples which have been reduced to an arbitrary datum. The depths shown are "safe" depths in the sense that the actual depth of water at any given time will usually exceed the depth shown on the map. Where there is a tide, the hydrographic depth will be the depth of water at one of the lowest low water; where there is no tide, the depth shown is related to the average water level. The first step in order

to obtain a manageable value of the mean depth D which can be utilized in our calculations is to divide the river bed into short sections of fairly uniform depth and width and average the hydrographic soundings over it. The gauge readings at a particular station like King Edward Pier, Frontenac and so on, give the height of the water surface above a given reference level at this particular location. It varies all the time and we have to take its average to compare it with the average depth obtained from the chart soundings for the channel around the station. For instance, the 10 year average at the gauge site of King Edward Pier gives 6.62 meters while the average depth of water in the corresponding channel turns out to be 6.19 meters. We establish the correspondence:

6.62 meters at King Edward implies a depth of water of 6.19 meters in the corresponding channel and so on.

In 1969, the mean annual level read on the gauge at King Edward was 6.86 meters. I surmise that the depth of water in the channel was 6.44 meters. In the regions where there is a tide, the average depth deduced from the soundings should be increased by the mean half range of the tide strictly speaking. In practice though, the actual mean soundings may be kept as a useful representative quantity since taking account of the change of depth during the tidal cycle introduces second order effects which it is not the purpose of this paper to investigate. We are more interested in having first a more general view of the basic mechanism present in the river.

The quantity I is even more difficult to extricate from the gauge readings and the soundings. At King Edward, the

10 year average gauge reading indicates that the surface of the water at this station stands at 6.62 meters above mean sea level (which we assume to coincide with IGLD). From the soundings, we had calculated a mean depth of 6.19 meters for the portion of the river bed enclosing King Edward; we deduce that the river bed stands at .43 meters above msl at this location. The Sorel gauge indicates a 10 year average of 4.69 meters while the average of the soundings in the vicinity gives 7.40 meters; we conclude that the river bed lies at 2.71 meters below msl. In this fashion we may obtain the mean inclination of the river bed between King Edward and Sorel. The distance between the two stations is 68,600 meters and therefore: (see Fig. 2):

$$I=3.14/68600 = 4.6 \times 10^{-5}$$

With the same information we deduce

$$-h_X=S_X=1.8/68600=2.6x10^{-5}$$

We notice that the river bottom slopes more rapidly than the river surface in this portion of the river. In practice we are not too concerned with I since it cancels out of (5).

#### 4. The inclination of the water surface

The gauge readings give us a direct indication of  $S_X$ , the inclination of the water surface. Such gauge readings are published by the Canadian Hydrographic Service under the Tides and Water Levels Section. We consider that under average

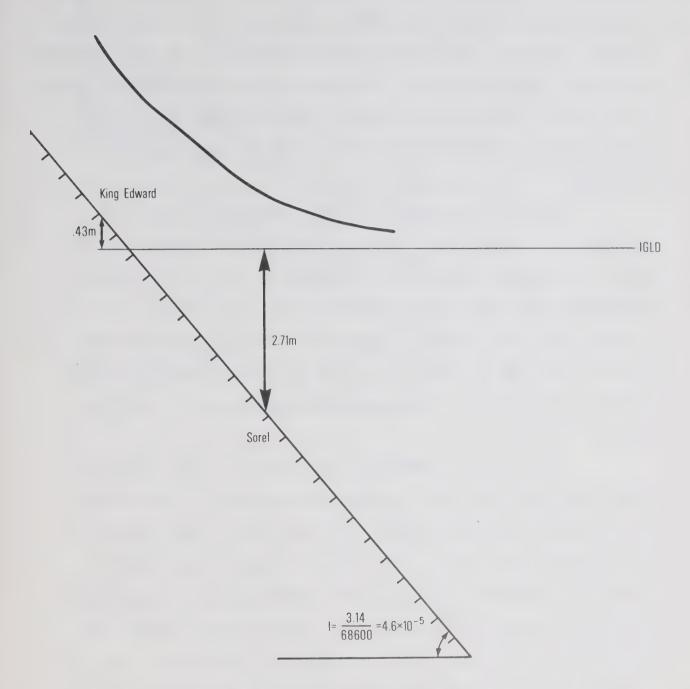


FIG.2 THE INCLINATION OF THE RIVER BED

conditions steady state prevails between Montréal and Sorel; therefore we calculate  $S_X$  over this portion of the river. Fig. 3 shows the mean profile S between the years 1962 to 1969; the rate of change of this profile with x gives  $S_X$ . We have entered with each curve the value of the mean yearly discharge  $Q_O$ ; there is an adequate correlation between the discharge  $Q_O$  and the overall depth of water in the river, but it certainly is not absolute.

What is remarkable in Fig. 3 is that in spite of D varying quite appreciably with  $Q_0$ , the profile itself seems little affected. It differs from point to point like between King Edward and Sorel and between Lanoraie and Sorel, but it changes only very slightly from year to year.  $S_X$  oscillates between 2.48 and 2.74x10<sup>-5</sup> for 1964 to 1969 between Frontenac and Lanoraie while D itself undergoes variation of over one meter.

Would indicate that  $Q_0^2/D^3$  is nearly a constant. This may be ascertained by plotting  $Q_0$  vs  $D^3/2$  for some of the stations. Fig. 4 shows this type of plot for Frontenac, King Edward, Lavaltrie, Lanoraie and Sorel from observations on the annual mean taken between 1962 and 1969. The horizontal scale for Lavaltrie is different and this has been indicated on the diagram. This fact in turn indicates that the depth varies directly with the discharge according to the relation:



FIG. 3. SLOPE OF THE WATER SURFACE  $S_x$ 

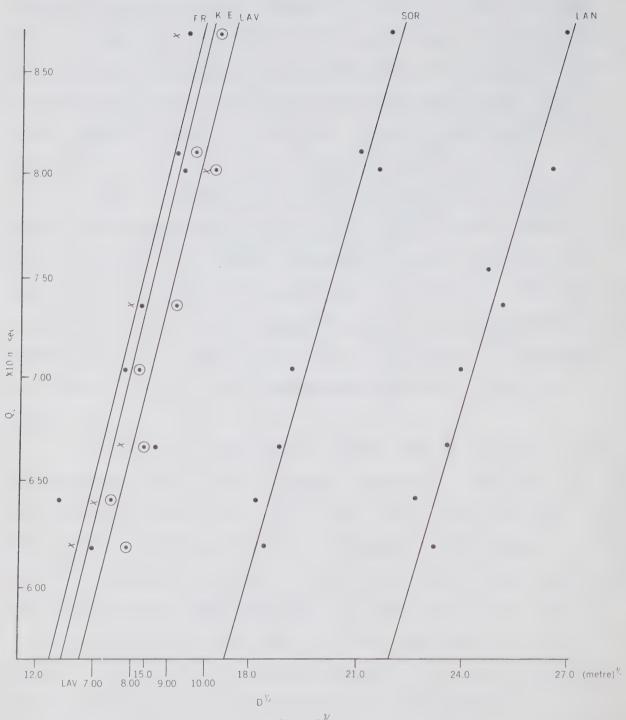


FIG. 4 PLOT OF  $\mathbf{Q}_0$  VERSUS  $\mathbf{D}^{\frac{3}{2}}$ 

#### 5. The mean level at various stations

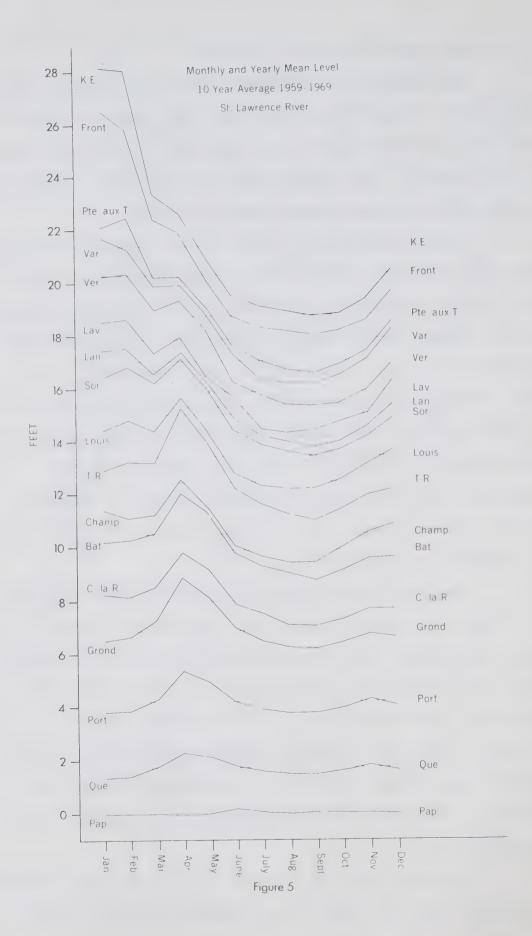
The gauges also indicate directly the water levels at the points where they are located. The mean value of the levels has a climatological value and Fig. 5 shows this mean level on a monthly and yearly basis for a 10 year average at the principal stations strung along the river.

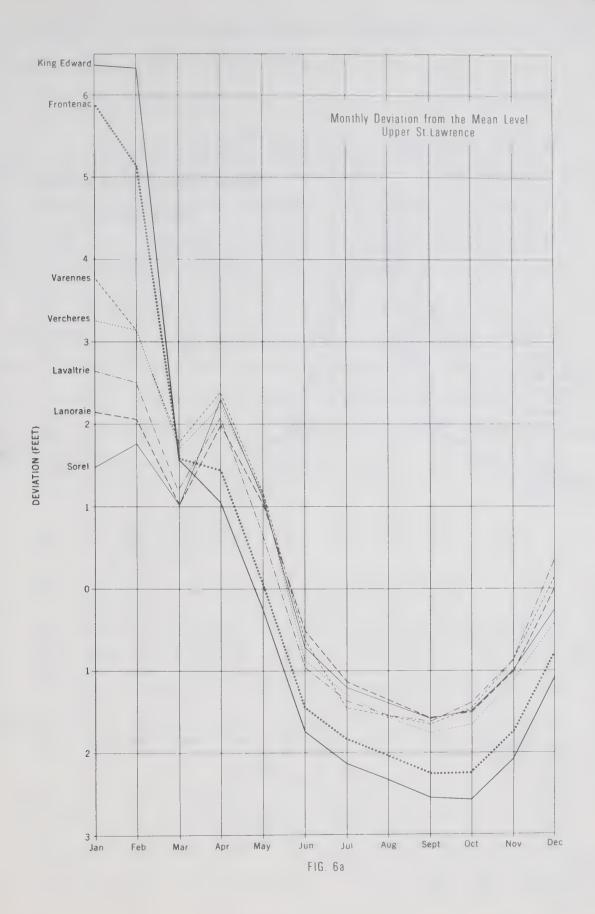
The fluctuations in mean levels decrease gradually downstream till they become nearly imperceptible at Pointe au Père. At King Edward and Frontenac, there is a large drop in level in January and February following the freeing of the channel from the ice. The next feature common to all the curves is the increase in level in March and April created by the melting of the snow cover. These variations in level are maximum between the mouth of Lake St. Peter and Québec.

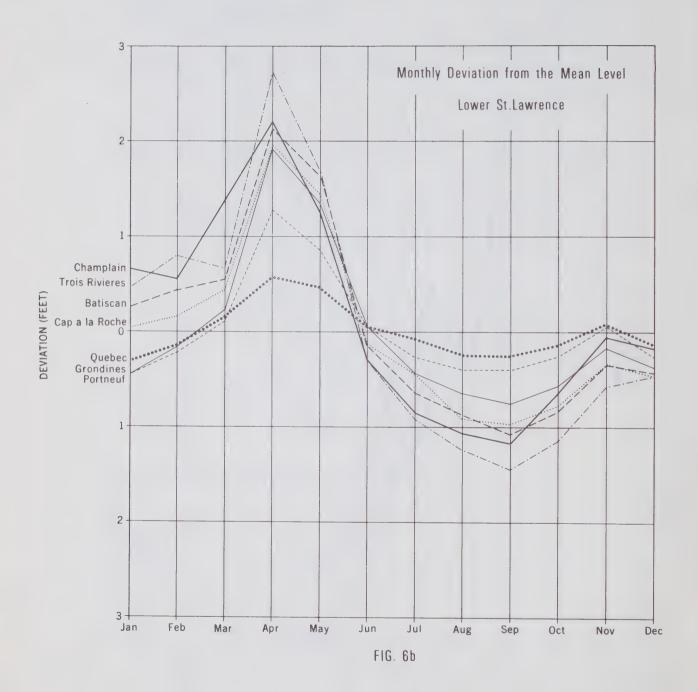
Afterwards, the level drops uniformly to reach a minimum. value in September when they start increasing on account of the increased precipitation.

Figs. 6a and 6b show the deviation from the mean yearly level of the monthly levels. As expected these deviations enclose each other like an envelope since the bed of the river expands downstream and is increasingly less affected by the fluctuation in the discharge. The only exception to the rule is Trois Rivières during April which exhibits then larger deviations from mean level than the stations on the other side of Lake St. Peter such as Sorel and Lavaltrie.

The mean level may deviate by as much as 1.3 meters in Montréal while at Québec its deviation seldom exceeds 20







centimeters. Ice free conditions in Montréal should reduce the fluctuations in level appreciably.

#### 6. The Chézy coefficient

The term  $Q_0^2/c^2B^2D^3$  is meant to represent the effect of friction; it does so but very imperfectly. It is adequate for the purpose of our investigations. It depends fundamentally on the value of C, the Chézy coefficient. There exist theoretical estimates for this parameter (Dronkers: Tidal Computations p. 157) but in practice it has to be determined from observations. (5) may be utilized for this purpose since  $S_X$  and B are known and since we have established that  $Q_0/D^{3/2}$  is approximately a constant of value 483 m $^{3/2}$ /sec according to Fig. 4. Between Montréal and Sorel,  $S_X \sim 2.6 \times 10^{-5}$  and the average width of the river is 1850 meters so that

 $c \sim 53 \text{ m}^{1/2/\text{sec}}$ 

# 7. The rate of change of depth with distance, $\text{D}_{\text{X}}$

According to (3),  $D_X$ , the change of depth of water as one moves downstream may be positive, negative or zero depending on the relative magnitude of the two quantities on the right hand side.

 $D_{\rm X}{<}0$  over a longish stretch of the river would imply that the depth of water would eventually reach zero, a situation which definitely does not occur in the St. Lawrence River.

 $D_{\mathbf{X}} \! > \! 0$  implies that D increases steadily downstream so that in turn the friction term on the right hand side diminishes

steadily and the  $D_X \rightarrow I$  for large x which means that the surface of the water eventually becomes horizontal. An inspection of Fig. 1 shows that this is what actually occurs in most rivers and in the St. Lawrence in particular.

The special case  $D_X=0$  corresponds to the situation considered by Chézy in which the depth of water is constant throughout the channel. Experiments seem to indicate that such a state of flow is unstable. (Dodge and Thompson: Fluid Mechanics. p. 247). This special case is illustrated by the dotted line in Fig. 1. In this particular instance  $S_X=I$  and therefore

$$S_{x} = Q_{0}^{2}/C_{B}^{2}D^{3}$$

which in turn implies

$$S_X^{\frac{1}{2}} \alpha Q_O$$

I did plot  $S_X^{\frac{1}{2}}$  vs  $Q_O$  in the upper reaches of the St. Lawrence; I could find no correlation at all between  $S_X^{\frac{1}{2}}$  and  $Q_O$ . It is therefore unlikely that  $D_X=0$  holds anywhere in that area.

8. The convective term and the critical velocity

We have dismissed the convective term in (2) as being negligible in the St. Lawrence. We may now verify this.

Throughout the paper we schematize the river bed as having a rectangular section of width B and depth D. We may then write

Qo=BDuo

uo being the mean velocity of the current along the channel under consideration. The convective term becomes

$$u_0^2 D_x/gD$$

assuming that only the depth varies downstream. We had noted that between Montréal and Sorel

$$I \sim 4.6 \times 10^{-5}$$
  $S_X \sim 2.6 \times 10^{-5}$   $D_X = I - S_X \sim 2 \times 10^{-5}$ 

 $u_{O}$  nowhere exceeds 2 m/sec and D  $^{\circ}$  6 m. so that

$$u_0^2/gD \sim 1/90$$

so that the convective term

$$u_{O}^{2}D_{X}/gD \sim D_{X}/90 << D_{X}, I$$

Similarly

$$Q_0^2/C^2B^2D^3=u_0^2/C^2D=(u_0^2D_X/gD) \cdot (g/D_XC^2)$$

Now  $g/D_xC^2 \sim 181$  so that  $Q_0^2/C^2B^2D^3 \sim 181$  times the convective term. The convective term cannot always be neglected. For larger velocities than those quoted for  $u_0$  it has to be retained

in equation (2). An analysis of the complete equation (2) including the convective term indicates that  $D_{\mathbf{X}}$  steepens to infinity for

 $u_0 = (gD)^{\frac{1}{2}}$ 

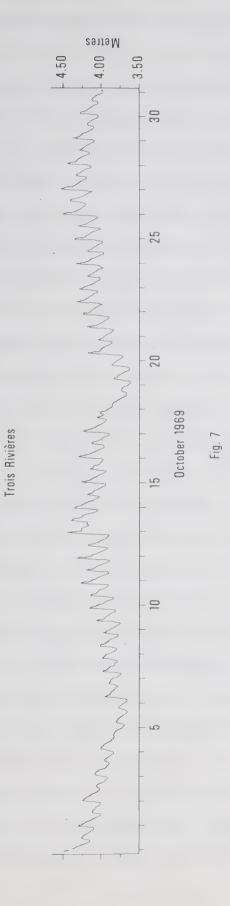
which is known as the critical velocity. At this velocity a hydraulic jump may occur and waves dowstream are effectively prevented from moving upstream. The velocity of the currents in the St. Lawrence is well below the critical velocity; this is the reason why the tide can move so far into the river.

#### TIDAL MOTION

1. The tide in the St. Lawrence River

From Trois Rivières at the mouth of Lake St. Peter, onwards to the Gulf of St. Lawrence the river flow and the water level become increasingly affected by the tide.

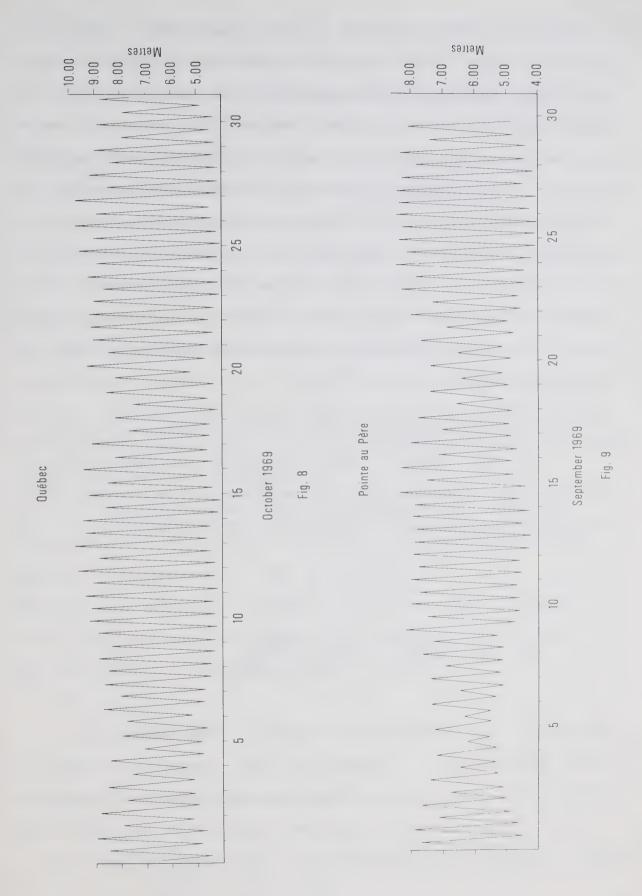
First, the modifications are slight and are appreciable only over intervals of two weeks or more. Fig. 7 shows the changes in level during a month at Trois Rivières. The daily tidal fluctuations ride on the back of a more important semimonthly oscillation. Downstream of Trois Rivières the diurnal and semidiurnal oscillations become predominant. The change in level is the shape of saw teeth when plotted on a time scale as can be seen for instance in the plot of  $\Delta S$  in Fig. 12 with a steep and rapid increase in height followed by a more gradual lowering of the level till the cycle is repeated with varying



intensity (upside down). Such a curve necessitates many harmonics to describe it.

The tide reaches its maximum amplitude in the vicinity of the Isle aux Couldres; downstream it decreases and takes a more symmetrical oceanic character. Figs. 8 and 9 show the change in level at Québec and Pointe au Père during one month. Québec is a station where the tide has a shallow water character although the tight time scale of the diagram does not allow to fully realize this fact (look instead at Fig. 12) while at Pointe au Père the tide has a smooth oceanic character.

The surface slope  $S_X$  studied in the previous section is no longer constant and it varies from instant to instant. In contrast to (4), it may become negative i.e. the water level is higher downstream than upstream, for instance before high water. The river flow Q as well, becomes an oscillating quantity and below a certain point, located in the vicinity of Grondines and Batiscan, the water may flow upstream during appreciable intervals of time. As the capacity of the river bed increases downstream, the steady flow Qo contributed by the discharge of the river and its tributaries becomes imperceptible compared to the oscillating flow Q created by the tide. There is an intermediate zone though, which stretches between Trois Rivières and St. François (Isle d'Orléans) where the river discharge Qo and the oscillating flow are of the same order of magnitude. In this area Qo and Q will interact with each other through the influence of the friction represented by the Chézy term in the equations of motion. The tide will be damped by the friction of



the river bed and its progress will be slowed down by the discharge. High and low water will move more and more slowly upstream and they rapidly decrease past the constriction of Québec City; the interval between low and high water will be shortened while the interval between high and low water will be increased. The interaction between the tidal constituents creates slower oscillations of a period of two weeks and one month which increase in amplitude upstream and eventually dominate the whole tidal motion in the vicinity of Trois Rivières. Eventually these slow oscillations and the more rapid diurnal and semidiurnal oscillations become extremely weak upstream of Lake St. Peter.

2. The Equations of Hydrodynamics for time dependent motion

A fair amount of the motion in the St. Lawrence where the tide is of importance may be described and understood with the help of the simplified Equations of Hydrodynamics:

$$Q_X + BDt = 0$$
 (6)

$$I - D_{X} - \frac{Q|Q|}{C^{2}B^{2}D^{3}} = \frac{Q_{t}}{g BD} + \frac{2QQ_{X}}{g B^{2}D^{2}}$$
 (7)

(Dronkers: Tidal Computations 196-197). The symbols have already been defined. t stands for time. Equations (6) and (7) reduce (1) and (2) when the time dependent terms are dropped.

Our first task is to investigate which of these terms may be evaluated from observations. The flow  ${\tt Q}$  is the total

flow across a vertical section of the river; it embodies the river discharge Qo and the oscillations due to the tide. The flow can only be measured directly with the help of current meters; this instrument measures the current at a given point. To obtain the flow Q from such current measurements it is necessary to investigate the horizontal and vertical distribution of the current in the vertical section which unfortunately differs for each stage of the tide and to do a space integration. A sensible technique for getting an idea of Q is to make a set of preliminary measurements at many points across the section at various depths and try to locate the point where the current is approximately equal to the average current going across the section. A self recording current meter can be installed there to study the time variation of Q from observations of the current at a single point of the section.

The width B and depth D can be deduced from the schematization of the river and readings of the water level gauges. The gradient in flow  $Q_{\rm X}$  can only be found if values of the flow are measured at sections which are closely spaced, a costly undertaking. The equation of continuity (6) allows us to re-express  $Q_{\rm X}$  in terms of Dt since

$$Q_X = -B D_t$$

Therefore, the useful form of the equation of motion is

$$S_{X} - Q|Q| = \frac{Qt}{c^{2}B^{2}D^{3}} = \frac{Qt}{gBD} - \frac{2QDt}{gBD^{2}}$$
(8)

 $\mathbf{S}_{\mathbf{X}}$  can be inferred immediately from the differences in readings between the various tide gauges strung along the river.

2.1 The order of magnitude of the terms contained in the Equations of Hydrodynamics

Keeping in mind that all the quantities present in (8) vary at each instant of time, it is still desirable to compare the order of magnitude of the terms present in (8).

If we use the observations of the Ship Channel Division at Québec Bridge in 1968 we notice that the difference in level between Basile Low Light and Québec Bridge oscillates between + and - 40 cm over a distance of two nautical miles (3.7 km), so that

$$S_{\rm X} \sim 40/3.7 \times 10^5 \sim 10^{-4}$$

To estimate the maximum values of the remaining terms we use the fact Q=BDumean

and

 $C=52 \text{ m}^{\frac{1}{2}}/\text{sec}$ 

B=1650 m

D=21 m

 $u_{\text{max}} = 2m/\text{sec}$ 

 $u_{t_{max}} = 2m/sec/2$  hours

Dt=3m/6 hours

then

$$\frac{Q|Q|}{C^2B^2D^3} = \frac{u|u|}{C^2D} \sim 7x10^{-5}$$

 $Qt/gBD=ut/g \sim 10^{-5}$ 

 $2QD_{t}/gBD^{2} = 2uD_{t}/gD \sim 3x10^{-6}$ 

The terms are of the same order of magnitude with the exception of the convective term 2uDt/gD which is smaller than the rest. The instantaneous values of the terms considered may be vastly different from the values quoted but here we have tried to estimate their maximum possible values.

### 2.2 An elementary solution

Equations (6) and (7) cannot be solved analytically; still there is little point in embarking into involved numerical computations before getting some insight into the physical situations they may describe. We may drop the convective term without more ado, now that we have convinced ourselves that it is relatively small. We will replace the troublesome friction term by a linearized term of the form rQ; this is a mathematical artifice which renders (6) and (7) tractable analytically although this excludes from our possible solutions, physical situations which we will have to face eventually such as the creation of shallow water constituents.

(6) and (7) become in their simplified form

$$Q_{x} + BD_{t} = 0 \tag{6a}$$

$$I-D_{X}-rQ = Q_{+}/qBD \tag{7a}$$

The D at the denominator of (7a) makes things unmanageable as well and we treat it as a constant depth  $D_0$ . This latter approximation is adequate in relatively deep water where  $D_X$  and  $D_t$  are small compared to the mean depth of the water column  $D_0$ ; finally we drop I compared to  $D_X$  since retaining it increases the analytical manipulations without improving our physical understanding. We are now left with

$$Q_X + BD_{\dagger} = 0 \tag{6b}$$

$$D_X + rQ + Qt/gBD_O = 0$$
 (7b)

At this stage, the time and space dependence of D and Q can be treated as separate factors of the form X(x)T(t) and since we work with tides which are periodic oscillations, we write

We consider a channel or rectangular section of constant depth  $D_{\rm O}$  with width  $B_{\rm O}$ . The solution of (6b) and (7b) for such a channel are

where

$$K = \sqrt{i\sigma B_0 (r + i\sigma/gBD_0)} = \sqrt[4]{\sigma^2 B_0^2 r^2 + \sigma^4/g^2 D_0^2} \exp \left[i\frac{1}{2}arc \tan B_0 rgD_0/\sigma\right]$$

K has the form  $K=\alpha+i\beta$  and  $e^{\pm}Kx$  contains a factor  $e^{\pm}\beta x$ .  $\beta$  depends essentially on the friction r and it would vanish if r vanished. Linear friction automatically creates an exponential damping in the elevation and currents associated with the propagation of the tide (the exponential increase does not make physical sense). The  $\alpha$  part of K may be interpreted as a wave number; without linear friction it would be  $\sigma$  /  $(gD_O)^{\frac{1}{2}}$ . The wave length would depend on the frequency  $\sigma$  and the depth  $D_O$  of the channel. With the addition of linear friction the wave length depends also on r; since  $\sigma^2 B_O^2 r^2$  is essentially positive, the wave length is shortened with the occurrence of linear friction.

We could write out explicitly the solution to (6b) and (7b) using the basic relations we have just derived. There is little point in doing this as these solutions are given in innumerable textbooks (Dronker: Tidal Computations p. 225 ff. Defant: Physical Oceanography Vol. 2 p. 142 ff. Stoker: Water Waves pp. 37-54). All that we wish to retain is that linearized friction implies an exponential decay in a tidal wave progressing upstream and a shortening of its wave length.

The solution of (6b) and (7b) is possible for more complicated channels such as a channel whose width and depth vary exponentially, linearly or quadratically. Besides the analytical complications, the physical results are the same. Linear friction damps the amplitude and currents associated with the travelling tidal wave and there occurs a shortening of the wave length.

2.3 The creation of shallow water constituents and their selective damping up the river

We now have to look at the new features brought by the non-linear friction term. A frontal attack cannot do as this term makes (6) and (7) intractable, but one may get some information by using some devious means. We substantiate, with the help of a little bit of algebra, the two following statements:

- 1) the shallow water constituents are created by the interaction of the tidal constituents with each other and with the river discharge through the friction term  $u|u|/C^2D$
- 2) The low frequency constituents can travel further up the river than the higher frequency constituents before being damped away.

From Trois Rivières to Montréal, the current is one directional and the tidal streams become quite small compared with the steady current due to the river discharge. In this part of the river, we may therefore write:

$$u | u | /C^2D = u^2/C^2D$$

If we consider only two tidal constituents along with the river discharge, we get:

$$u=u_0+u_1\cos(\sigma_1t-\alpha_1)+u_2\cos(\sigma_2t-\alpha_2)$$
 (9)

and we assume

$$u_1^{<<}u_0$$
  $u_2^{<<}u_0$ 

The friction term becomes explicitely:

$$\mathbf{u^2/C^2D} = \mathbf{u_0^2} + \frac{1}{2} \left( \mathbf{u_1^2} + \mathbf{u_2^2} \right) + 2\mathbf{u_0u_1} \cos \left( \sigma_1 \mathbf{t} - \alpha_1 \right) + 2\mathbf{u_0u_2} \cos \left( \sigma_2 \mathbf{t} - \alpha_2 \right)$$

$$+ \frac{1}{2} u_1^2 cos2 \left(\sigma_1 t - \alpha_1\right) + \frac{1}{2} u_2^2 cos2 \left(\sigma_2 t - \alpha_2\right)$$

$$+u_1u_2 \{\cos[(\sigma_1+\sigma_2)t-(\alpha_1+\alpha_2)] +\cos[(\sigma_1-\sigma_2)t-(\alpha_1-\alpha_2)]\}$$
 (10)

Even taking into account of our very restrictive assumptions about the components of u, (9) still cannot be a solution to equation (8) if the friction term is given by (10). In an iteration process however, assuming that it will converge, we may consider (9) as a first order solution obtained without the friction term while an improved solution may be obtained by retaining the linear terms in (10) and by dropping the non-linear ones:

$$u_2/C^2D \sim u_0^2 + 2u_0u_1\cos(\sigma_1t - \alpha_1) + 2u_0u_2\cos(\sigma_2t - \alpha_2)$$
 (11)

The terms in  $u_0u_1$  and  $u_0u_2$  represent the interaction of the tidal constituents with the river discharge; the term  $u_1u_2$  represents the interaction of the tidal constituents with each other. If, for the sake of an example, we take as the two

tidal constituents,  $M_2$  and  $S_2$ , (10) indicates that their interaction through the friction would create shallow water constituents such as  $M_4$ ,  $S_4$ ,  $MS_4$  and MSf. The terms in  $u_0u_1$  and  $u_0u_2$  will modify  $M_2$  and  $S_2$  themselves. The use of (11) allows us to separate the solutions for  $u_0$ ,  $u_1$  and  $u_2$ . The solution for constituents 1 and 2 has the form:

u,h exp [iot 
$$\pm \frac{i \sigma}{(gD)^{\frac{1}{2}}}$$
 (1 - 2igu<sub>O</sub>/C<sup>2</sup>σD)<sup>1/2</sup> x] (12)

the amplitude of the wave moving up the river is modulated by the function

$$\exp \left[-i\sigma/(gD)^{\frac{1}{2}} (1 - 2igu_O/C^2\sigma D)^{\frac{1}{2}} x\right]$$
 (13)

which differs from the usual sinusoidal modulation factor:

$$\exp \frac{-i \sigma}{(qD)^{\frac{1}{2}}} x$$

for a wave travelling upstream without friction.

The argument of (13) has the form

$$-\underline{i} \sigma \qquad (\alpha - i\beta) \tag{14}$$

 $\alpha$  is akin to the wave number  $K=\sigma/(gD)^{\frac{1}{2}}$  but it is modified by the friction.  $\beta\sigma/(gD)^{\frac{1}{2}}$  reduces the amplitude of the wave exponentially.

The explicit expression for  $\alpha$  and  $\beta$  is

$$\begin{pmatrix} \alpha \\ \beta \end{pmatrix} = [1 + (2u_0) \frac{2g^2}{C^4\sigma^2} D^2]^{\frac{1}{4}} \quad \begin{pmatrix} \cos \\ \sin \end{pmatrix} \quad (\frac{1}{2} \text{ arc tan } 2u_0 g/C^2\sigma D)$$

$$\simeq (2u_0g/\sigma DC^2)^{\frac{1}{2}} \begin{pmatrix} \cos \\ \sin \end{pmatrix} (\frac{1}{2} \text{ arc tan } 2u_0g/C^2\sigma D) \tag{15}$$

for  $\sigma << 1$ , which holds for the tidal constituents.

We may insert some numerical values into (13) and check if the calculated phase retardation and the attenuation agree at all with the observations.

Between Lavaltrie and Trois Rivières, the mean depth D  $\sim$  6.6 meters, while between Trois Rivières and Batiscan D  $\sim$  6.2 meters. We take  $u_O \sim$  55 cm/sec, C  $\sim$  52m $^{\frac{1}{2}}$ /sec

$$\sigma_{\text{M}_2} = 1.41 \text{x} 10^{-4} / \text{sec}$$
  $\sigma_{\text{O}_1} = .73 \text{x} 10^{-4} / \text{sec}$   $\sigma_{\text{MSf}} = .049 \text{x} 10^{-4} / \text{sec}$ 

The distance between Lavaltrie and Trois Rivières is 82.4 kilometers while the distance between Trois Rivières and Batiscan is 31.9 kilometers.

Table 1 shows a comparison between the observed and calculated values of the phase retardation and of the attenuation.

The attenuation between Batiscan and Trois Rivières is larger than our estimate but it must be understood that the assumptions underlying (11) do not hold in this area. But even there our qualitative deduction that the damping of MSf should be much less than that of M2 and O1 still holds true.

Table 1

Phase Change and Attenuation Between Batiscan and Lavaltrie for  $\mathbb{M}_2$ ,  $\mathbb{O}_1$  and  $\mathbb{M}_2$ f

MSf	A $\beta$ $\beta$ ' $\alpha$	7.8 62.7 .8x10 <sup>-2</sup> .4	13.8 75.4	11.6 88.8 .6x10 <sup>-2</sup> .4	7.5 92.0 .7x10-2 .3				.6x10-2 .3	C
01	β' α' A km <sup>-</sup> 1 deg/cr km	5 3.5x10-2 1.1 17	10	1.7x10 <sup>-2</sup> 1.5	2.4x10-2 1.1	2.2x10-2 .9	.2 140.8 2.9x10-2 1.6		2.3×10-2 1.2	
	a' A ø' deg/ cm deg/ km	0.6 4.5x10-2 1.1 7.9 331.5 3.5x10-2 1.1 17.8 62.7 .8x10-2	2.6 7.1	.9 2.4x10 <sup>-2</sup> 2.8 1.4 58.1	10-2 2.2 .6 67.2	10-2 2.2 1.5 74.0	4.7 3.8x10 <sup>-2</sup> 2.7 .2 140.	eres and Lavaltrie	2.7x10-3 2.5	
M2	A $\phi$ $\beta'$ cm deg km <sup>-</sup> 1	32.9 32	7.9 356.4	2.5 89.9 2.4x	1.3 128.2 3.1x10 <sup>-2</sup>	1.3 157.2 1.6x10 <sup>-2</sup>	.6 214.7 3.8x	Trois-Rivières and	2.7x	
Station	∧ × ∧ X	Batiscan -31.9	Trois- Rivières 00	Louiseville 33.8	Sorel 58.1	Lanoraie 71.9	Lavaltrie 82.4	Average between Tro		

.5x10-2

1.8x10-2 1.1

2.4×10<sup>-2</sup> 1.7

ø=phase

A=Amplitude

 $\beta$  = log(A<sub>2</sub>/A<sub>1</sub>)/ $\Delta$ x

 $\alpha' = \Delta \phi / \Delta x$ 

Calculated

#### 2.4 The speed of propagation of the tide

A wave travelling in a frictionless channel should propagate at a speed:

$$c = (gD)^{\frac{1}{2}} \tag{16}$$

In the lower St. Lawrence the tidal oscillation is intermediate between a standing wave and a travelling wave. Upstream of Québec the tidal wave takes a definite progressive character.

In a frictionless channel when one is in presence of a mixture of standing and travelling waves, the crest of high water should always travel at a speed in excess of (16). If the crest, anywhere in the channel, travels at a speed less than (16), one is in presence of a damped travelling wave. The use of (12) confirms this fact at least qualitatively, keeping in mind the limitations of its validity. From (12) we deduce that the speed of displacement of a damped wave is

$$c' = (gD)^{\frac{1}{2}}/r$$
 (17)

where

r=1.89 for M2

=2.13 for  $O_1$ 

=8.17 for MSf

Therefore in the presence of friction, free waves travel at different speeds depending on their frequency and they always travel more slowly than an undamped wave.

Also, an inspection of the speed of travel of high and low water in the upper St. Lawrence indicates that the low water travels more slowly than the high water. The fact that the depth of water D in the channel is less at low water than at high water may explain this fact as the difference in depth between high and low water equals the range of the tide.

However, the low water travels much more slowly than would be predicted by simply inserting the appropriate value of the depth in (17) thus indicating that our algebra does explain many facts qualitatively but not quantitatively. A more refined numerical analysis is necessary in order to obtain some quantitative agreement.

The distance between Québec and Grondines is about 93 km and the crest of high water takes an average of  $2\frac{1}{2}$  hours to reach Grondines from Québec; therefore, the high water travels at about 10.3 m/sec. The average depth between the two localities is 12 m. and therefore the speed of the free wave in the absence of friction should be c=10.9 m/sec. The actual wave travels more slowly than the free wave and it therefore consists of a damped travelling wave.

To conclude, we give a table of the time of travel for the high and low water at neaps and springs between various stations (Table 2). These averages are based on the occurrences during April and October, 1964. The low water at neap is higher than the low water at spring; this may explain why it travels appreciably faster. The table indicates that the high water at spring reaches Québec before St. François on the Island of

Table 2

Average Time of Travel of the High and Low Water Between Pointe au Père and Various Stations Upstream in Hours.

	A	В	С	D	E	F	G	Н
High Water Spring	0	4.3	4.1	5.5	5.9	6.6	6.7	7.3
Neap	0	4.0	4.3	5.4	6.0	6.8	6.8	7.9
Low Water Neap	0	4.1	5.1	6.4	7.4	8.4	8.8	10.3
Spring	0	4.8	5.9	7.4	8.1	9.3	9.5	11.2

- Α Pointe au Père
- В St. François
- Québec C
- Neuville D
- E Portneuf
- F Grondines
- Cap à la Roche Trois Rivières G
- Н

Orleans. This is hard to believe and even more difficult to explain, but many observations confirm this fact.

# 2.5 The critical points $\Delta S=0$ , u=0, $u_t=0$

For a channel of constant width B and depth D, (8) reduces to

$$S_{x}=u|u|/C^{2}D+ut/g - 2uDt/gD$$
 (19)

u being the mean current across the section.

 $S_X$  may be written in finite difference as  $\Delta S/\Delta x$ , where  $\Delta x$  is the distance between two tide gauges while  $\Delta S$  is the difference in level between them. When  $\Delta S=0$ , the water level is horizontal all along the portion of the channel considered and this implies  $S_X=0$ . From (19) we deduce that if  $S_X=0$ , the friction term  $u|u|/C^2D$  should approximately balance the acceleration term ut/g, neglecting the convective term. In theory therefore one may deduce the Chézy coefficient C by evaluating u and ut at the instant when  $\Delta S=0$ . In practice though it is not easy to get an accurate value of ut.

Finally when  $u_t=0$ , the currents run at their maximum velocity either at flood or ebb and there should be an approximate balance between the surface gradient  $S_X$  and the friction term  $u \mid u \mid /C^2D$ ; this case gives us our only real chance to evaluate C when there is tidal motion.

2.6 A review of the current observations carried out by the Ship Channel Division and G.C. Dohler

There is now some material available for the study of (19), namely the work carried out by the Ship Channel Division near the Québec Bridge in 1968 and the current survey of the Upper St. Lawrence carried out by G.C. Dohler in 1960.

In finite difference form (19) becomes

$$\frac{\Delta S}{\Delta x} - \frac{u |u|}{C^2 D} = \frac{\Delta u}{\Delta t} \frac{1}{g} - 2 \frac{\Delta D}{\Delta t} \frac{1}{gD}$$

We obtain values for  $\Delta S$ ,  $\Delta D/\Delta t$ , and u by setting up tide gauges at both ends of the channel and a current meter in the middle where we assume that we effectively measure Q.

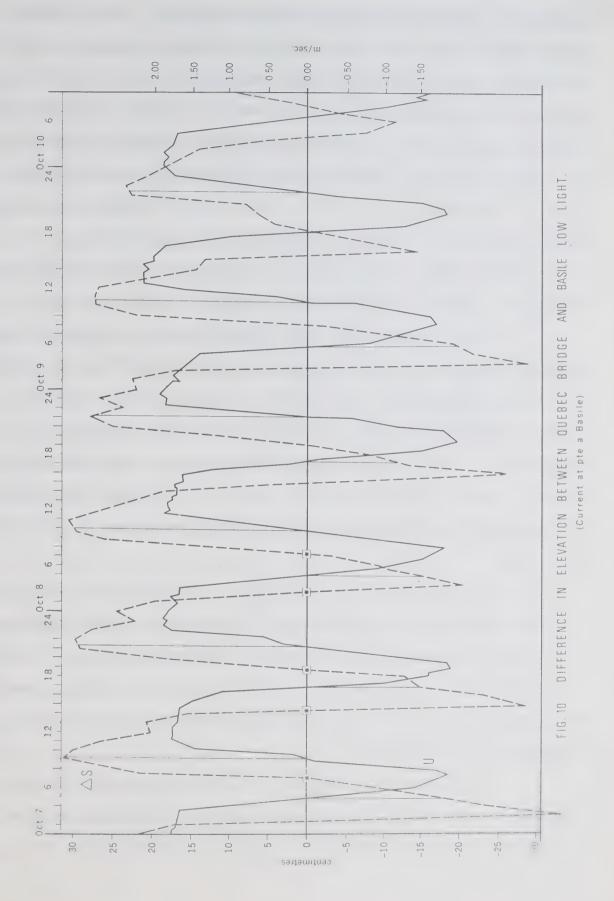
In 1968, the Ship Channel Division set up two tide gauges, one at Basile Low Light and the other at Québec Bridge, located exactly two nautical miles apart; they also installed one at Pointe à Basile, located halfway between the two stations. At the same point they installed a self recording current meter at the appropriate depth. Such a set up is the ideal for our present investigation. During the interval of observations the current meter worked quite adequately but unfortunately the results obtained from the tide gauges are questionable.

The tide gauge installed at Pointe à Basile which at first sight looks superfluous could have been used to check the other gauges but unfortunately its indications cannot be used at all: the high and low water which it records occur one or two hours after those at Basile Low Light and Québec Bridge. There

is something basically wrong with the clock mechanism or the abstraction of the data. The gauge at Basile Low Light seems to have worked properly during the interval of interest but I have the definite impression that the gauge at Québec Bridge was damped and had an appreciable time lag behind the true water level. This has unpleasant consequences as the quantity  $\Delta S$ , the difference in level between Basile Low Light and Québec Bridge is very slight, seldom exceeding 30 cm, and a small error in one or the other water level suffices to make it very nearly impossible to verify (20) in this portion of the river where it would have been extremely interesting to put it to test.

Although the two stations are barely two nautical miles apart there is an appreciable difference in mean level between them: the observations give a mean of 64 cm at Basile Low Light for October, 50 cm at Québec Bridge while the mean level at Québec (Lauzon) was 40 cm. The difference of 14 cm between Basile Low Light and Québec Bridge seems a bit high to me and there might be an error in levelling at one or the other gauge sites considered.

I did plot  $\Delta S$  and u for this portion of the channel on a time scale for a few days in October when the two gauges seemed to work not too badly, although one of the gauges gave some indication of drifting. The plot of u and  $\Delta S$  is shown in Fig. 10. The current curve is sensible enough. We notice that the interval of flooding is shorter than the time of ebb and that the current runs at maximum flood only for a very short time whence it turns suddenly to ebb during which it flows at



nearly the same velocity for about four hours. The current curve exhibits the saw tooth structure which we mentioned in the paragraph on the tide in the St. Lawrence River. The difference in level  $\Delta S$  between Québec Bridge and Basile Low Light has the proper shape as well; the level at Québec Bridge is higher than the level at Basile Low Light when flooding is about to take place. However, we will see that the  $\Delta S$  curve does not seem right on the time scale; it appears to be displaced by a few hours on the left, a fact which I attribute to the lag and the distorted response of the gauge located at Québec Bridge.

Slacks occur when the current curve u intersects the zero abscissa; these points have been marked in Fig. 10. The friction terms  $u|u|/C^2D$  vanish then and the surface gradient  $\Delta S/\Delta x$  should balance the acceleration term  $(1/g)\Delta u/\Delta t$ . By inspecting the diagram we notice that  $\Delta u/\Delta t$  is approximately the same for two consecutive slacks while the corresponding  $\Delta S$  is quite irregular, especially for the slack water which precedes the flood.

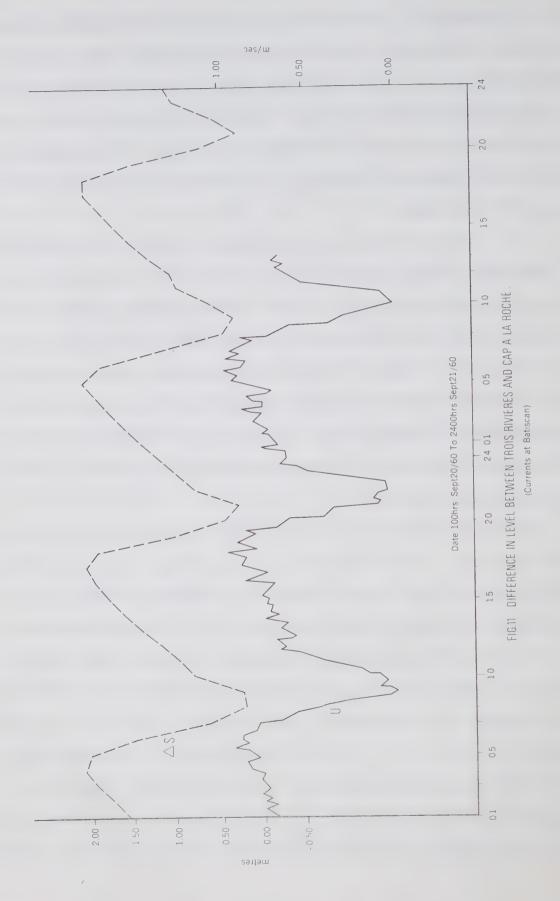
At maximum current, either at flood or ebb,  $\Delta u/\Delta t=0$  or is very small; there should be an approximate balance between the surface gradient  $\Delta S/\Delta x$  and the friction  $u|u|/C^2D$ . In Fig. 10 we notice that at maximum flow at ebb,  $\Delta S$  has pretty regular values while at flood it is virtually impossible to obtain consistent values of  $\Delta S$ .

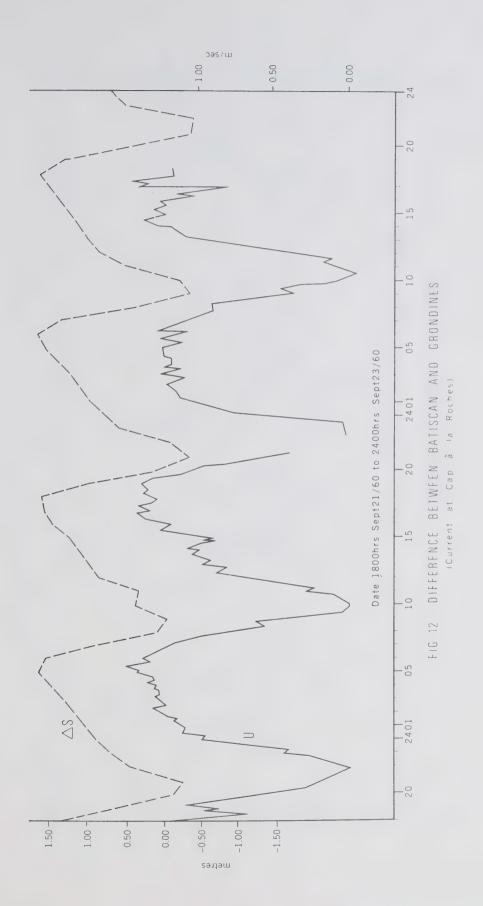
Finally when there is no surface gradient ( $\Delta S/\Delta x=0$ ), the acceleration (1/g) $\Delta u/\Delta t$  should balance the friction term  $u|u|/C^2D$ . At flood,  $\Delta S=0$  coincides with  $\Delta u/\Delta t<0$  while the

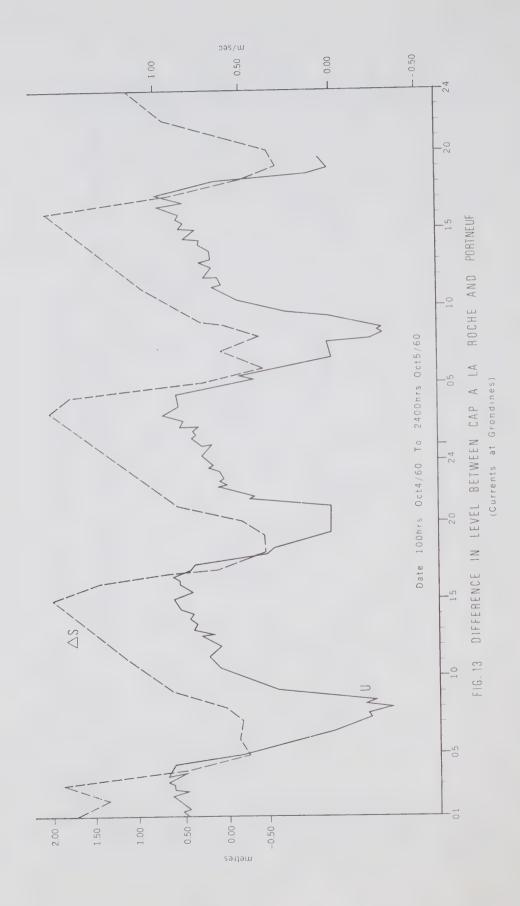
friction term is positive; therefore there is not even agreement in sign. Naturally  $\Delta u/\Delta t$  varies very rapidly near maximum flow at flood but still the general lack of balance between the various terms of the equation of motion suggests to me that one of the two gauges was not working properly.  $\Delta S=0$  should fall after the maximum flow at flood and not before as the record seems to indicate.

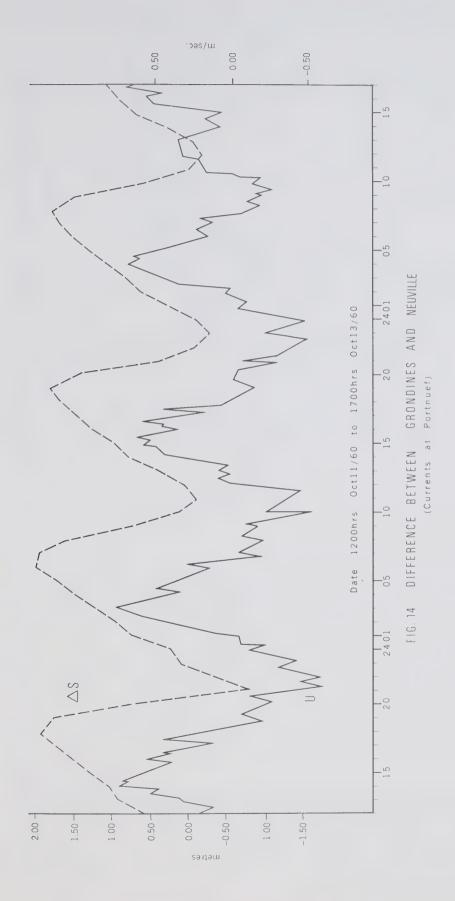
Now we turn our attention to the set of observations collected by G.C. Dohler and his assistants contained in the publication "Current Survey St. Lawrence River 1960". The currents were measured manually with Ott and Neyrpic type of instruments at various depths and locations; the duration of the observations seldom exceeded 48 hours and the observations showed marked fluctuations. Many of these observations are not contained in Mr. Dohler's publication and they have been kindly communicated to me.

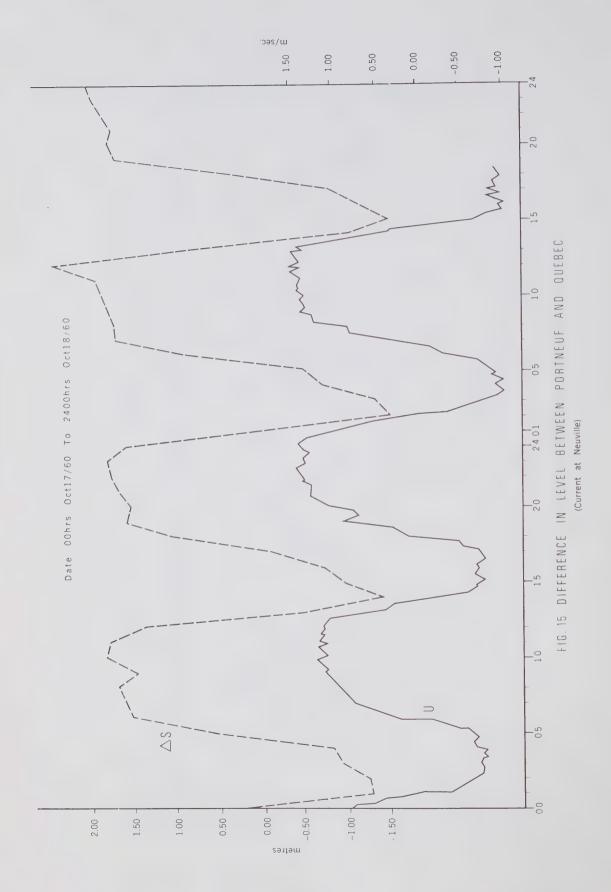
In order to test the equations of motion with such observations I used the hourly values of the water levels at the stations for which tides gauges are operated. These gauges sometimes are more than 50 km apart and the approximation to (19) by finite differences cannot be too good. The current observations show so many irregularities that it is usually impossible to obtain even a fair estimate of  $u_t$  and at times even of  $u_t$  when  $u_t=0$ . Still I found that the observations were of some use at the points  $u_t=0$ ; I have plotted  $u_t$  and  $u_t$  on a time scale and these are shown in Fig. 11 to 15.  $u_t$  is deduced from the records of hourly heights kept in the records of the











Tides and Water Levels section for the following stations:

Trois Rivières

Batiscan

Cap à la Roche

Grondines

Portneuf

Neuville

Québec (Lauzon)

The current plotted is the current measured by Dohler and his group at a point intermediate between the two stations. For instance in Fig. 11, we show the currents measured at Neuville on October 17 and 18, 1960 and we plot them along with the difference in level observed between Portneuf and Québec. Only at Neuville are the currents observations smooth enough to allow us to make some deductions from the case  $\Delta S=0$ . Table 3 lists the values of the Chézy coefficient we could deduce from the observations of Dohler.

The values of the Chézy coefficient which we deduce are rather consistent; the only exception is Grondines. Near Grondines, the river bed experiences a marked bend and constriction and there is a radical modification of the tidal regime in this area. Therefore, the high values of friction deduced for this area might be true.

# 3. Spectral Analysis

An analysis reveals the amplitude (and in tidal analysis, the phase as well) of all the oscillations contained in a sequence of observations over the range of frequencies

Table 3

Values of the Chézy Coefficient Deduced for the Section of the River Between Trois Rivières and Ouébec.

River Between Trois Stations	Rivières and Q Current at		Critical Point	Cflood m <sup>1</sup> /sec m	
Portneuf-Québec	Neuville	62	$u_t=0$	56	62
			∆S=0	60	62
Grondines- Neuville	Portneuf	49	u <sub>t</sub> =0	62	56
Cap à la Roche- Portneuf	Grondines	33	$u_t=0$	52	51
Batiscan- Grondines	Cap à la Roche	24	u <sub>t</sub> =0		63
Trois-Rivières- Cap à la Roche	Batiscan	54	u <sub>t</sub> =0		58

Average  $C_{\text{mean}}$  is 58  $m^{\frac{1}{2}}/\text{sec.}$ 

investigated. If the oscillations have approximately the same amplitude for any frequency, we are in presence of white noise which is of little interest to us. Whenever the tide is present in a record, some oscillations of given frequencies will dominate the spectrum, these frequencies lying in the semidiurnal and diurnal bands. Along the river, the phase of such oscillations will vary gradually.

The low frequency band of current and water level observations is filled mainly with noise; it contains the slow fluctuations of the steady current or the weekly and semimonthly variations of the mean level. In this band the resolution between the various oscillations becomes increasingly difficult. I personally do not subject this part of the spectrum to a harmonic analysis. I prefer to have a visual inspection of the low pass first; a power spectrum afterwards will indicate which frequencies dominate in this part of the spectrum.

In the semidiurnal and diurnal bands of the spectrum the tidal constituents dominate. The main tidal constituent is M2 which has a period of half a lunar day; the next constituent in importance is S2 which has a period of half a solar day. Then come N2, K1 and O1. N2 is linked to the variation of the moon's distance from the earth which goes through a complete cycle each month. K1 and O1 are diurnal and they are relatively small in the St. Lawrence system.

The five constituents mentioned make up the bulk of the tide. They are of astronomical origin but they are felt in the St. Lawrence system through the cooscillation of the gulf

and the river with the tide present in the Atlantic ocean which imparts some of its energy through Cabot Strait and Belle Isle Strait. These constituents are modified by the friction and they interact amongst each other. The interference of M2 with S2 will create a new shallow water constituent, MSf, which has a period of two weeks, while M2 and N2 will interact to give Mm which has a period of one month. The interference of K1 and O1 creates Mf which also has a period of two weeks. Many more constituents are created by the interaction of M2, S2, N2, K1 and O1 such as M4, MS4, S4, etc., but in a river the slow constituents MSf, Mm and Mf have a particular importance as they are not as rapidly damped as the constituents with higher frequencies.

3.1 The data available for analysis in the St. Lawrence River

Observations on water levels in the St. Lawrence River have been carried out since the days of Bell Dawson in 1895. They are usually abstracted every hour as such a time step is appropriate to tidal analysis and they have been accumulated in numerous records. Since 1961 these records have been put on punched cards, microfilms and magnetic tape at the Tide and Water Levels section. As such they constitute a wealth of most precious information for the study of the mean level, the tide and its propagation, wind and storm surges, etc.

Current observations have been much more sporadic since these are so difficult to carry out. The observations of Dohler in 1960 although very useful in the study of the dynamics of the river are too short to be subjected to a meaningful analysis.

On the other hand the material accumulated by the Ship Channel Section in 1968 at Québec Bridge and Isle aux Couldres is extensive enough to yield to an analysis, and we give the results of our calculations in one of the paragraphs which follow. Dr. W.D. Forrester has made current measurements between Pointe au Père on the south shore and Pointe à Michel on the north shore in 1965 and the results of his work are contained in the publication "Currents and Geostrophic Currents in the St. Lawrence Estuary" BIO 67-5 1969. Cdr. W.I. Farquharson has carried out similar current observations in the vicinity of Pointe des Monts in 1963 (St. Lawrence Estuary Surveys. BIO 66-6).

#### 3.2 Water levels

observations on water levels which stretch almost continuously over many years. We may use such long series of observations to follow the slow variations in the mean level and to separate very sharply most of the tidal constituents. We consider first the low frequency spectrum of the following stations:

Pointe au Père
Rivière du Loup or Tadoussac
Québec
Trois Rivières
Cap à la Roche or Batiscan
Frontenac

These stations are located at strategic points along the river where the tidal regime is modified from an oceanic type to a

shallow water type to eventually become a straight river flow. Figures 16(a), (b) and (c) show the low pass of the hourly values of the water level at the stations mentioned for the years 1967-68-69. For every year we have put the stations in their proper sequence on the same time scale so that the common oscillations can be picked up at sight.

Out of the initial confusion, for instance at Pointe au Père, there are little but random oscillations, we may notice that during the calm summer months, semimonthly oscillations due mainly to MSf, start standing out sharply at Québec and reach their maximum amplitude in the vicinity of Cap à la Roche and Batiscan. A little bit of attention shows that these slow tidal oscillations appear even in Montréal during the summer months but greatly attenuated.

This qualitative assessment of the situation may be put on a more quantitative basis by computing the power spectrum of the oscillations of the low pass. Mr. K.B. Yuen of our organization, has kindly performed this calculation and his results are shown in Fig. 17 for Pointe au Père, Québec, Trois Rivières and Frontenac. The peaks indicate the predominant oscillations and they are labelled by the proper constituent when we happen to know it. We note an oscillation of 1 cycle/9.6 days, which is caused by the interference of  $S_2$  with  $N_2$  which we could call SN and an oscillation of 1 cycle/week which may be due to a further interference of the slow constituents MSf or Mf with themselves. This would explain as well the extra peak at 1 cycle/4.8 days which would depend on SN.

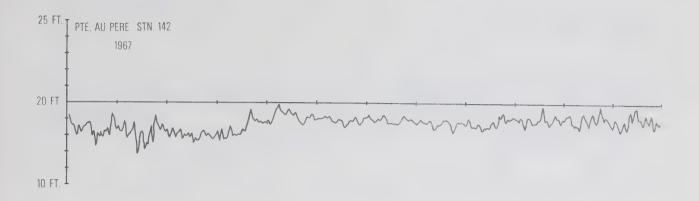






Fig. 16(a) Low Pass Filter of Hourly Heights, 1967.

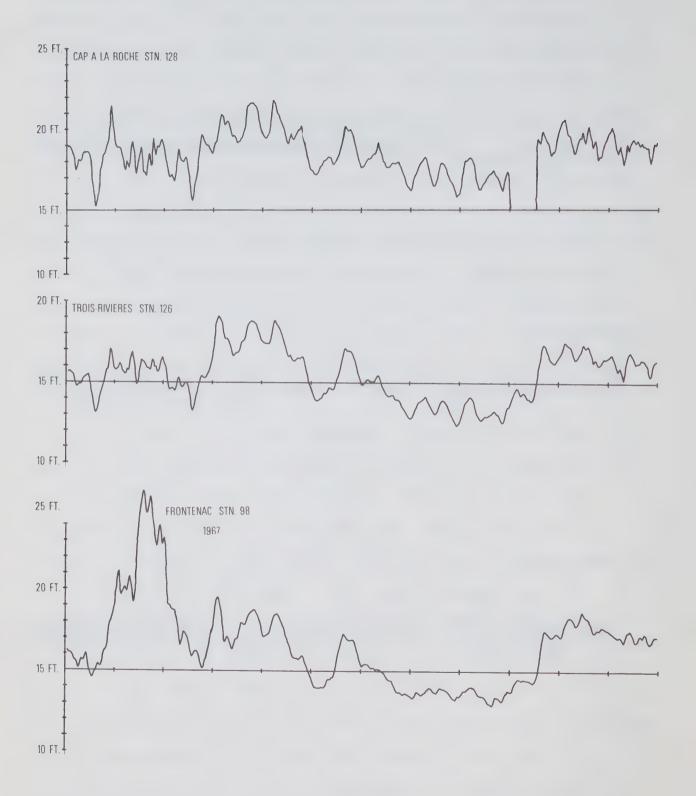


Fig. 16(a) (cont'd)



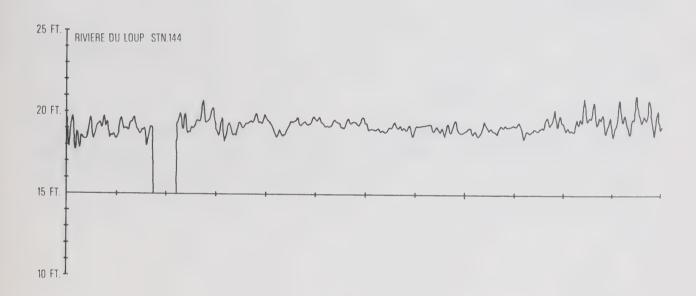




Fig. 16(b) Low Pass Filter of Hourly Heights, 1968.

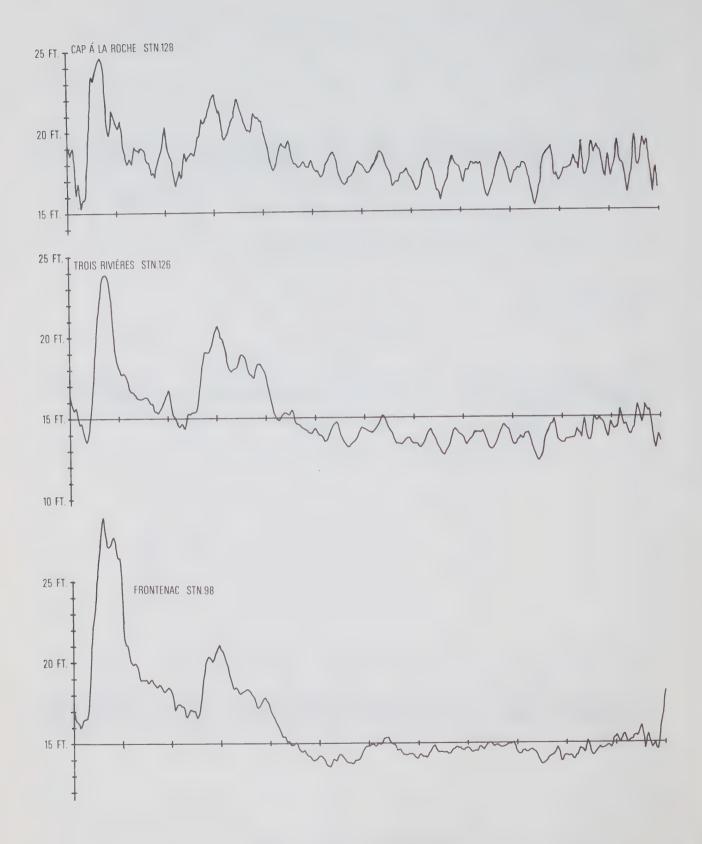


Fig. 16(b) (cont'd)

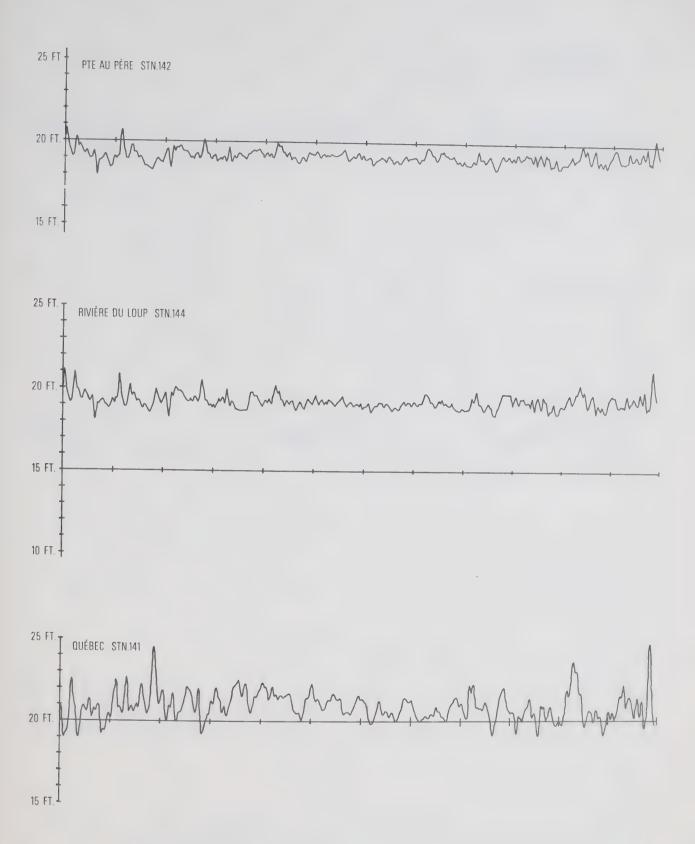


Fig. 16(c) Low Pass Filter of Hourly Heights, 1969.



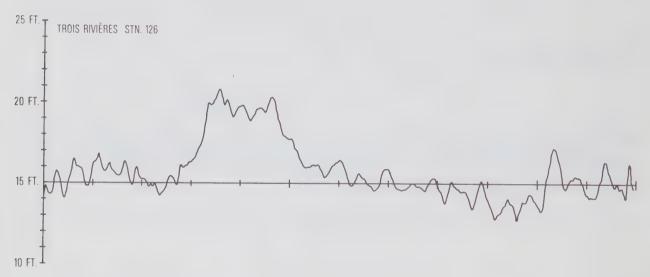




Fig. 16(c) (cont'd)

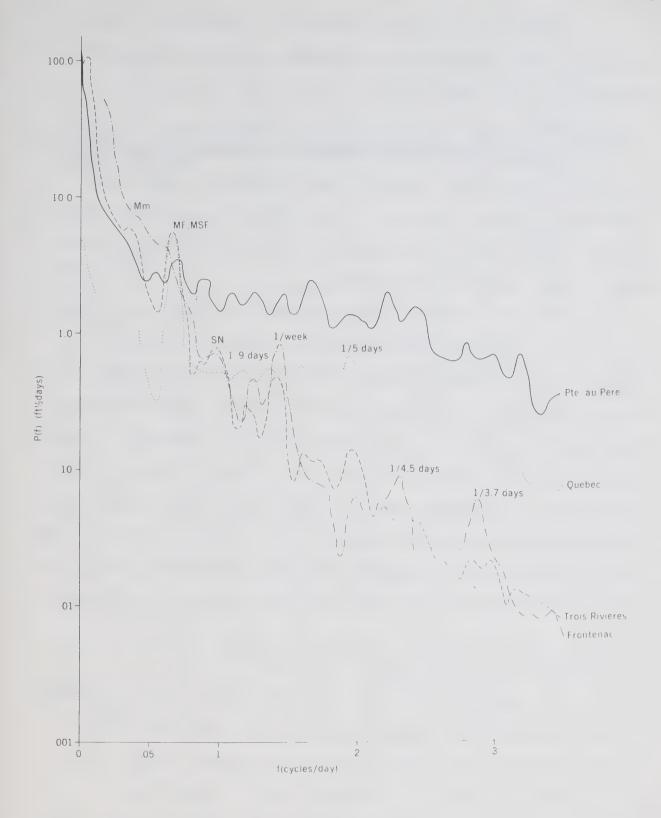


Fig. 17 Power Spectra of the Water Levels at Pointe au Père, Québec, Trois Rivières, and Frontenac. Low Frequency Band.

These oscillations, although of theoretical interest, are quite small compared to the diurnal and semidiurnal tidal oscillations and compared even to the random low frequency oscillations.

The bulk of the variations in the water levels is due to the semidiurnal and diurnal tides and Table 4 lists the amplitude in meters and phase in degrees of M2, S2, N2, K1, O1, MSf, Mm and Mf for the stations along the river for which one year of observations is available.

In Figs. 18(a) to 18(h) we show the same constituents plotted on a distance scale along the river: we give the amplitude and the phase change. Over the points of phase change, we enter the actual Greenwich phase lag of the constituent at the given locality. Only for Mm and Mf is it impossible to give any value for the phase change as there is much scatter in the observations. We notice that the amplitude of the constituents increases upstream to a maximum and then decreases rapidly. The decrease takes place past Québec for the semidiurnal and diurnal constituents and past Cap à la Roche for the slow constituents.

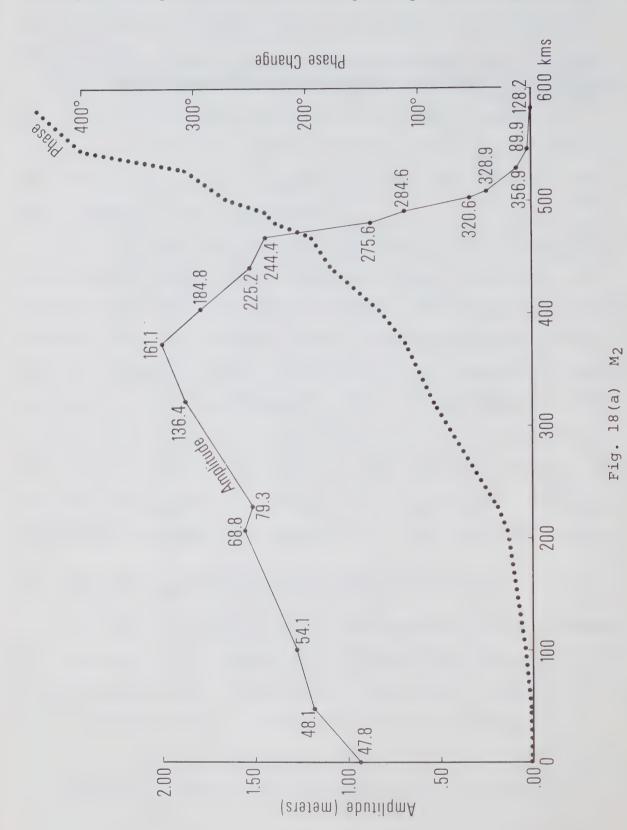
## 3.3 The analysis of the current observations at Québec Bridge and Isle aux Couldres

We have described in a previous publication our method for analyzing current observations (Godin: The Analysis of Current Observations, International Hydrographic Review, Vol. 44, No. 1, 149-164, 1967). It consists first in smoothing the observations if they have been taken at intervals smaller than one hour and then of low passing them to have a look at the

Table 4 Tidal Constituents of the Water Levels at Various Stations Along the St. Lawrence River Deduced from a One Year Analysis

	Y <sub>2</sub>	S <sub>2</sub>	$v_2$	<sup>(</sup> 1	ol	MSf	Уm	Уf
CHAMTON	1 7	V C	V C	Λ G	1 3	۸ G	, G	1 9
STATION  Ste. Anne Des Monts	.94 1,7.8	.29 815	.20 27.8	.21 204.7	.21 183.8	.03 102.6	.02 2711	.01 111.1
Raie Comeau	1.19 48.1	.38 810	.26 25.7	.23 200.9	.21 181.0	.02 94.3	.02 246.1	.01 216.8
Pointe-au-Père	1.29 54.1	.41 90.4	.27 29.9	.23 203.6	.22 181.0	.01 310.0	.02 277.4	.02 9.1
Tadoussac	1.57 68.8	.50 105.3	.33 16.4	.21. 207.1	.23 188.1	.0lı 273.9	.03 322.2	.nl. 217.5
Rivière Du Loup	1.52 77.3	.50 114.1	.32 54.5	.24 211.2	.22 190.8	.01 190.8	.02 292.2	.02 209.9
Jt. Jean Port Joli	1.88 136.4	.51 175.P	.34 111.2	.26 240.0	.24 219.5	.78 6.1	.ol. 31:0.8	.02 177.0
St François	2.02 151.1	.49 203.8	.36 137.1	.25 254.3	.24 234.0	.11 36.7	.06 57.5	.05 4.1.7
Québec	1.80 184.0	.42 208.1	.30 1/9.0	.23 270.5	.21 247.1.	.11 48.9	٠١٠ ١٠.3	.77 37.)
'euville	1.53 225.2	.31 274.3	.25 1713	.19 293.1	.10 260.3	.17 54.8	.10 34.3	. )^. ^ ).0
Portneuf	1.44 244.4	.30 289.2	.22 217.1	.18 306.5	.17 287.7	.20 52.2	.12 53.4	.07 41.1
Grondines	.87 275.6	.18 317.8	. U <sub>1</sub> 21 <sub>1</sub> 8.3	.11, 328.7	.13 303.7	.23 54.7	.v. 45.4	. 5 50.2
Cap à la Roche	.70 281,.6	.15 325.2	.11 253.5	.12 325.7	.12 310.3	.20 57.0	.13 52.1	.)8 1,7.2
Butiscan	.33 320.6	.08 355.6	.06 285.h	.08 394.6	.08 331.5	.18 62.7	.12 60.9	.05 50.6
Champlain	.24 328.9	.06 2.8	.05 277.0	.06 3.0	.06 338.1	.16 62.1	.11 52.9	.05 58.7
Trois Rivières	.08 356.9	.02 28.0	.02 318.7	.03 27.6	.03 7.1	.14 75.4	.10 /5.3	.05 48.3
Louiseville	.02 89.9	.01 119.6	.01 13.9	.02 88.9	.02 55.1	.12, 88.8	.09 72.5	.04 39.1
Sorel	.01 128.2	.00 158.5	.00 81.0	.01 129.9	.01 89.3	.08 92.0	.07 40.8	.03 72.1
Lanoraie	.01 157.2	.00 225.9	.00 101.3	.01 129.5	.01 71,.2	.17 326.0		
Lavaltrie	.01 211.7	. m 261:.2	.00 118.0	.01 153.6	.00 140.c	.10 326.3		
/ercher is	.01 215.0	.00 257.9	.00 13/.8	.01 1"1.1	.00 25).1	.10 322.3		
Varennes	.01 209.8	.01 286.1	.00 112.3	.01 121.7	.00 31/2.1	.09 310.3		
Longue Pointe	.00 229.7	.00 301.4	.00 01.2	.00 113.2	.00 243.9	.08 310.3		
Frontenac Street	.00 225.2	.00 292.6	.J) 106.2	.1) 131.6	.00 200.0	.08 304.8		
King Edward Pier	1 237.9	.:: 289.0	. 10 91.7	.00 93.1	. 0 2 3.2	.08 304.5		

Fig. 18 Amplitude and Phase Lag Along the St. Lawrence



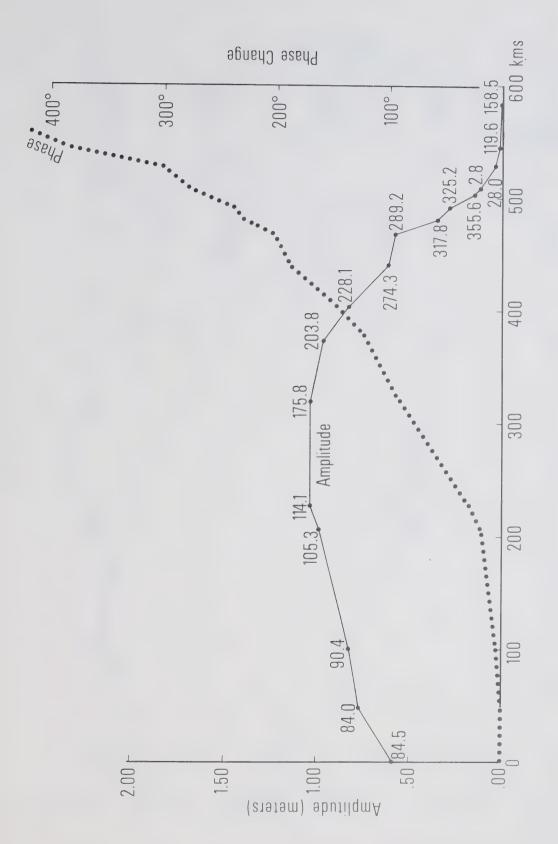


Fig. 18(b) S2

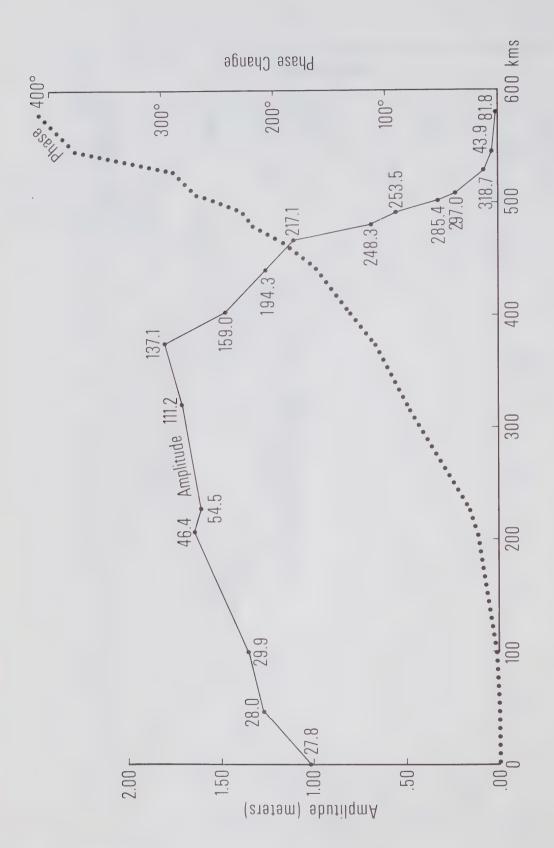
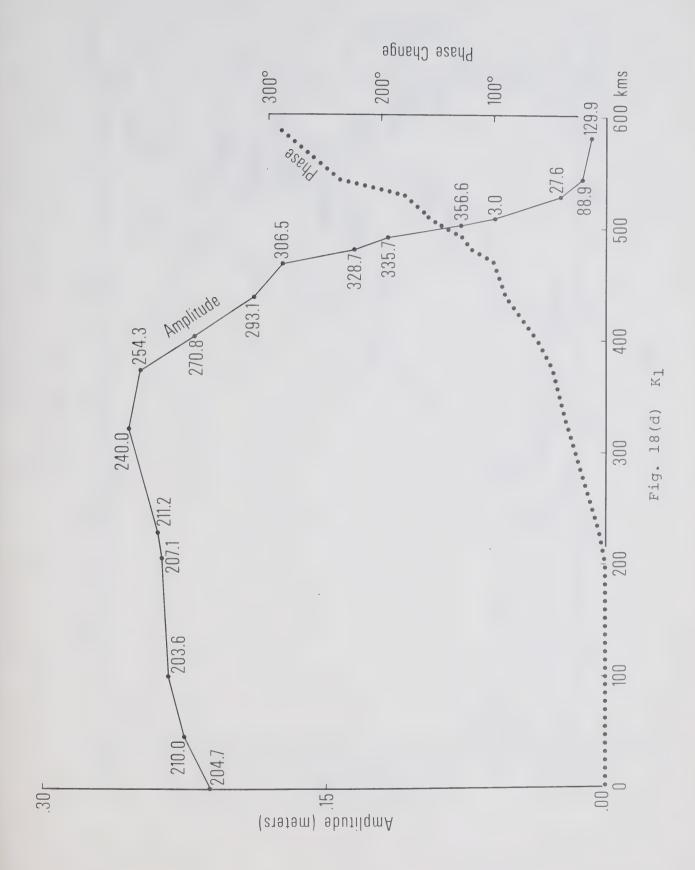
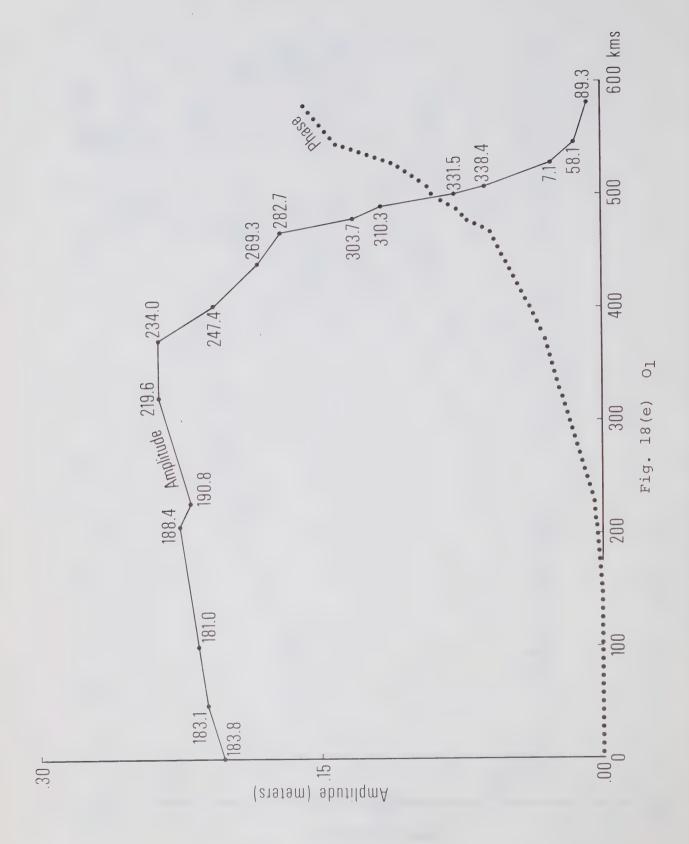


Fig. 18(c) N2





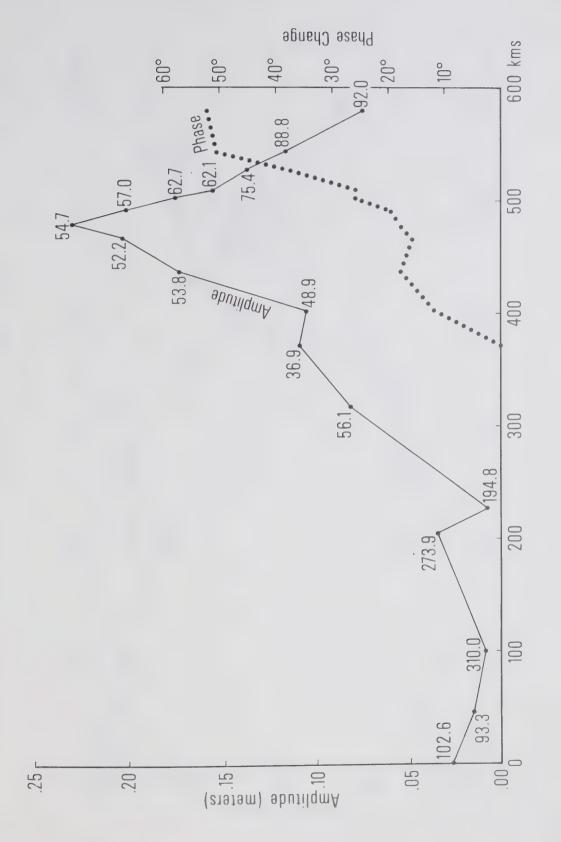


Fig. 18(f) MSf

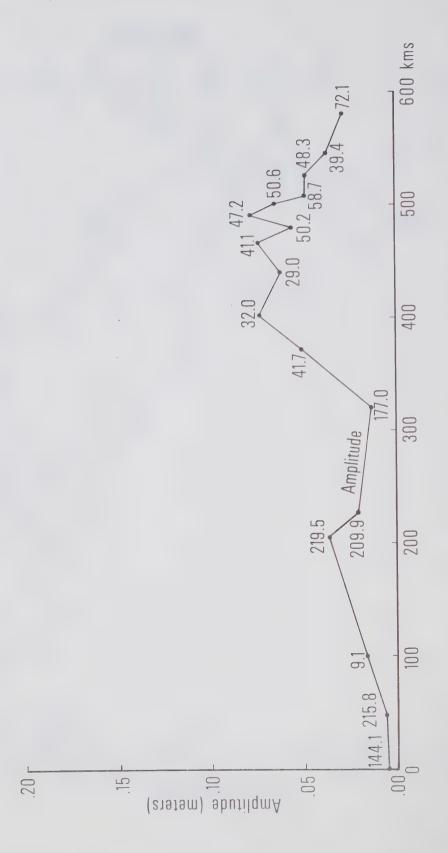


Fig. 18(g) Mf

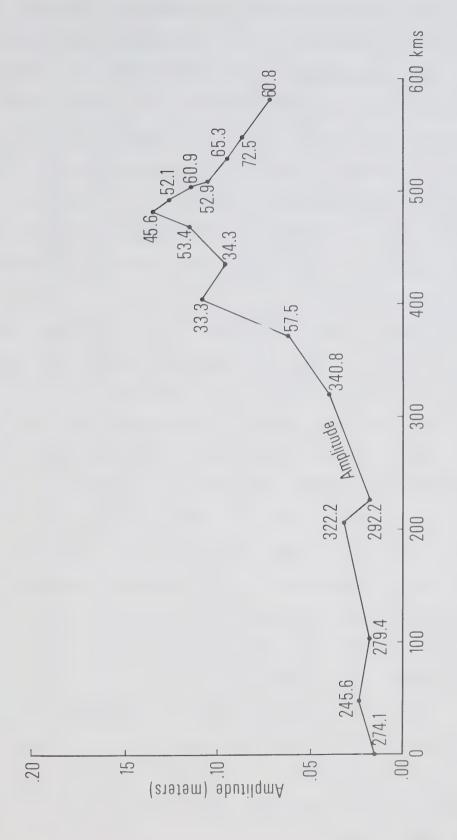


Fig. 18(h) Mm

fluctuations of the residual current, a part of the spectrum which is not worth subjecting to a harmonic analysis. Next the east and north components of the semidiurnal and diurnal bands are analyzed harmonically. This method takes in any amount of data and can handle as well discontinuous records which are more the rule than the exception. All the steps of the analysis are plotted in order to afford a visual interpretation of the success or failure of the analysis; a side profit is the quick discovery of the errors in the original data which could easily have escaped a thorough inspection of the numerical records and which can then be corrected.

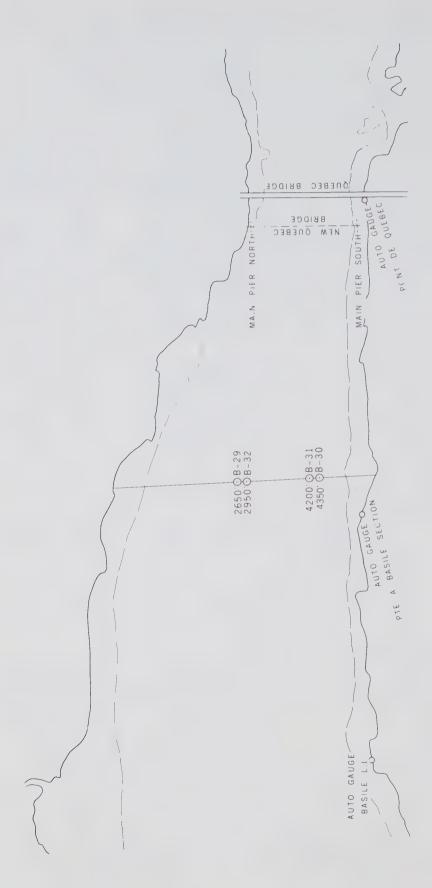
the Ship Channel Division at Québec Bridge and Isle aux Couldres during their survey of 1968. The table gives the code name of the station, the time interval at which the data were abstracted and the useful portion of the record which can be subjected to an analysis. Fig. 19 and 20 show the geographical location of the stations listed in Table 5. Stations B-29 and B-32 are located at nearly the same point of the river and their combined records may be analyzed as a single discontinuous record. The same thing holds for stations B-30 and B-31. The pair B-29 and B-32 is located exactly in the middle of the river while B-30 and B-31 are closer to the south shore.

For the Isle aux Couldres, station B-25 is located in the north arm of the river half way between St. Joseph de la Rive and Isle aux Couldres. The other stations are strung at approximately equal distances apart from the Isle aux Couldres to the south shore.

Table 5

Current Observations at Québec Bridge and Isle aux Couldres

Useable Portion of the Record			10.30 Oct. 2-15.30 Oct. 16 1968 12.30 Oct. 17-8.30 Nov. 3 1968	12.00 Oct. 2-13.30 Oct. 16 1968 12.00 Oct. 17-7.00 Nov. 5 1968		21.00 June 10-4.00 June 27 1968	17.30 June 18-17.10 July 10 1968 10.30 June 30-18.40 July 11 1968 0.00 June 21-0.30 July 4 1968 17.40 June 17-12.50 June 29 1968
Time interval between Observations	∆t minutes		0 0 8 8	30		30	10 10 10
Station Label		Québec Bridge	Middle B-29 B-32	South B-30 B-31	Isle aux Couldres	North B-25	South N-20 N-19 B-26 N-18



Location of the Current Meters and the Tide Gauges at Québec Bridge 19 Fig.

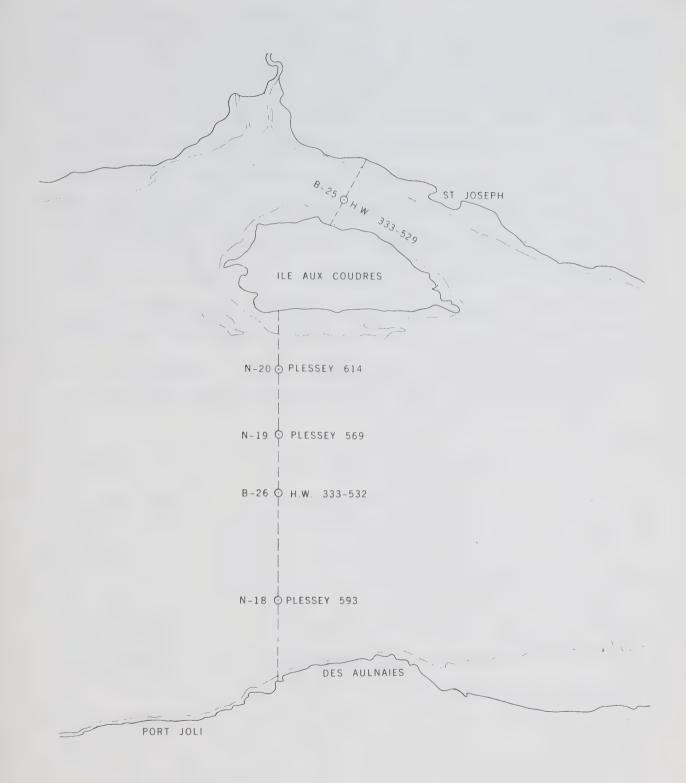
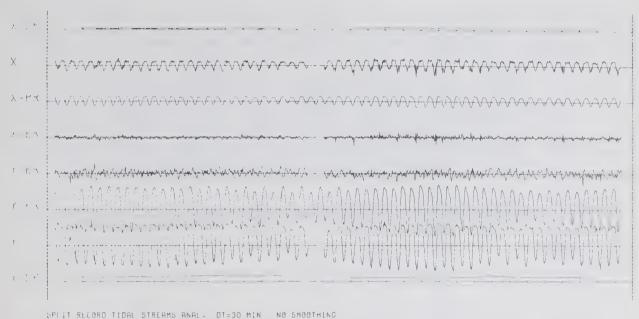


Fig. 20 Location of Current Meters at Isle aux Coudres

Figs. 21 and 22 show graphically the results for the analysis at Québec Bridge. X and Y represent respectively the east and north component of the current u in the raw. LP is the low pass and gives the residual current which here coincides with the steady current created by the river discharge. In the case under consideration the flow is remarkably constant. X-PR and Y-PR are the values of X and Y reconstituted from the analysis which is made up wholly of tidal harmonics. X-ER and Y-ER are the discrepancies between the observed current components and reconstituted components. The "error" graph shows more than random fluctuations: there are systematic high frequency oscillations which the harmonic analysis could not pin down. These are responsible for the failure of the reconstituted curve to reproduce the sharpness of the observed curve especially at neaps.

Fig. 23 shows the results for stations B-25, N-20, N-19, B-26 and N-18. The plots indicate that the analysis succeeds very poorly in reproducing the tide on account of the fact that the constituents N<sub>2</sub> and S<sub>2</sub> could not be separated from M<sub>2</sub> at most stations because of the short interval of observation. The only recourse left is to correct the analyzed values for these unknown constituents with the technique of inference, an unpleasant chore the results of which I have little faith in. For this purpose I have used the results of the analysis at the coastal station of St. Jean Port Joli which gives the following empirical relations within the subgroups of constituents:



RECORD 110H STREEMS HAND 11530 H.N NO SHOOTHING RECORD B-29 10-30HRS BCT 2768 TO 15-30HRS BCT 16/68 RECORD B-32 12-30HRS BCT 17/68 TO 8-30HRS NOV 3/68

Fig. 21 Analysis of Stations B-29, B-32 at Québec Bridge

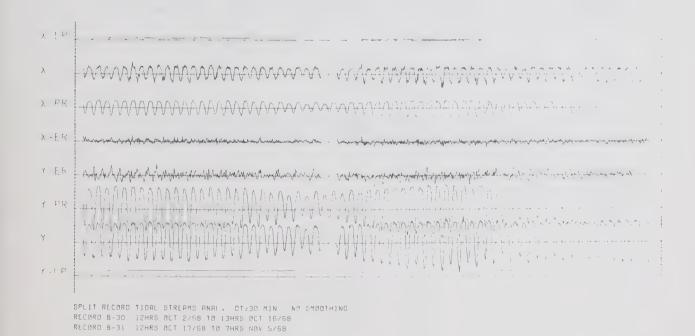


Fig. 22 Analysis of Stations B-30, B-31 at Québec Bridge

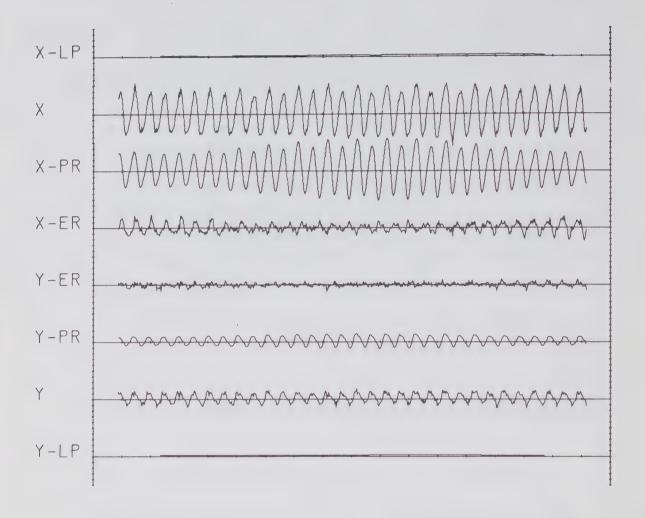


FIG. 23(a) #1. TIDAL STREAMS ANALYSIS DT=30 MIN NO SMOOTHING RC=.9 RECORD B-25 21HRS JUNE 10/68 TO 4HRS JUNE 27/68.

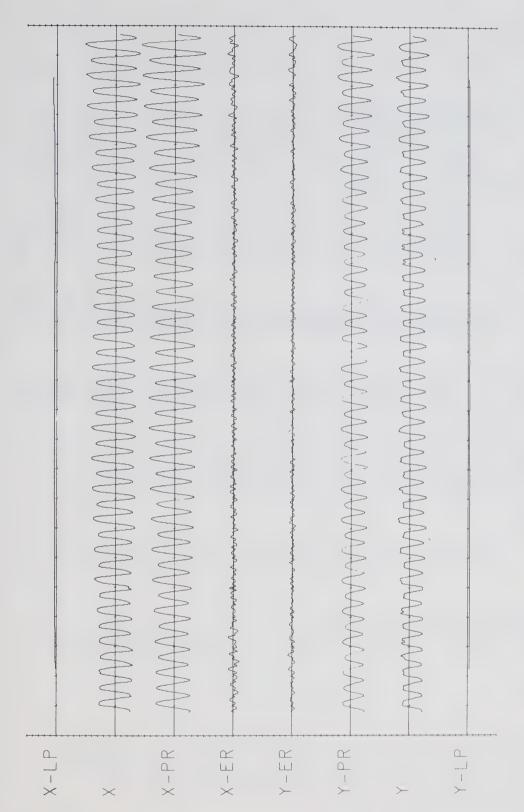


FIG. 23(a)#2. TIDAL STREAMS ANALYSIS DT=10 MIN SMOOTHING DT(ANAL)=30 MIN RC=0.8 RECORD N-20 17.40 JUNE 18/68 TO 18.50 JULY 11/68.

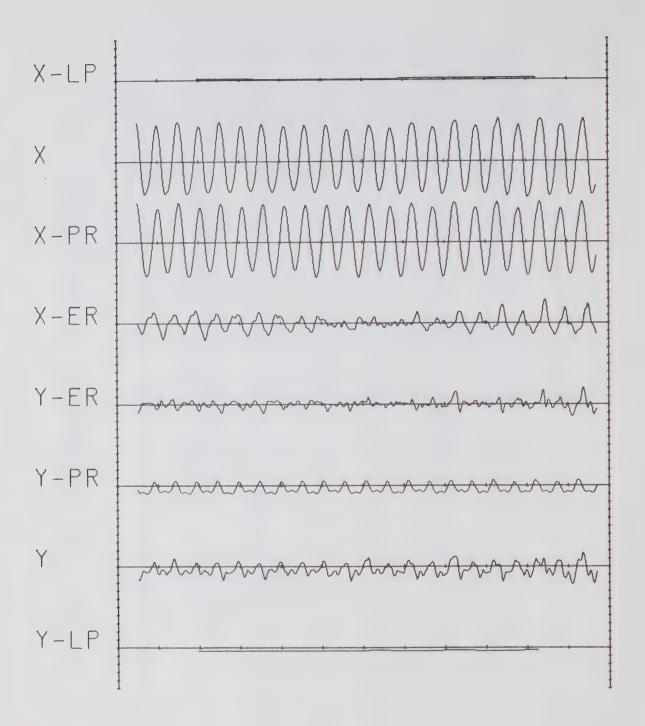


FIG. 23(a)#3. TIDAL STREAMS ANALYSIS DT=10 MIN SMOOTHING DT(ANAL)=30 MIN RC=0.9 RECORD N-19 10.30 JUNE 30/68 TO 18.40 JULY 11/68.

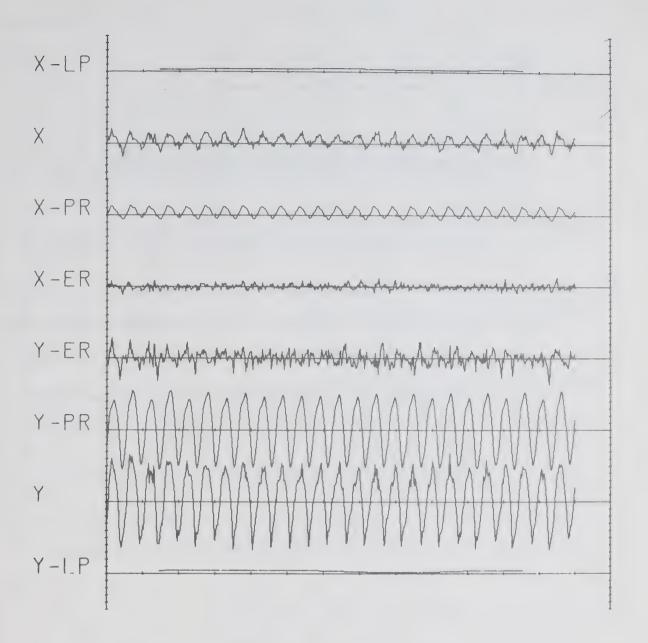


FIG. 23(a)#4. TIDAL STREAMS ANALYSIS DT=30 MIN NO SMOOTHING RC=.9 RECORD B-26 OHRS JUNE 21/68 TO OHRS JULY 4/68.

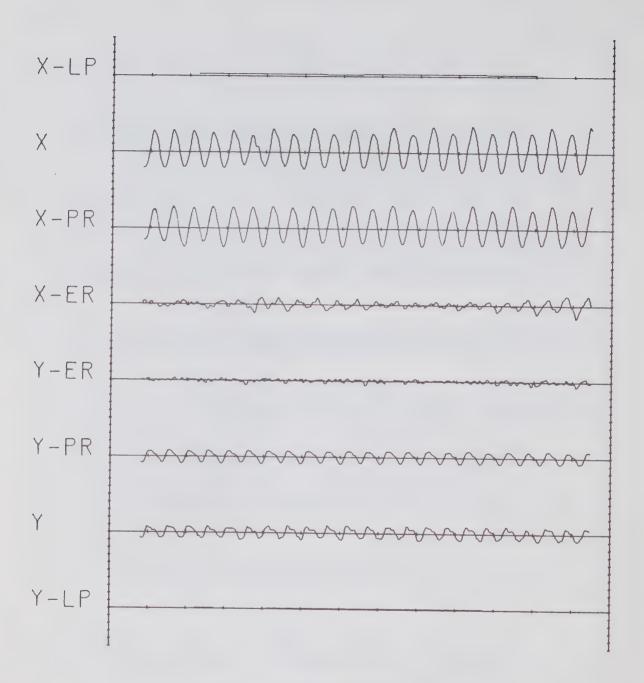


FIG. 23(a) #5. TIDAL STREAMS ANALYSIS DT=10 MIN SMOOTHING DT(ANAL)=30 MIN RC=0.9 RECORD N-18 17.40 JUNE 17/68 TO 12.50 JUNE 29/68.

## Constituents

M <sub>2</sub>	Name	Ratio	Phase Difference
	M <sub>2</sub>	1.00	00
	S <sub>2</sub>	. 27	+40°
	N <sub>2</sub>	.18	-25°
Кl	Кı	1.00	00
	o <sub>1</sub>	.93	-20°
	P <sub>1</sub>	. 27	0.0

In only one instance, station N-20, was it possible to lower the Rayleigh criterion to .8 in order to separate directly the important constituents. Fig. 23(b) shows the result of the inference of the constituents for the observations at Isle aux Couldres and the reduction of the Rayleigh criterion at N-20. Table 6 lists the results of the analysis for the major constituents. From the table, we see that the tidal streams reach their maximum intensity first on the shallow portions of the river and then in the channel; the tide is at its latest in the north arm of the channel, lagging by nearly two hours. The results on the steady current in the main channel seem absurd at first sight: at N-20 and N-19 the current flows almost at a right angle to the channel. A little bit of reflexion however, indicates that there may be an eddy motion set up by the tide and the steady currents reflect this general situation. More prolonged observations might give entirely different steady currents on account of the stronger averaging, while photogrammetry would indicate the instantaneous eddies set up during a tidal cycle.

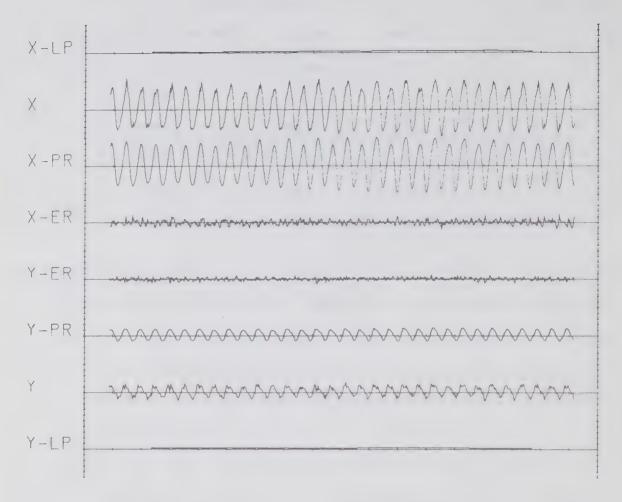
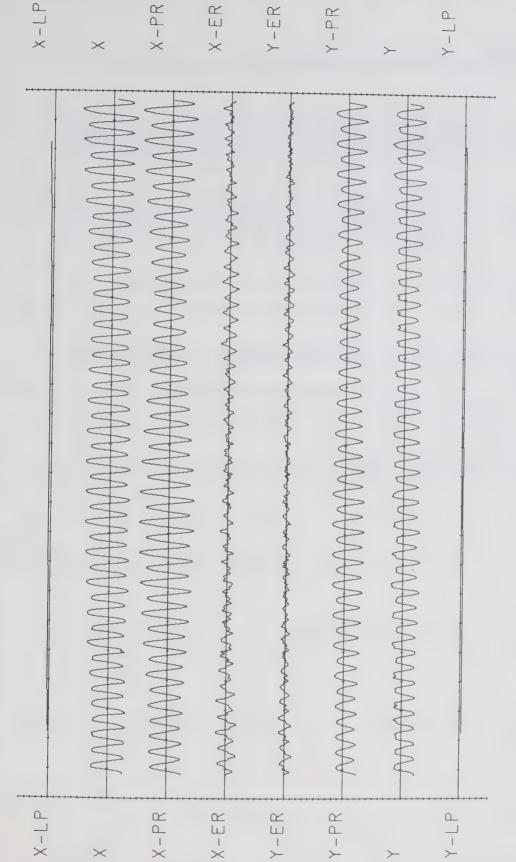


FIG. 23(b) #1. TIDAL STREAMS ANALYSIS DT=30 MIN RECORD B-25 21HRS JUNE 10/68 TO 4HRS JUNE 27/68.



DT=10 MIN SMOOTHING DT(ANAL)=30 MIN RC=0.9 RECORD N-20 17.40 JUNE 18/68 18.50 JULY 11/68. FIG. 23(b) #2. TIDAL STREAMS ANALYSIS

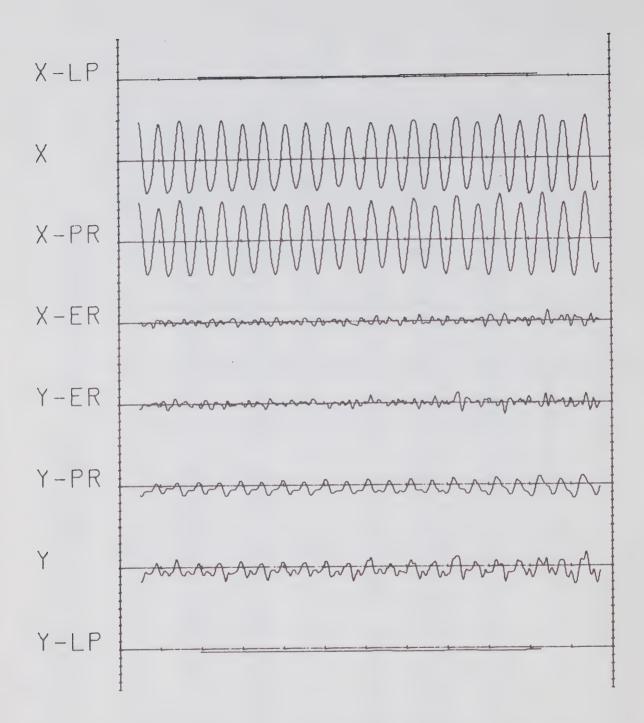


FIG. 23(b)#3. TIDAL STREAMS ANALYSIS DT=30 MIN RECORD N-19 10.30 JUNE 30/68 TO 18.40 JULY 11/68.

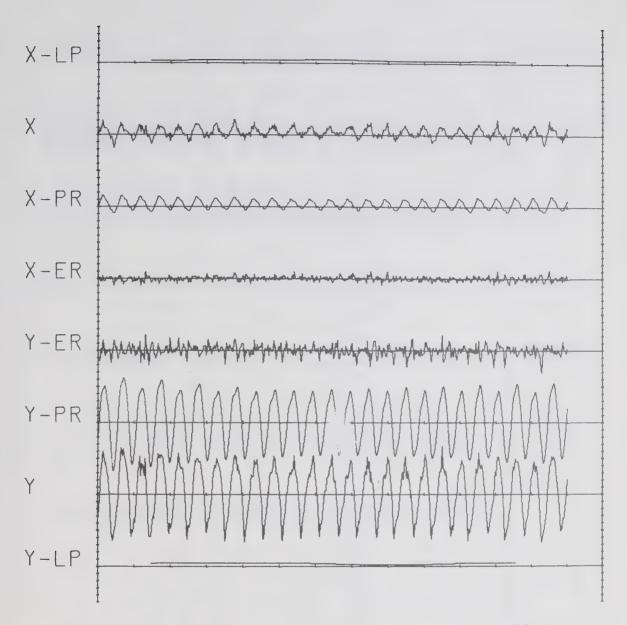


FIG. 23(b) #4. TIDAL STREAMS ANALYSIS DT=30 MIN RECORD B-26 OHRS JUNE 21/68 to OHRS JULY 4/68.

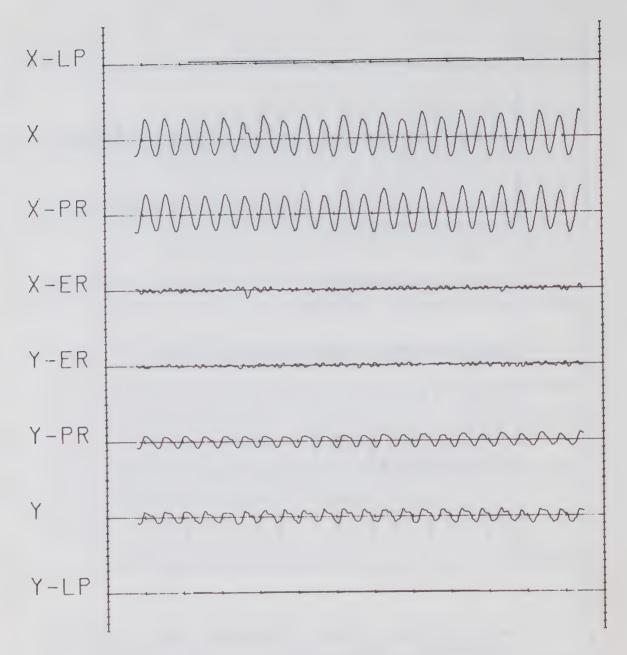


FIG. 23(b) #5. TIDAL STREAMS ANALYSIS DT=30 MIN RECORD N-18 17.40 JUNE 17/68 TO 12.50 JUNE 29/68.

the cur	Table 6	f the current observations taken by the Ship Channel Division at
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	Inc	86			75 304	
	p	349		184	290	248
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	×	. 22		.04	000	.05
	Inc	87		342	62 302 93 104	
	p	318 306		314	235 147 97 173	140
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	Inc	79		14	28 110 85 120	
_	מ	196		268	126 168@ 223 158@	163
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	×	13		0.0	.09 .08 .07	.03
	Inc	73		10	29 111 72 12	1
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	×	122.		.12	.07	.07
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	×	.40		.23@	.270 .250 .170	.22
	Inc	78		12	30 30 81@ 16@	
	deg	1843		147	83 93@ 107@ 86@	94
Souldres S2	sec sec	1.1		1	01 06@ 02@	
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and	deg	153		8 6	533	051
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ec E	Sela B	·				9
ŏner ŏ	sela M	1.67		1.27	1.30 1.51 1.39	1.26
	Inc.²	74		43	323 294 48	451
	m/sec	e	ldres	.10	41. 71. 70.	90.
	Station	Québec Bridge B-29, B-32 B-30, B-31	Isle aux Couldres	North Shore B-25	Main Channel N-20 N-19 B-26 N-18	Average

M=semimajor axis  $^{*}$  -solutive, indicates that the current vector rotates counterclockwise around the ellipse; otherwise, it is clockwise

Vector average of the component along the river bed which we assume orientated at 45° . The inclination is in trigonometric degrees, not in geographical degrees is the first of the for \( \frac{1}{2} \).

3=uncorrected for K2

The minor component of the tidal stream ellipse is very small; an indication that the tidal streams are very nearly rectilinear everywhere.

3.4 Comparison of the values of the tidal streams observed with those predicted by W.D. Forrester using the Equation of Continuity

W.D. Forrester (B.I. 69-2 and B.I. 67-5) has used the equation of continuity

$$Q_{X} = -BD_{t}$$
 (6)

in order to estimate the tidal streams in the St. Lawrence River from the observations derived from the tide gauges strung along the river, with the assumption that the tidal oscillation vanished around Lake St. Peter.

For a given harmonic, (6) becomes

$$Q_{X} = -i\sigma BD$$
 (21)

where  $\sigma$  is the frequency of the harmonic. In finite difference form (21) becomes

$$\frac{\Delta Q}{\Delta +} = -i\sigma BD \tag{22}$$

If the river is divided into a succession of channels of length  $\Delta x_j$ , width Bj, j labelling each channel, (22) becomes for a given channel:

$$Q_{j+1}-Q_{j}=-i\sigma B_{j}D_{j}$$
  $\Delta x_{j}=-i\sigma A_{j}D_{j}$  (23)

Aj being the horizontal area of the jth channel. (23) along with the boundary condition:

$$Q_{O}=0 \tag{24}$$

allows the value of the flow  $Q_j$  at section j to be evaluated for the vertical displacements upstream:

$$Q_{j+1}=-i\sigma \sum_{k=0}^{j} A_k D_k$$
 (25)

For a harmonic constituent with amplitude  $|D_k| = d_k$  and phase lag  $\delta_k$ , (25) may be written as

$$Q_{j+1} = \sigma \sum_{k=0}^{j} A_k d_k e^{-\frac{1}{2}\pi}$$
 (26)

from which the magnitude and phase of the flow Q at the entrance of the j+lth channel may be computed. We refer to the publication "Tidal Transports and Streams in the St. Lawrence River and Estuary" B.I. 69-2 by Dr. Forrester for specific values of Q and u for the various tidal constituents along the river.

In Table 7 we entered the value observed at Pointe des Monts by Farquharson (BIO 66-6), at Pointe au Père by Forrester

Table 7

Value of the Harmonic Constituents of the Currents Observed Compared with the Value Calculated by Forrester Using the Equation of Continuity

Steady Current m/sec	ent m/	M2 sec o	m	S 2 s e c	/ш 0	N2 sec o	K m/s	0 ec 0	m/sé	0 G 0	M/se	0 0 0	MS4 m/sec	0
N.														
	91.	339	90.	015	.04	316	.018	118	.016	101	.005	353	.001	00
	.12	356	.03	031	.03	330	.014	124	.012	107	.005	354 021	.001	05
aux Couldres														
	1.16	058	.32	100	.21	033	.08	157	.07	141	.00	159	.04	21
	1.55	154	 	205	30	117	.12	224	11.	209	.40	312	.22	34
rre	*S <sub>2</sub> has been corrected	for the	in	fluen	ice of	K2.	N2	has b	een	corre	cted	for	۰2٠	

40

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∞ 0

(BIO 67-5), Isle aux Couldres and Québec Bridge by the Ship Channel Division (1969) and we compare them with the values deduced by Forrester using (26) which are based directly on the equation of continuity.

The values predicted agree to within 10 percent of the value observed; it must also be realized that the "observed" values are uncertain by this amount as well on account of the relatively short duration of the observations and the highly variable character of the currents.

This brings to the fore the question of the wisdom of taking current observations in a river, an undertaking which costs hundreds of thousands of dollars, when quite adequate values of the mean flow may be deduced from the tide gauge network if the latter is dense enough. The fine structure of the current flow across a given vertical section cannot be predicted in this way. To find out about it, one needs to set up a network of current meters at various depths and locations across the river during one or two tidal cycles at neap and spring tide. The fine structure is of no use in a one dimensional study of the river and has to be integrated away in the first place. It can only be required on rare occasions to solve specialized problems of hydrodynamics or of hydraulics.

# 4. A mathematical model of the St. Lawrence River

If the equations (6) and (7) hold everywhere in the St. Lawrence there is good reason to believe that their integration over the whole river bed using the actual widths B and depths D would yield all the possible modes of motion.

There is much wishful thinking in this statement since (6) and (7) are only approximations to the more exact Equations of Motion and there is no a priori reason why the river should "feel" a width B and a depth D obtained by schematizing its channel into parallepipeds of rectangular cross section. There is also the added complication that the discharge is not wholly present at the head in Montréal but that it is gradually but irregularly augmented all along by the contribituon of the tributaries like the Richelieu, the St. Maurice, the Saguenay and the Manicouagan. There is the problem of the advection of salt water into the fresh water régime which forces a vertical distribution in the densities and the velocities, with which (6) and (7) cannot cope. Finally, as the river widens there are definite gradients in the elevations and currents across the width of the channel.

Still it is obvious that we do not want to face all the complications in a first approach study and that it is desirable to integrate (6) and (7) over a rectangular schematization of the channel, neglecting the gradual changes in the discharge. The results of such calculations should give us an overall view of the first order type of motion occurring. Comparison with observations will show that such calculations model adequately the main features of the tidal motion upstream and its interference with the discharge.

Also, we will then be in a position to obtain definite information on the influence of some of the basic parameters of the river on the propagation of the tide upstream. We will be

able to investigate the effect of an increase of the discharge on the tide moving upstream and to investigate the effect of closing the river at various crucial points like at Montréal, at the mouth of Lake St. Peter or at the constriction of the site of Québec Bridge. The conclusions which we will reach cannot be essentially modified by the increased complications in the river flow, which we have mentioned at the beginning of this paragraph.

#### 4.1 The schematization

The bed of the river has been subdivided into 53 subchannels each of which is schematized as a channel of rectangular section of width B and depth D. The width B is the average width of the section while D is the overall average of the depth over the section. The criterion for initiating or terminating a subchannel is the occurrence of a geographical accident, an island, a constriction or widening or a change of depth. Working in this way from King Edward Pier in Montréal to Pointe des Monts, we ended up with 53 sections of various lengths each of these sections being relatively homogeneous in depth and width, as shown in Figs. 24 and 25.

There exists no single map showing the whole of the river bed between Montréal and Pointe des Monts; it is therefore not practical to attempt to show our subdivisions. However, Table 8 lists, besides the section number, the name of the coastal stations which are located within a section; this may give to the reader an idea of the location of the subdivision.

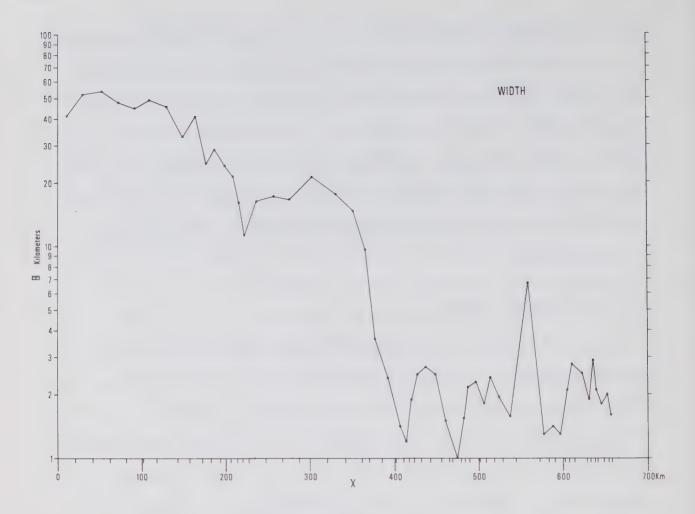


Fig. 24 Width B of the Schematized St. Lawrence River

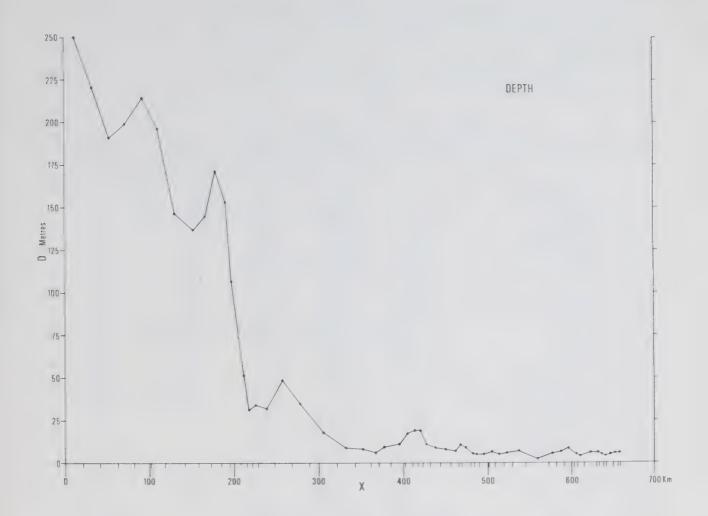


Fig. 25 Depth D of the Schematized St. Lawrence River

Table 8
Schematization of the St. Lawrence River

Section		Horizontal Area x106 meter2			D	Accumulated distance
				meters	meters	x10 <sup>3</sup> meters
0-1	Pointe des Monts Petite Matane	883	220353	41520	249.5	21.2
1-2	Matane	1077	237424	51788	220.3	42.1
2-3	Baie Comeau	1141	218171	53740	191.2	63.3
3-4	Manicouagan	859	169882	48054	197.7	81.2
4-5	Mont-Joli	914	195475	45166	213.8	101.4
5-6	Pointe au Pêre	896	175774	49314	196.2	119.6
6-7	Rimouski	997	146326	46024	146.8	141.3
7-8	Bic	558	76508	33131	137.2	158.1
8-9		600	86751	40693	144.7	172.8
9-10	Escoumain	289	49563	24272	171.6	184.7
10-11	Trois-Pistoles	308	47041	29059	152.8	195.3
11-12		282	30358	23723	107.6	207.2
12-13	Tadoussac	171	8828.0	21305	51.8	215.2
13-14		82.6	2657.0	15615	32.2	220.5
14-15	Cacouna	148	4987.5	11378	33.8	228.2
15-16	Rivière du Loup	413	13424	20030	32.5	248.8
16-17	•	310	14872	16455	48.0	267.6
17-18	Pointe au Pic	349	12375	17160	35.5	288.0
18-19	St. Joseph de la Rive	543	8969.8	16628	16.5	320.6
19-20	St. Jean Port Joli	474	4470.3	21438	9.4	342.7
20-21	Montmagny					
	Isle aux Couldres	337	2612.4	17633	7.8	361.9
21-22		154	1002.4	14565	6.5	372.4
22-23	St. François IO	117	872.10	9601	7.5	384.6
23-24	St. Laurent IO '	74.6	829.82	3641	11.1	405.1
24-25	Lauzon	13.8	241.91	2390	17.5	410.9
25-26	Wolfe Cove	9.29	175.74	1445	18.9	417.3
26-27	Pont de Québec	7.77	149.91	1194	19.3	423.8
27-28		15.7	191.81	1879	12.2	432.1
28-29	Neuville	30.1	280.33	2456	9.3	444.4
29-30		32.5	254.62	2741	7.8	456.3
30-31		23.3	168.77	2509	7.2	465.5
31-32	Portneuf	5.41	59.016	1492	10.9	468.2
32-33	Deschambault	8.98	82.147	803	9.1	480.4
33-34	Grondines	5.84	33.551	1555	5.7	484.1
34-35		13.7	64.069	2217	4.7	490.3
35-36	Cap à la Roche	25.6	123.34	2336	4.8	501.3
36-37	Batiscan	14.7	95.413	1841	6.5	509.2
37-38	Champlain	23.6	128.39	2439	5.4	518.9
38-39	Trois-Rivières	20.1	131.56	1943	6.5	529.3
39-40	Port St. François	25.1	205.45	1591	8.2	545.1
40-41	Louiseville	188	541.10	6794	2.9	572.8
41-42	- 3	11.7	71.639	1304	6.2	581.7
42-43	Sorel	14.1	104.33	1429	7.4	591.6
43-44 44-45	Lanoraie	14.6	123.94	1320	8.5	602.6
45-46		10.1	52.539	2078	5.2	607.5
45-40	Lavaltrie	0.2		2052	4 0	61.6
46-47	Contrecoeur	21.6	92.242	2852	4.3	615.1
47-48	Verchères	30.0	175.03	2047	5.8	629.7
48-49	Répentigny	6.16	35.131	1932	5.7	632.9
49-50	Varennes Pointe aux Trembles	11.6	55.736	2093	4.8	638.5
50-51	Longue Pointe	6.35	24.850	1757	3.9	642.1
51-52	Laurier	13.7 9.11	70.631	1992 1579	5.2	649.0
52-53	King Edward Pier	3.03	54.947	1015	6.0	654.7
32 33	King nawara rice	3.03	18.797	1013	6.2	657.7

Table 8 lists the various sections along with their relevant parameters such as the horizontal area, the volume, the average width, the average depth and the accumulated length of the schematized sections.

4.2 The rewriting of the Equations of Hydrodynamics into systems of finite differences

The equations (6) and (7)

$$Q_X + Bht = 0$$

$$-h_x-Q|Q|/C^2B^2D^3 = Q_t/gBD + 2QQ_x/gB^2D^2$$

have to be rewritten as

$$h_{+} = -Q_{X}/B \tag{27}$$

$$Q_{t} = -gBD (h_{x} + Q|Q|/C^{2}B^{2}D^{3}) + 2QQ_{x}/gB^{2}D^{2}$$
 (28)

in order to be able to go from past to future values of the variables h and Q.

The river may be cut up into subsections which we label as

At one point we compute h and the next one we compute Q. At an interior point (27) and (28) can be written out in terms of finite differences as

$$h(n,t+2\Delta t) = h(n,t) - (2\Delta t/B_n)(Q(n+1,t)-Q(n-1,t))/(\Delta x_{n-1}+\Delta x_n)$$
 (29)

 $Q(n+1,t+2\Delta t)=Q(n+1,t)$  -

$$2g\Delta tB_{n+1}D_{n+1}[(h(n+2,t)-h(n,t))/(\Delta x_{n+1}+\Delta x_{n+2}))$$

$$+Q(n+1,t)|Q(n+1,t)|/C^2B_{n+1}^2D_{n+1}^3)$$

$$+2Q(n+1,t) (Qn+3,t)-Q(n+1,t))/(\Delta x_{n+2}+\Delta x_{n+3})$$
 (30)

These relations hold for an inner point. The values of B and D are averages of the values they have in the two sections delimited by the three points.

At the mouth we impose an h boundary condition so that h(1,t) is specified at all times. At the head we impose a condition on Q. The differential for  $Q_X$  in the convective term is computed in a biased fashion but the term is quite small anyway. The boundary condition on h at the mouth is  $h(1,t)=1.03\cos(28.98^{\circ}t+89^{\circ})+.33\cos(30.00^{\circ}t+273^{\circ})+.23\cos(28.44^{\circ}t)$ 

 $+107^{\circ})+.22\cos(15.04^{\circ}t+38^{\circ})+.20\cos(13.94^{\circ}t+79^{\circ})$  meters The Q boundary condition at the head is  $Q(54,t)=0,8.5,11.3,14.2\times10^{3}m^{3}/sec$ 

The h boundary condition embodies the approximately observed values of M2, S2, N2, K1 and O1 at section 1 and the phase has been chosen to correspond to the instant 00 hours, September 1, 1964. t is measured in hours. At Montréal the discharge is taken to be a constant at all times.  $Q=8.5\times10^3/\text{sec}$  corresponds

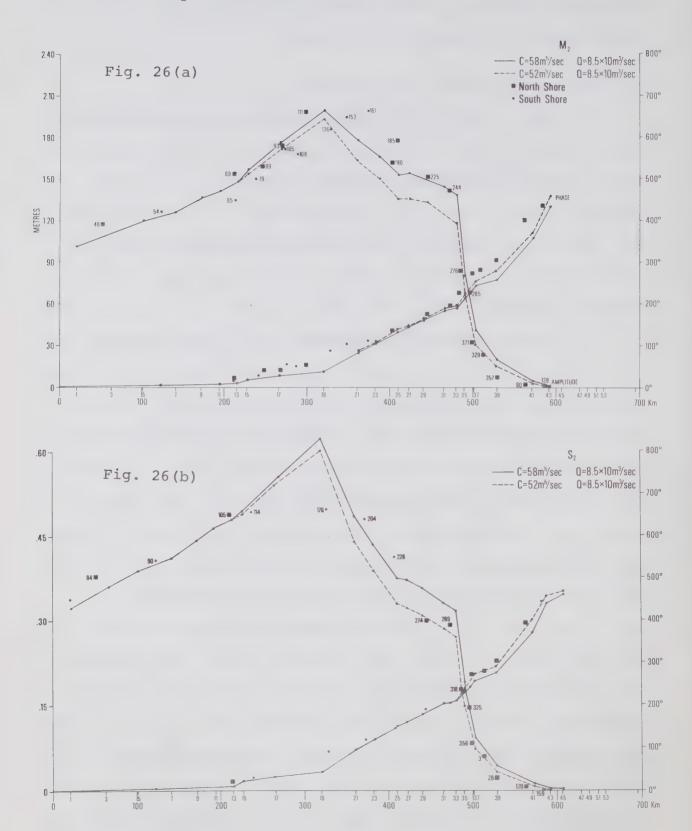
to a little bit more than the average discharge in Montréal. The higher values may prevail during the spring run off. Q=0 corresponds to a closure of the river.

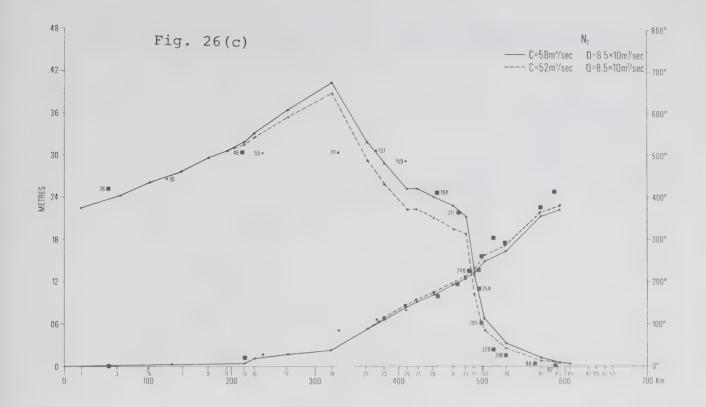
#### 4.3 The results of the calculations

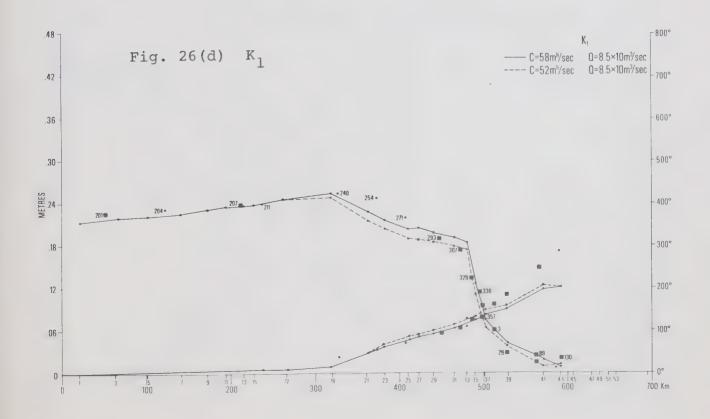
Plots 26(a) to 30 show the computed amplitudes and phase changes of some constituents compared to their observed values derived from the analysis of observations in the St. Lawrence River. Most of the calculations have been performed for C=58 m½/sec since our previous investigations in par. 2.6 indicated that such a value of the friction may hold in the lower St. Lawrence.

M2 being the most important constituent, we shall look at it first. Fig. 26(a) like all the others which follow, shows as continuous curves the results of the calculations using our schematized model of the river. One curve shows the variation of the amplitude up the river while the other shows the phase change upstream of the particular constituent. The analyzed values based on observations are shown by dots; those which are squared were derived from stations located on the north shore. The Greenwich phase lag of the constituents is written explicitely over the point and allows the location of the observed phase lag points over the phase curve. For M2, all the stations observed have been entered even if the observations covered less than one year. For the other constituents we have entered only those based on a one year analysis since the shorter observations bring in much scatter. Fig. 26(a) shows the calculated amplitude and phase change of M2 for C=58 m<sup>1</sup>2/sec

Fig. 26 Amplitude and Phase Change for Two Values of the Chézy Coefficients







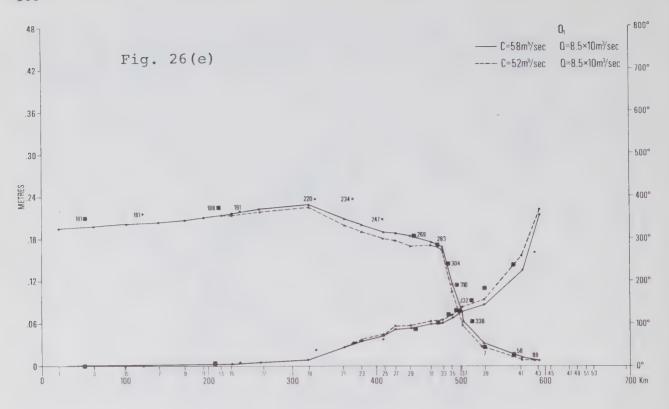


Fig. 27 The Fast Shallow Water Constituent  $M_4$  for Three Values of the Discharge

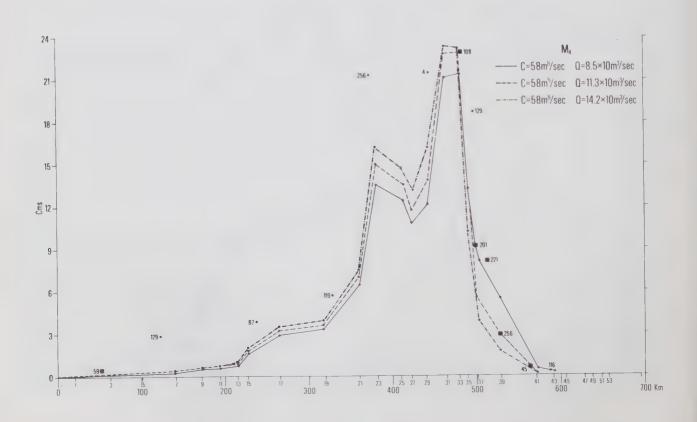


Fig. 28 The Slow Shallow Water Constituents MSf for Three Values of the Discharge and Two Values of the Chézy Coefficient

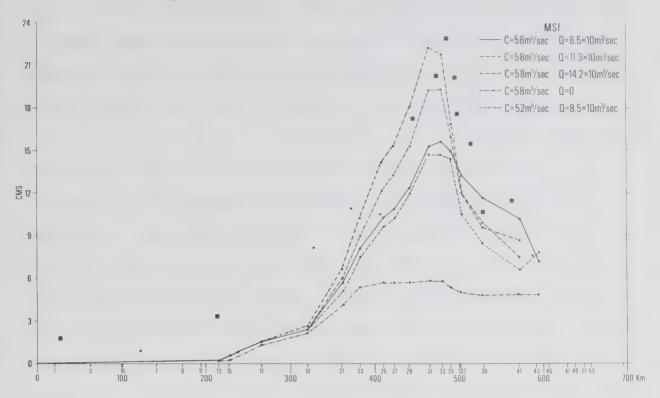
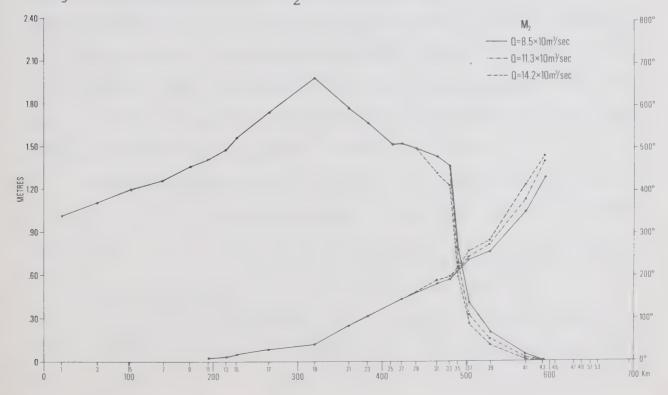


Fig. 29 The Constituent  $M_2$  for Three Values of the Discharge



at a head discharge of 8.5x10<sup>3</sup> m<sup>3</sup>/sec; it underlines the fundamental weaknesses of our model which we have discussed in the introductory paragraph. We had to leave out of our schematization the Saguenay River whose mouth stretches over sections 11 to 13 and we schematized the double channel around by the Ile d'Orléans by a single channel which stretches between sections 22 and 25.

At Saguenay, some of the tidal energy is diverted into the fiord and we notice that upstream of Saguenay our calculated phases of M2 fall consistently below those observed. Similarly our analyzed amplitude does not model the observed amplitude adequately over sections 21 to 24. Still the overall agreement between the values derived from our model and those observed is adequate and suffices for a study of the fundamental parameters which influence the propagation of the tide upstream.

It is obvious that we get better agreement for  $C=58m^{\frac{1}{2}}/\text{sec}$  with the observed amplitudes, while the difference in phase between C=52 and  $C=58m^{\frac{1}{2}}/\text{sec}$  is slight. The value of  $C=58m^{\frac{1}{2}}/\text{sec}$  corresponds to the value of C we deduced from the observations of the Ship Channel Division.

The subsequent plots of  $S_2$ ,  $N_2$ ,  $K_1$  and  $O_1$  show as well the overall adequacy of our model for describing the general features of the tidal motion upstream.

#### 4.3.1 The shallow water constituents

One of the attractive features of our model is that it creates as well, shallow water constituents through the non-linear friction term  $Q|Q|/C^2B^2D^3$ . Out of the interaction

of M2, S2, N2, K1 and O1 we should get new constituents such as M4, MS4, S4, 2MS6, M8, MSf, SN, Mm and Mf amongst others. Our model does yield some values for these quantities which may be compared with observations. In Figs. 27 and 28 we show plots of M4 and MSf. The three curves in the diagram for M4 show the calculated M4 for various values of the discharge; M4 increases first for larger values of the discharge up to Québec, then the relation is reversed. On the way to Montréal it falls off sharply as our rudimentary algebraical calculations had indicated. Our calculated M4 is significantly smaller than the observed values; there is a discrepancy of up to 12 cm in the Isle d'Orléans. MSf has been calculated for three values of the discharge and two of friction, which we know are possible in the river. For Q=8.5x103m3/sec which corresponds to the mean discharge, our calculated MSf falls quite short of the observed values; increasing the discharge does enhance MSf very appreciably though. A further indication that MSf is fundamentally dependent on the value of the discharge is that for Q=0, MSf is very slight throughout. The calculated values for the other shallow constituents, fast and slow, are similar to those we have shown for M4 and MSf.

We conclude that our model provides a qualitative agreement with the observed values of the shallow water constituents but certainly not a quantitative agreement. It is my personal feeling that such a quantitative agreement would be very difficult to obtain even with a very sophisticated model as these constituents are created through a very subtle and sensitive interaction of the fundamental parameters of the river.

# 4.3.2 The influence of the discharge on the propagation of the tide

An inspection of Figs. 27, 28 and 29 shows the influence of the discharge on the propagation of the tide upstream. In Fig. 29 we see that the fundamental constituent M2 decreases in amplitude and is slowed down in its progress upstream for increasing values of the discharge. The effect of the change in discharge is not felt till beyond the constriction of Québec City. The modification in amplitude amounts to at most 10 to 20 percent of the tidal amplitude; the time of arrival of the M2 tide however may be retarded by up to two or three hours. The shallow water constituents created by the interaction of the tide with the discharge and bottom friction are significantly increased in the downstream portion of the river for increased values of the discharge but they become more rapidly extinguished upstream of Québec City, basically because the original tidal constituents are themselves more rapidly damped.

# 4.3.3 The influence of quadratic friction

Figs. 26(a) and 28 show the effect of the quadratic friction. In Fig. 26(a) we notice that for lesser values of the friction, the tide is damped less rapidly and it suffers less phase retardation, a fact which we could demonstrate algebraically in par. 2.2 in the case of linearized friction. Now we see that this holds as well in the case of quadratic friction.

The effect of quadratic friction on the shallow water constituents is not that apparent; the use of linearized friction does not even give us a clue as to their existence. Fig. 28 indicates that MSf is lessened for increased values of friction;

the same holds for the fast shallow water constituents such as M4 and MS4 although we have not shown this explicity in diagrams. The use of our intuition would lead us to believe that the stronger the friction the sharper the shallow water constituents. The facts show that this is not the case; an indication once again that non-linearities escape the grasp of the human mind. I believe that if in our case the non-linear friction were to be varied slowly from zero, the shallow water constituents would increase initially but beyond an optimum value they would start diminishing as the primary wave starts wasting itself into more complicated additional modes of motion. Most likely between the values of C=58 m½/sec and C=52 m½/sec, we must have passed beyond this optimum value of C.

## 4.3.4 The effect of the closure of the river

Our model allows us to investigate the effect of closing the river wherever we please; it seems interesting to see what would happen to the propagation of the tide if we closed the river at the constriction where the Québec Bridge has been laid, or at the mouth of Lake St. Peter or at Montréal itself.

Even before attempting any calculations, a little bit of reasoning indicates that the tide originates from the Atlantic Ocean and that it becomes progressively wasted in the Saguenay River and up the St. Lawrence. The erection of any barrier would allow more tidal energy to stay downstream. Our calculations support this intuitive conclusion; they also delineate the portion of the river where the dyke would modify

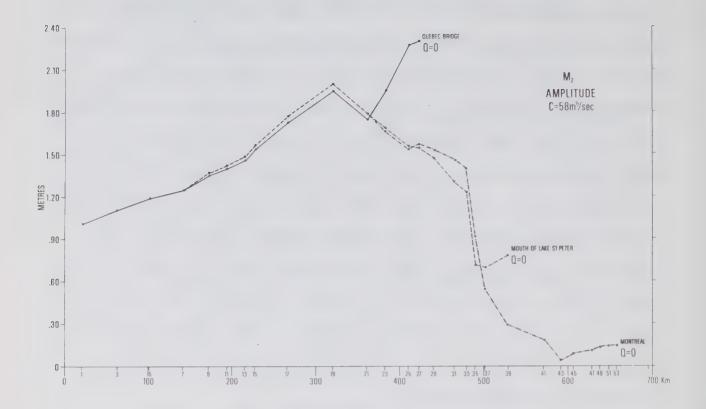


Fig. 30 The Amplification of  $M_2$  due to the Closure of the River at Québec Bridge, Trois-Rivières and King Edward Pier

appreciably the tidal régime. The results of such calculations are shown for  $M_2$  in Fig. 30.

If we blocked the river at Québec Bridge, the tide between Isle aux Couldres and Québec Bridge would be significantly enhanced by values which would reach up to 50 percent of the present value of M2 at the dyke site. At Québec Bridge M2 would have an amplitude of 2.31 meters compared to its actual value of 1.53 meters.

Blocking the river at section 39 (Trois Rivières) has the same effect. In terms of absolute magnitude, the enhancement of the tide is not as dramatic. M2 would take an amplitude of .78 meters compared to its value of .08 meters. Finally setting up an opaque barrier at King Edward Pier would once again enhance the tide. In fact, Montréal as it is presents an effective wall against the tide but the discharge present in the river has a marked damping effect as we notice in Fig. 26.

This brings to the fore the fact that we cannot build opaque barriers in a river where the mean discharge hovers around 8.5x10<sup>3</sup> m<sup>3</sup>/sec; this water has to be accommodated out. If it were at all possible in practice to let the water out at descending tide while impeding the flow of the tide upstream, the actual amplitude of the tide downstream of the dyke would be significantly less than the values predicted by our model. An added problem in the St. Lawrence is that water is not always in its liquid form and that the outflow of ice would create major problems in the design of semi-opaque dykes. A natural solution would be to take advantage of the increasing effects of pollution.

It would perhaps be possible in the near future to raise the temperature of the water to a point where it would never freeze in winter.

## CONCLUSIONS

Steady state conditions prevail on the average between Montréal and Sorel. The parameters present in the Equations of Hydrodynamics may be uniquely defined and derived in a physically significant fashion from the hydrographic charts, gauges and current meter readings. The gauge readings on their own supply, when averaged, the mean inclination of the water surface. The mean levels derived from gauge readings and hydrographic soundings supply data on the mean water depth at a station; this depth varies throughout the year but with lesser amplitude as one moves downstream. The motion everywhere in the river is subcritical and the tidal wave may move upstream at all times.

The tide moving up the St. Lawrence reaches its maximum amplitude between Isle aux Couldres and Isle d'Orléans. Shallow water constituents become significant between Isle d'Orléans and Sorel. The tide behaves as a mixture of standing and travelling waves between the mouth at Pointe des Monts and Isle d'Orléans; from there on it has the character of a travelling wave.

Use of the Equation of Hydrodynamics for time dependent motion in conjunction with current observation water level readings helps to delineate the mean value of the Chézy coefficient which holds in the area lying between Trois Rivières and Québec Bridge. The analysis of water level observations gives the values of the major constituents which should be reproduced in a mathematical model of the river. An analysis of the current observations taken at Isle aux Couldres and Québec Bridge by the Ship Channel Division reveals that these quantities could have been deduced theoretically with adequate accuracy using the equation of continuity as suggested by W.D. Forrester.

A one dimensional model of the river, neglecting the variation in discharge, the existence of the Saguenay River and the bifurcation of the channel at Isle d'Orléans reproduces adequately the major tidal constituents for a value of the Chézy constituent of C=58m½/sec which was deduced from the current observations of the Ship Channel Division and of G.C. Dohler. The model also gives values for the shallow water constituents which are created along the river but these values cannot be considered as accurate. Increases in the discharge do not significantly alter the tide between Pointe des Monts and Isle d'Orléans; beyond, the tidal amplitude is slightly reduced and the progress of the tide is retarded by up to 3 hours. The closure of the river by opaque dykes at any point would always enhance the tide downstream; semi-opaque dykes would increase the tidal amplitude as well, but to a lesser degree.

## ACKNOWLEDGEMENTS

Dr. F. Henry has read and criticized the manuscript. Mr. F.G. Barber has provided editorial comments and editorial assistance. Mr. J. Cookson has drawn a good part of the diagrams. Miss H. Bureau and a group of summer students have abstracted the mean widths and depths of the St. Lawrence River necessary for the elaboration of the numerical model. Messrs. R.W. Engley and J.D. Taylor have written the programs and carried out the computations pertaining to this undertaking.

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No. 19

The restoration of beaches contaminated by oil in Chedabucto Bay, Nova Scotia

E.H. Owens



Marine Sciences Branch

Department of Energy, Mines and Resources, Ottawa



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# Manuscript Report Series No. 19

# THE RESTORATION OF BEACHES CONTAMINATED BY OIL IN CHEDABUCTO BAY, NOVA SCOTIA

E.H. Owens

## ABSTRACT

Following the wreck of the "Arrow" in Chedabucto Bay, Nova Scotia, more than 150 miles of shoreline were polluted with Bunker C oil and a beach restoration programme was established which involved 30 miles of coastline. The beaches of this region vary from low-energy marsh environments to complex shingle spit systems, as well as many areas of eroding rock and till deposits. Only a few miles of beach are made up of sand-size material.

The contaminated sand beaches were cleaned relatively easily but although various manual and mechanical methods were implemented, no effective or efficient method of removing oil from shingle beaches was found. Oil on active shingle beaches was often buried up to a depth of 5 feet with clean and contaminated sediments having been thoroughly mixed by wave action. The amount of oil in these sediments was as low as 10 ppm and restoration of this type of shore involved the removal of large volumes of beach material. Where oil remained as a surface layer on the beach, a front-end loader proved to be effective in removing the contaminated layer.

In active beach environments normal wave processes will clean the beaches naturally unless there is so much oil or wave energy is so low, that the movement of sediments is prevented. On these paralysed beaches some attempt should be made to break up the surface so that wave action is able to rework and clean the beach sediments.



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# 1. INTRODUCTION

Following the wreck of the tanker "Arrow" on Cerberus Rock, February 4, 1970, the oil spill of Bunker C affected more than 150 miles of shoreline in Chedaducto Bay, which has a coast of 380 miles from Point Michaud in the north-east to Cape Canso in the south-east (Figures 1 and 2). The Project Oil Task Force initiated a beach restoration operation and a total of 24 miles of coast was initially selected for the work programme which was carried out by the Department of Public Works during the period April to October, 1970. In April, a coastal geomorphologist was assigned to aid the restoration project and this paper reports the results of the observations and investigations of that part of the programme.

This operation was the first attempt to restore non-sand beaches without the use of dispersants. On certain beaches the work was followed very closely in order to gain knowledge which would be of value to future projects of this nature. These instances are reported in detail along with a broader discussion of the more general aspects of the operation. A full account of the entire beach restoration programme which deals with each area where work was undertaken and the techniques employed has been prepared by the Department of Public Works site engineer (MacKay, 1970). This report is complimentary to that of MacKay but is concerned primarily with the geological problems rather than the engineering, economic, or social factors, although these are discussed in the final section. Most of the investigations were of a qualitative nature, though data collection was undertaken wherever possible to supplement the subjective observations.

The use of earth-moving equipment for beach restoration was not recommended in this area and was found to be only partially adequate in removing all contaminated material. Sand beaches were cleaned by manual and mechanical methods with success but no satisfactory mechanical methods were available for the restoration of cobble and shingle beaches. This lack of effectiveness can be attributed to the fact that it was not possible to remove all the contaminated material from the polluted cobble beaches and to the movement of oil alongshore and offshore which led to the recontamination of several of the restored beaches.

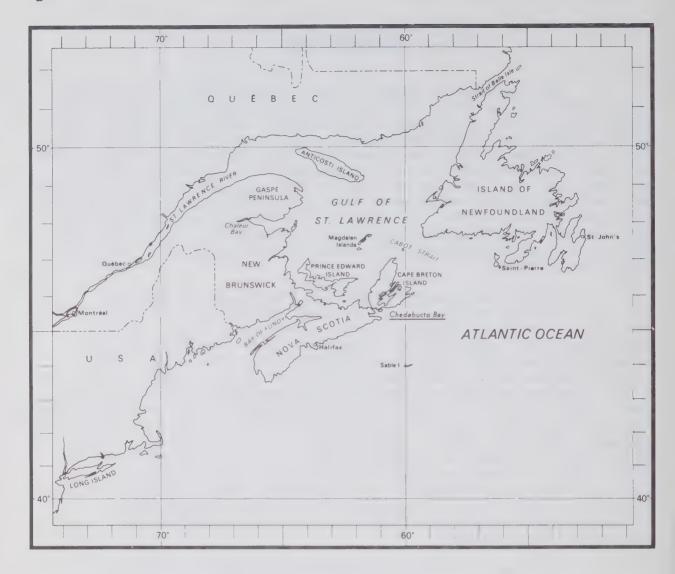


Figure 1. Chedabucto Bay location diagram.

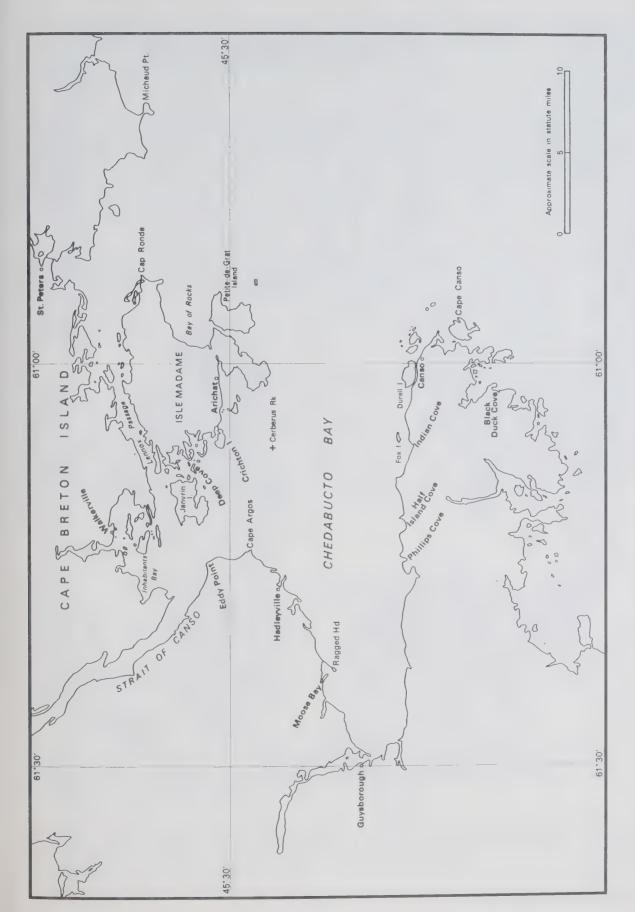
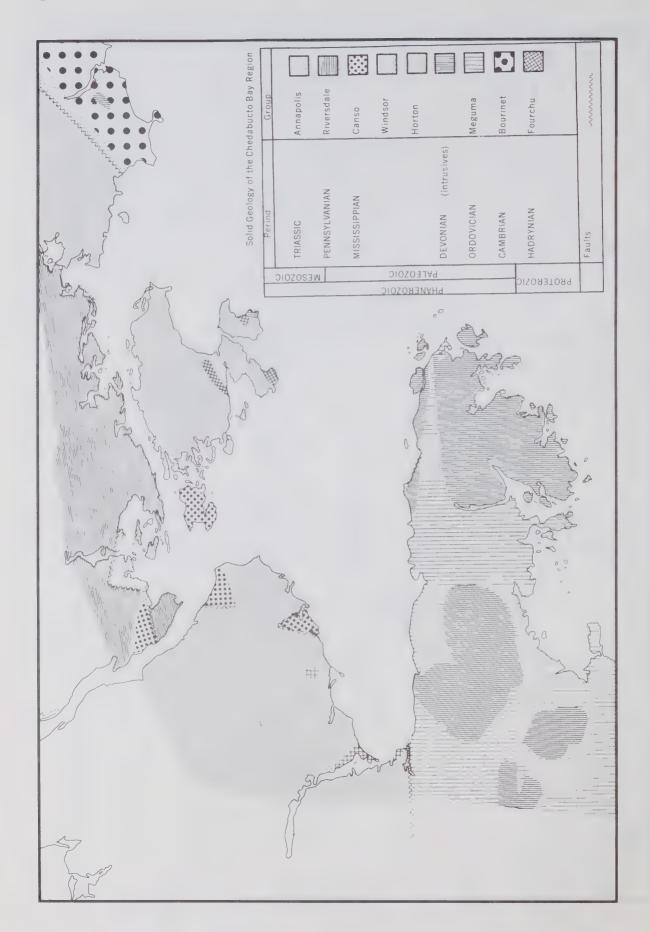


Figure 2. Locations within the study area.



# 2. PHYSICAL BACKGROUND

# 2.1 Geology

The major structural feature of this region is the Chedabucto fault or fault complex of the Acadian orogeny which separates the resistant metamorphosed rocks of the Meguma group in the south from the relatively less resistant late Paleozoic sediments which lie to the north (Figure 3). The fault zone defines the south shore of the Bay and can be traced eastwards onto the Nova Scotian Shelf (King and Maclean, 1970).

South of the fault zone the stable Meguma platform is made up of tightly folded Ordovician quartzites and slates which have been intruded by Devonian granites (Webb, 1969). This block has been eroded to give an uplifted peneplain which slopes gently southwards and the rise of sea level has led to an indented shoreline which is in marked contrast to the long straight coast of the fault zone.

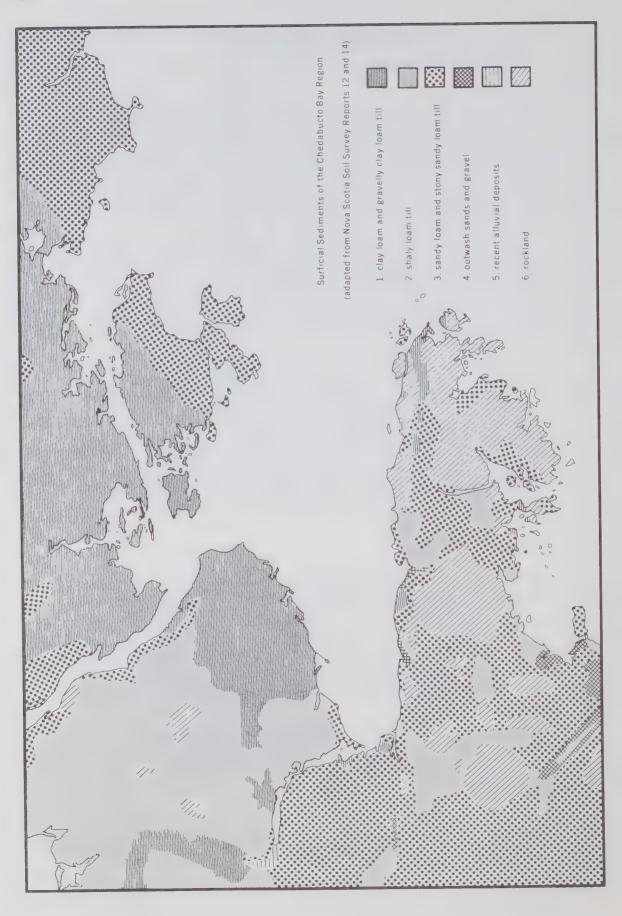
The Carboniferous beds which characterize the area north of the fault consist in part of a small basin within the uplifted basement and are relatively less resistant conglomerates, sandstones, shales, and limestones which have been folded and metamorphosed. Erosion along the fold axes provides a clear example of the general north-east/south-west trend in the Lennox Passage area. The erosion of these late Paleozoic rocks has led to the development of an undulating lowland area which has been drowned to produce an irregular coastline.

# 2.2 Surficial Sediments

The areas of bedrock exposure (Figure 4, unit 6) are confined largely to the more resistant uplands south of the fault zone. Where sediments are present in this area they are often very stoney and not deep. The remaining areas west of the Strait of Canso have a cover of locally derived glacial till (units 1, 2 and 3) with a few local deposits of outwash material (unit 4) or post-glacial alluvial sediments (unit 5). These latter units account for less than 3 percent of the land area (Hilchey et. al., 1964).

East of the Strait of Canso the till deposits are derived from the Carboniferous rocks and differentiation on a general basis produces a simple textural pattern related to the bedrock parent material.

The till deposits, which constitute the major unit of surface sediments in this region, were laid down directly by the ice as an unstratified and unconsolidated mixture of clay, silt, sand, gravel and boulders.



The earliest ice advances which have been traced (Grant, 1971) indicate a west to east movement across Isle Madame. This was superceded by a flow to the north and north-west from a centre on the Scotian Shelf while the final phase of glaciation is related to the recession of an ice cover which had moved south over the area and which disturbed much of the existing till deposits.

#### 2.3 Tides and Waves

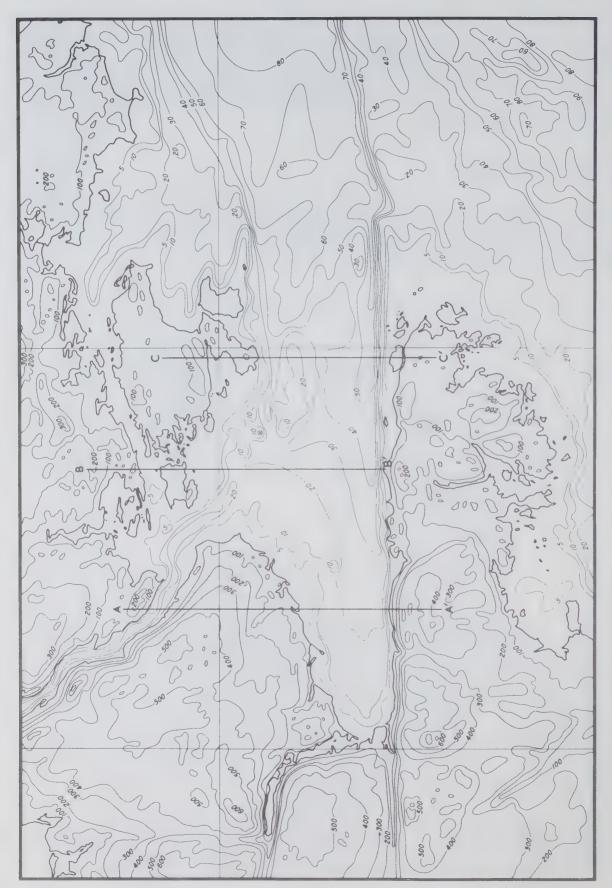
The tides of this area are semi-diurnal, with mean and maximum ranges in the order of 4.3 and 6.6 feet respectively (Table 1). The tidal range is an important characteristic of the littoral zone as it affects the width of beach over which wave action can take place.

Table	1	Tidal	ranges	in	Chedabucto	Bay	(Anon.,
		1970)				_	

Location	Mean Tidal Range (ft)	Large Tidal Range (ft)
Canso	4.2	6.6
Guysborough	4.5	6.4
Port Hawkesbury	4.4	6.9
Arichat	4.2	6.3
Petit de Grat	4.4	6.6
St. Peters Bay	4.3	6.1

The east coast of Canada is a storm wave environment (Davies, 1964) where the most important waves in the littoral zone are those generated by local storms and a brief review of wind data and the occurence of storms provides an indication of the seasonal variations in wave intensity. This part of Nova Scotia is exposed to the full force of waves from the east and south and those coasts which are directly open to these waves have high energy littoral environments. The relatively shallow areas in the north and west of the Bay (Figure 5) refract the incoming swell waves so that beach orientation is often a reflection of the offshore topography.

The summary of wind data for the meteorological station at Canso (Table 2) shows that the highest wind speeds are mainly from the north-west and occur in the period December to February. The generalized data indicate that the extent of the winter



depths Topography and bathymetry. Heights in feet: in fathoms. 5. Figure

season, in terms of the wave climate and beach form, is probably from November to April, while the period of accretion and summer profiles is May to October when winds are from the south-west.

This agrees in general with Neu (1971) who states that from December to May "most of the wave energy occurs in the 7 to 12 seconds (wave) period band propagated from SSW to WNW". As part of the Project Oil investigation, Neu recorded wave data in Chedabucto Bay for 33 days in March and April, 1970. The results indicated that waves with a period less than 9 seconds were generated locally or over the adjacent shelf areas but that ocean waves with periods longer than 9 seconds prevailed 80 percent of the time.

Table 2 Summary of the wind data, 1964 to 1970, Canso, Nova Scotia (Anon., not dated)

	Average of the Monthly Mean Speeds (mph)	Prevailing Direction	Average of the Monthly Maximum Recorded Speed	Dominant Direction
Jan.	14.3	NW	42	NW/NW
Feb.	14.8	NW	40	NW
Mar.	13.8	NW	35	NE
Apr.	13.4	NW	33	NW/SW
May	12.7	SW	33	SW
June	11.6	SW	29	SW
July	10.6	SW	27	SW
Aug.	11.9	SW	29 .	SW
Sept.	11.4	SW	30	SW
Oct.	13.0	SW	34	SE
Nov.	13.7	NW/SW	36	NE
Dec.	14.8	NW	38	SE

### 3. THE COAST OF CHEDABUCTO BAY

The basic coastal trends are structurally controlled while the actual detail of the shoreline results from the erosion and submergence of the fault block. Grant (1970) estimated that submergence is taking place at a rate of about 0.5 feet per century through a combination of the eustatic rise in sea level and crustal subsidence. This has led to the drowning of lowlying coastal areas to produce a complex and irregular shoreline for much of the region (Owens, 1971a).

#### 3.1 South Shore

The south shore of Chedabucto Bay between Canso and Guysborough is a straight, steep coast with a narrow offshore shelf (Figures 5 and 6). This is a resistant erosional shoreline composed largely of rock platforms and low cliffs with pocket beaches of shingle and coarse sand. The amounts of sediment in the littoral zone increase noticeably from east to west as indicated by the presence of spits and bars in the Salmon River - Guysborough area. South and west of Canso the irregular coast has resulted from the drowning of the southward sloping peneplain. Although this area was not examined in detail it was apparent that rock platforms with little beach material dominate the character of the shore. Occasional pocket beaches and spits interrupt the pattern but otherwise this area is an excellent example of a submerged resistant More sediment is available for reworking in the littoral zone of this region than along the south shore of the Bay, but in both cases there is a general scarcity of beach material.

#### 3.2 West Shore

Between Guysborough and the Strait of Canso the relatively uniform south-east slope of the subaerial and submarine topography (Figure 5) gives rise to a generally straight shoreline. Few bedrock exposures occur and the coast is made up of actively eroding till cliffs (Photograph 1) alternating with wide accretional shingle beaches (Photograph 2). In an attempt to develop an equilibrium shoreline, beaches have built up across river exits and embayments while erosion is active where there are headlands and higher relief along the shore.

The till cliffs, which have a maximum height of 60 feet, are easily eroded by subaerial and marine processes but provide the beach zone with relatively little sediment. The till is composed largely of clay and silt sized material which is removed from the base of the cliff as a suspended sediment, leaving a few cobbles and boulders in the littoral zone. The beaches of the area are generally long and wide with

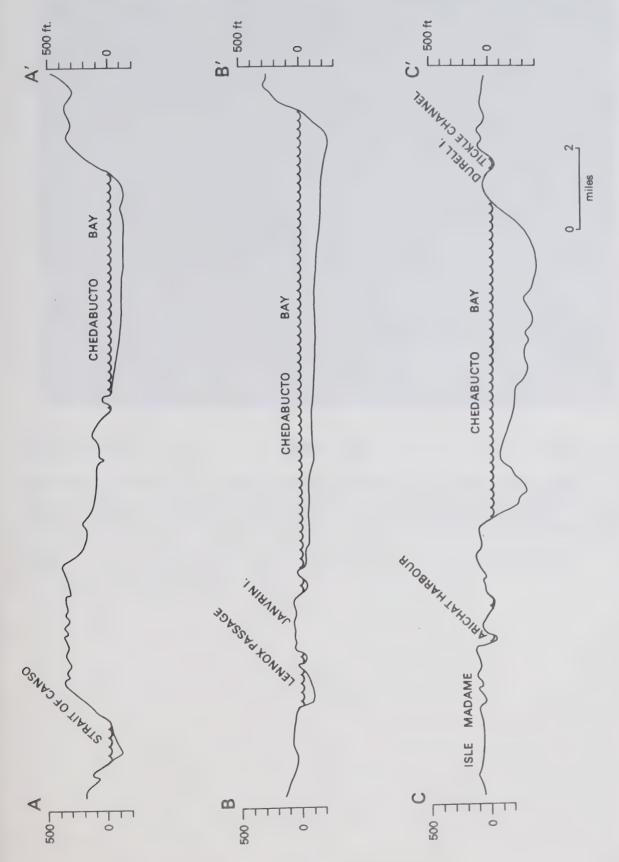


Figure 6. Profiles across Chedabucto Bay.



Photograph 1

CAPE ARGOS

March 27, 1970

Air view of a partially eroded drumlin and an actively eroding till cliff with a narrow but heavily polluted beach at its base. The beach material here is almost entirely derived from the erosion of the cliff.



Photograph 2

MOOSE BAY

June 9, 1970

Air view of the main section of the beach at mid tide before restoration. The main storm ridge is backed by a vegetated swale which in turn is replaced by a higher vegetation sequence. A squad of "slick pickers" is removing contaminated material manually (see also photograph 18), these are located by an arrow. No sections of this beach were paralyzed; although originally heavily polluted, most of the oil was buried or removed by normal wave action.

well-developed storm ridges and would appear to owe their growth to the landward movement of sediments from the offshore zone with the rise of sea level. It is doubtful that rivers and erosion of the till cliffs contribute significant amounts of coarse sediment under present conditions. Subsequent reworking and longshore sediment transport has produced some complex spit systems at several places along this coast. Where the till cliffs have narrow beaches at their base, waves attack the cliff directly and remove subaerially eroded material, but if the beaches are wide and there is a large sediment build-up, erosion is unlikely.

#### 3.3 North Shore

East of the Strait of Canso, as far as Point Michaud, the drowning of an undulating lowland has given rise to a complex series of islands and inlets in an area of little local relief. Much of the coast is protected from direct wave action and beaches have not developed in these sheltered zones because of the low energy conditions.

The River Inhabitants is a particularly good example of a drowned valley. It flows into Inhabitants Bay which was the original flood-plain but is now a shallow bay with poorly developed beaches in a sheltered wave environment. South of Inhabitants Bay a series of well-developed bars and spits results from the greater availability of sediments and the higher energy level of the littoral zone, as this area is exposed to wave action from the south and east. The great variety of coastal types in this region results from the irregular nature of the drowned lowland area and the pattern of islands, bays and inlets. The beaches and spits of south Janvrin Island and Jerseyman Island are a direct contrast to the adjacent but sheltered areas of Port Royal or Arichat Harbour which do not receive the full force of waves from the Atlantic.

For most of the sheltered areas the coast is typified by a narrow beach, less than 100 feet wide, of coarse sediments resting on a till platform. This is usually backed by natural vegetation or a low till cliff which may not be subject to marine processes under normal conditions. There is rarely any evidence of storm ridge development, though bars and spits of limited size occur in most areas.

On the exposed shores of south Isle Madame and Petit de Grat the resistant Proterozoic and Devonian outcrops produce a rocky coast devoid of sediments which is very similar to the south shore of the Bay. To the north, in the Bay of Rocks, this gives way to a series of wide sand and cobble beaches oriented towards the refracted incoming waves from the east. Along the north shore, to the east of Lennox Passage, a large number of complex spits, bars and tombolos result from the

reworking of sediments in an energetic wave environment. As with the West Shore there is little indication that terrestrial sources are a major contributor of coarse material so that again the offshore zone appears to have been the major source for littoral sediments.

## 3.4 Summary

Apart from a few exceptions, in the north-east part of the Bay the beaches consist mainly of cobbles and boulders which are thrown up by wave action onto the higher parts of the beach to form a storm ridge. When these investigations commenced in April, steep winter profiles characterized most areas but by July these had been replaced by the more gently sloping accretional summer profiles, and these were still evident on all beaches in November. With this change in profiles came a marked increase in the amount of coarse sand and gravel in the intertidal zone.

Sand beaches are restricted to the Bay of Rocks, Point Michaud, and Blackduck Cove area, all at the eastern end of the Bay. These are wide, long beaches with shallow offshore zones and have well-established dune vegetation in the backshore areas.

The shoreline of Chedabucto Bay may be described as an erosional coast of rock and till exposures with pocket beaches. The growth of large accretional features is evident along the exposed north and west coasts but these beaches are supplied with little coarse material under present conditions. The sheltered coasts along the north shore have a limited sediment supply and are in a low energy wave environment with evident erosion in many areas as the region undergoes submergence.

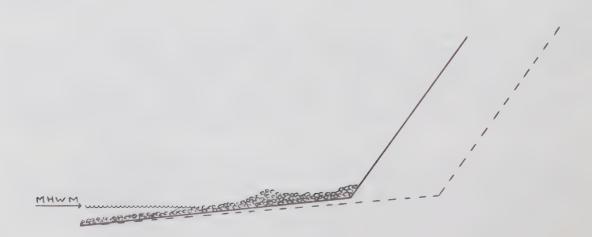
### 4. SHORELINE STABILITY

"A beach is nothing more or less than a protective apron of rock waste fronting the land. Where beaches are wide and in equilibrium, erosion is unlikely. Where they are narrow and starved of material, wave attack is directed against the land with little to absorb its full force" (Kidson, 1966).

Coarse beach sediments, larger than 3 inches, were moved landward from the offshore zone with the rise in sea level and any sediment loss by erosion or excavation must be equated with the capability of natural replacement. Under present conditions, material of this size does not move into the littoral zone from offshore, thus the supply of coarse material is very limited and large-scale sediment losses cannot be replaced naturally.

#### 4.1 Till Cliffs

At the base of till cliffs which are in active retreat there is usually a narrow cobble and boulder beach (Photograph 1). This material is derived largely from the erosion of the free face, which is a limited supply source as most of the till is clay and silt sized material. The retreat of the cliff leads to the development of a small platform at the cliff base on which rests the thin cover of beach material. In the sketch below it can be seen that on removal of the beach mantle the till base is exposed and becomes subject to erosion by marine agents. The protection that is afforded by the thin beach mantle is not great but with the loss of this mantle the level of the till platform is lowered by marine erosion and this in turn leads to a temporary acceleration of cliff retreat until an equilibrium condition is achieved by erosion of the free face.



Those till cliffs which have wide beaches at their base are not normally subject to direct wave action and erosion is a result of subaerial agencies, although the talus may be removed by littoral processes. These cliffs would be unaffected by beach loss unless large volumes of sediments are excavated to the extent that marine processes could begin to erode the base of the cliffs.

In sheltered areas, such as Inhabitants Bay, wave action is limited so that often, marine processes are not active in cliff erosion except under storm conditions. In these instances, subaerial rather than marine erosion may be consistently active though the infrequent storm would produce more dramatic results in the long term. The removal of the beach sediments which front these cliffs will have the same effect as on more exposed shores.

### 4.2 Beaches

Beaches are constantly changing in response to a variety of processes which may alter in intensity with the season, tidal cycle, or weather conditions. This is a complex environment which is still only partially understood.

The beaches of the region are made up largely of cobbles and pebbles and these sediments are being eroded, reworked and transported as the littoral zone strives towards the equilibrium which the processes demand. Although sand and fine sediments are fed into the coastal zone, particularly during the summer months, these are not a major contributor to the growth of beaches and large constructional features, except for areas at the east end of the Bay.

The loss of large volumes of cobble and larger size material would be harmful to the shore environment in this region. An instance where this has occured in a similar situation is reported by Robinson (1961) at Hallsands, on the south coast of England. As with most of the British coast, this beach owes its origin to the reworking of deposits during the post-glacial rise in sea level and is not now fed with coarse sediments from the offshore zone. At this locality there is no significant sediment supply from rivers or coastal erosion so that the beaches consist mainly of "fossil" material. Almost 500,000 cubic yards of shingle was removed in the period 1897 to 1902 along a half mile length of shore. This led to a recession of up to 20 feet between 1907 and 1957 in those sections of schist cliffs which had little protection at their base. The actual beach area to the north of the cliffs was lowered by as much as 12 feet and only the most southerly sections appear to have recovered in any way (Figure 7, profile 6).

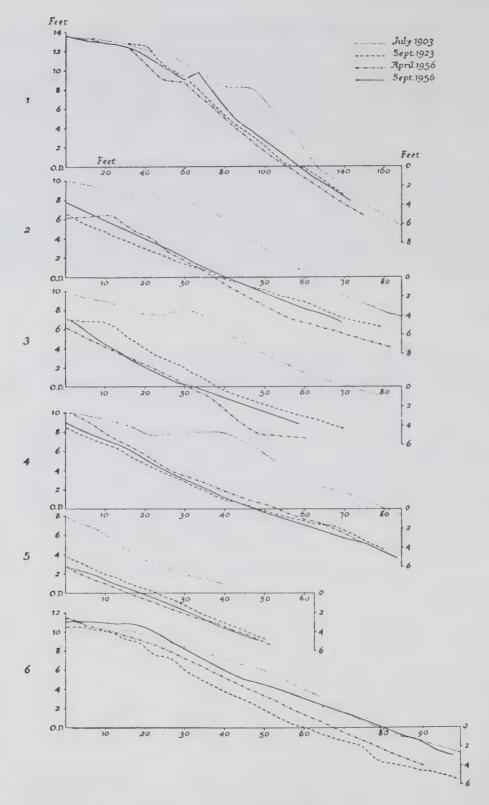


Figure 7. Beach Profiles from Hallsands, U.K. (Robinson, 1961)
The changes in beach profiles along Sections 1-6
between 1903 and 1956. Vertical exaggeration
for Section 1 is 1:5 and for remainder 1:2.5.

Although material removed from a beach may be replaced by longshore drift, providing the loss is not too great, there would be a net decrease for the shore as a whole in a region such as Chedabucto Bay, as this material cannot be replaced rapidly. On beaches which have a marked movement of sediment in one direction, should material be taken from the up-drift end, there is a danger that the beach may not fully recover as this is a section which would not normally receive much "new" material. In the same way, removal of sediments from a spit near the point of attachment may be harmful as most of the accumulation is concentrated at the distal end. This would apply to any constructional feature such as a bar, tombolo, or foreland which is formed by the longshore movement of material.

While it is important to consider the loss of material in terms of the plan form of the beach, more critical perhaps is the effect on the beach in profile. The storm ridge, or the sand berm, is built above the normal high water mark and material lost from this zone will only be replaced during the infrequent occasions when wave processes are active on these sections of the profile, providing that material is available for replacement.

## 5. OIL ON THE SHORE

#### 5.1 Distribution

Those coastal areas which were contaminated by Bunker C from the "Arrow" are outlined in Figure 8. This diagram is only accurate for the Bay itself as many of the areas in the north-east and to the south and west of Cape Canso were only surveyed at the reconnaissance level and only those areas where oil was actually observed are indicated. The shore affected by oil was more extensive than is shown and it should be noted that contamination extended as far as St. Esprit, on the north shore some 5 miles east of Point Michaud, and Sable Island, which lies over 100 miles south-east of the Bay (Figure 1). The major spills took place in February and March but oil leaked from the wreck throughout the spring and summer.

The distribution diagram does not take into account the severity of the contamination or the frequency of reoiling. The shores of Janvrin Island, Crichton Island and Jerseyman Island and the south shores of Isle Madame and Petit de Grat were heavily polluted in February and subject to frequent reoiling in the following months. The area of Inhabitants Bay was not contaminated by the end of February but was subsequently heavily oiled. To the north and east of Petit de Grat contamination was relatively light and not continuous. The shore of the east section of Lennox Passage was not polluted due to the construction of a dam in the middle section.

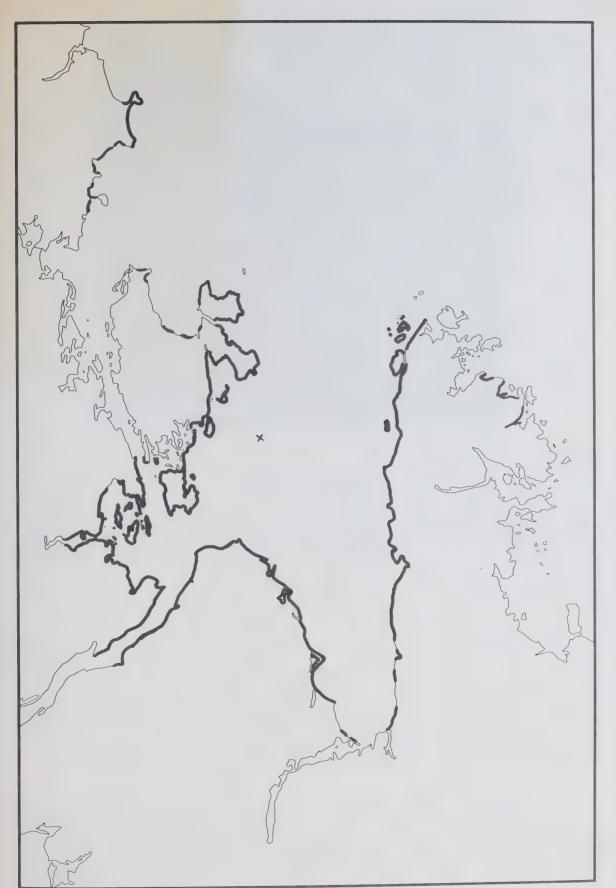
Along the west coast of the Bay, from Cape Argos to Guysborough, most of the shore was contaminated but the amounts of oil on the beach decreased towards the south-west. Similarly on the south shore, the area around Canso was severly polluted while to the west of Half Island Cove the degree of oiling decreased and some sections were only lightly oiled.

South of Cape Canso several of the inlets and embayments were badly polluted though a full survey of this area was not carried out. The distribution diagram does not indicate the true extent of oil on this section of coast.

## 5.2 Behaviour of Oil on the Shore

Field observations by Asthana and Marlowe (1970) and by Drapeau (1970) during February and March were carried out, in part, to provide an understanding of the nature and behaviour of oil on the coast. Their conclusions are summarized below.

Floating oil striking the shoreline behaves differently on different types of shore material.



Cerberus Rock Distribution of oil on the coast. Figure 8.

- a) Bedrock is coated by a uniform layer of oil in the intertidal zone. The oil is not removed during tidal submergence but when exposed, flows into crevices and hollows.
- b) Boulders and boulder beaches are affected in a similar way.
- On sand and gravel beaches the oil remained on the surface and behaved as discrete sedimentary particles with unique hydrodynamic characteristics. The oil particles tended to float free and become concentrated along the high water line. Oil did not permeate sand but acquired a surface coating of sand particles (Photograph 3) while on gravel it was observed to have permeated as much as 18 inches. Where oil had been buried and later exposed by the normal accretion and erosion processes, the oil layers were less readily eroded and formed ephemeral ledges (Photograph 4) which were seen to crumble under the influence of gravity and sunlight.

Drapeau also noted that within one month of contamination, gravel bars directly exposed to wave processes cleaned themselves naturally and effectively, while in protected areas the beaches remained polluted. This is particularly evident in Photograph 5 of Blackduck Cove. The lagoon remained heavily polluted throughout the summer while the exposed side of the spit was "cleaned" by the end of March. In many of the exposed areas the intertidal zone was rapidly cleaned leaving a concentration of oil where it had been deposited above the high water mark beyond normal wave action (Photographs 6 and 7). Although oil polluted much of the upper part of the intertidal zone, wave action is continuously active in this zone and any oil laid down above mean high water mark would remain undisturbed by normal processes. During periods of storm waves or spring tides the remainder of the beach would be subject to wave action so that eventually all contaminated areas would be cleaned naturally. On cobble and shingle beaches, oiled material originally in the intertidal zone was often moved up the beach by storm or swell waves and deposited as part of the storm ridge. This may have occurred on several occasions so that with subsequent exposure the contaminated material would appear as layers interbedded with clean sediments (Photograph 4).

Exposed beaches are generally capable of self-cleaning (Johnston, 1970) but there is a threshold beyond which beaches become "paralyzed" (Drapeau, 1970). This situation may arise as a result of either heavy contamination or insufficient wave action. The beaches of Arichat Harbour and Inhabitants Bay, for



Photograph 3 HADLEYVILLE No. 2 July 10, 1970

The east end of the beach was made up of coarse sand and the oil remained as large pans which had a surface layer of fines or was buried to a depth of several feet. No oil was visible in the intertidal zone.



Photograph 4 HADLEYVILLE No. 2

July 10, 1970

On the west end of this beach at low tide. The layers of oil, when exposed by wave action, are more resistant to erosion but quickly collapse through sunlight and gravity.



Photograph 5

BLACKDUCK COVE April 25, 1970

Air view of spit, lagoon, and cove from the south at high tide. The lagoon is virtually cut off from the sea at low tide. This area was heavily polluted though the exposed section of the spit has been virtually cleaned by wave action whilst the beach in the lee of the spit remains paralyzed.



Photograph 6 HADLEYVILLE No. 2 July 10, 1970

The central section of the beach was characterized by a thick, continuous layer of oil ten to fifteen feet wide, above the high water line. This upper part of the beach, rarely affected by waves, remained paralyzed while the intertidal zone was cleaned by wave processes.



Photograph 7 JERSEYMAN ISLAND June 21, 1970

A badly contaminated cobble beach near the high water mark. Although this beach has been partially immobilized, wave action has begun to clean the area below the high water line by abrasion and burial. With storm waves the remainder of the beach will be combed down and the cleaning process will be extended to include all the contaminated material.

example, were paralyzed by a combination of both these factors while the heavily polluted but exposed beaches of Crichton and Jerseyman Islands remained mobile.

Four samples were collected for analysis of oil content by volume to determine the oil/sediment ratio on different beaches. Two were taken from the exposed, mobile beach at Moose Bay near the high water mark. The first was at a site which appeared clean and the second from material which appeared badly oiled. The results indicated 9 and 90 parts per million of oil, respectively. Two more samples were taken from the intertidal zone of a "paralyzed" beach at Arichat. These provided values of 4 and 5 per cent of oil. All the samples were from cobble beaches with a material size range of 1 to 6 inches.

## 5.3 Previous Beach Restoration Projects

The only restoration projects before the Chedabucto Bay operation resulted from the "Torrey Canyon" and "Santa Barbara" spills. Neither of these are directly comparable as the coasts of south-west England were cleaned with dispersents while the wide, firm, sandy California beaches were restored with relative ease.

The work of the Ministry of Technology, Warren Spring Laboratory, dates from 1960 and their involvement in the "Torrey Canyon" clean-up provided a practical demonstration of the various restoration techniques which are reported by Wardley Smith (1968a, 1968b, 1969, 1970). The fundamental conclusion from their experience is that no one method is adequate because of the variety of coasts and the amounts and types of oil involved. Despite this, it was stressed that the best defence is preparedness so that if oil does reach a shore the correct action may be taken immediately.

The various methods of restoration which were discussed are summarized below.

- a) Burning. In most instances this is not satisfactory because of inefficiency and cost. Most oils, particularly Bunker C, do not burn readily and require a great deal of assistance; even then the oil burns very slowly. Heating often makes the oil more mobile so that it penetrates deeper into the beach and contaminates more material.
- b) Absorbtion. The use of sawdust, peat moss, or similar material was found to be very useful but costly in terms of the labour required to spread and retrieve the absorbant.
- c) Surface Coating. A crust could be formed on the surface of the oil with the application of large

quantities of powder or fine particles. This stabilized the surface but was not satisfactory for preventing the oil from moving and did not contribute to the removal of the oil.

- d) Mechanical Removal. This method was found to be satisfactory for the removal of surface oil but not all the contaminated material was removed by the machinery. In combination with dispersants this method proved effective on cobble beaches. Machinery was developed to pick up contaminated material because of the unsuitability of available excavating or earth moving equipment, but this was elaborate, sometimes expensive, and not always successful.
- e) Manual Removal. Effective on most beaches but expensive because of labour costs. Techniques include raking, shovelling and spreading absorbant materials.
- f) Dispersants. Very effective in removing the oil, especially on cobble beaches, but toxic to marine life.

Each of these methods was considered in terms of applicability to different shoreline types.

- a) Marshes. Manual removal would be best but burning may be effective during the non-growing season. However, unless the oil is thick the marsh is best left so that it could recover naturally.
- b) Mud Flats. Although these areas would be left alone where possible, if restoration is necessary mechanical scraping or dispersants could be used for thick deposits.
- c) Sand. Again it would be best to leave cleaning to natural processes if possible but if action is necessary, manual or mechanical raking would be effective. If the beach is heavily polluted, material could be pushed into the sea at low tide and the beach then sprayed with emulsifier. The sand would be reworked, cleaned, and returned to the beach without any sediment loss to the littoral zone.
- d) Shingle. Apart from natural processes only dispersants could rid the beach of all contaminated material. No machinery was available which could clean the beach without the aid of dispersants.

e) Rock. These areas are best left to the action of waves and weather but burning or dispersants could be applied if necessary.

Following the Santa Barbara spill on the California coast the United States Federal Water Pollution Control Administration (FWPCA) initiated an evaluation of earth moving equipment in oil contaminated beach restoration operations (FWPCA, 1970). This project was carried out on wide, firm, flat sand beaches which had oil on the surface or in the surface layers. The results of this technique analysis are given below.

- a) Grader/Scraper. This was found to be the best of the techniques evaluated. Contaminated sediments were pushed into wind-rows by the grader and then lifted by the scraper. This technique removed the least uncontaminated material but spillage from the scraper required a following pick-up crew. The grader became stuck on coarse sand unless expensive flotation tyres were in use and accuracy decreased if traction was low.
- b) Scraper. Used on its own the scraper had a high spillage and like the grader, it required a flat beach and became stuck easily.
- c) Grader/Front-end Loader. The loader was used to remove the wind-rows. Its performance in general is outlined below.
- d) Front-end Loader. This machine was the least efficient of those tested as it removed too much uncontaminated material and had a high spillage. The same defficiencies apply to bulldozers and tracked loaders as bulldozers were found to "grind the oil several feet into the sand".
- e) Ramp-Conveyor System. This method was developed to remove the material after the grader formed the wind-rows and was found to be valuable for very large operations.

Neither of these evaluations are directly applicable to the problems which were faced in Chedabucto Bay, as no wide, sand beaches required mechanical restoration and dispersants were ruled out of the programme. The conclusions from the work in Britain indicated that natural cleaning is best but even if restoration is necessary, shingle beaches were not cleaned properly by machines alone. The evaluation project by the FWPCA showed that the grader and bulldozer were the least efficient of the earth moving equipment tested. The performance was particularly bad if the vehicle had tracks, as clean and contaminated sediments were mixed by spillage and grinding.

### 5.4 Recontamination

Following the grounding and sinking of the "Arrow", large slicks contaminated over half of the shoreline in the Bay and leaks from the wreck continually spilt small amounts of oil throughout the spring and summer. These slicks were a great deal smaller and thinner than those of February and March but this oil did recontaminate many of the beaches along the north shore area.

Other sources of oil which led to shoreline recontamination were from the coastal areas themselves. The reworking of contaminated sediments by wave action led to the release of oil onto the water in small amounts. This was exemplified by the action of tracked vehicles working in the intertidal zone in the early phases of the project before better machine-operation techniques were developed. In particular at Arichat, which was the first area to be worked, a substantial amount of oil was released into the sea and this led to the reoiling of adjacent beaches (Section 6.4.1).

Oil was often contained in rock hollows and crevices above the limit of normal wave action. With spring tides or storm waves, some of these pools of oil were flushed, leading to recontamination of adjacent alongshore areas. This was judged to have been the reason for the reoiling of Indian Cove (Section 6.4.3).

## 6. BEACH RESTORATION

# 6.1 Deep Cove, February 1970

In mid-February, shortly after the sinking of the "Arrow", the Canadian Armed Forces bulldozed part of the bar at Deep Cove in an attempt to clean this beach (Photograph 8). It was estimated that some 3,000 cubic yards of sediment were excavated and removed during this operation (Asthana and Marlowe, 1970). This trial was not successful as oil and sediments were thoroughly mixed and not all the contaminated material was removed. This operation was also ineffective in terms of recontamination as this area was reoiled within a few tidal cycles and on several occasions subsequently.

The beach which was subjected to bulldozing is a narrow bar which joins two islands and the removal of large volumes of sediment was considered very damaging to the stability of the foreshore (Asthana and Marlowe, 1970). As a result of this experiment a recommendation was made to the Task Force that bulldozing should not be continued and that beach restoration should be restricted to manual methods (see also 6.4.6).

### 6.2 Selection of Beaches for Restoration

Before the assignment of the coastal geomorphologist to the programme, a series of beaches had been selected for restoration. These beaches included all accessible nationally rated shorelines (these are rated on a recreational basis), as well as community beaches in the area. The information regarding the location and extent of these beaches was supplied by the Emergency Measures Organization. Additional sections of coast were included in the programme as a result of public requests. In all, 30 miles of coast were restored by the Task Force.

The restoration programme was carried out by personnel of the Department of Public Works who determined whether beaches could be restored by manual or mechanical methods. Lightly oiled beaches were restored by squads of "slick pickers" (Section 6.3) while contracts were drawn up for the restoration of heavily contaminated sections. These were awarded to private companies after bids had been tendered.

No geological criteria were included in the decisions regarding which beaches were to be restored by the Task Force. Only after a section of coast had been designated were these criteria considered and the geological input was largely restricted to recommendations concerning those beaches actually under contract rather than advice related to the restoration programme as a whole in terms of the selection of beaches.



Photograph 8

DEEP COVE

March 1, 1970

Air view from the north of armed services restoration work. This is a long, narrow bar which joins two islands and is open to wave action from the south and east. At the east end of the beach a spit has grown north to partially close off a small lagoon which is being infilled with mud and silt transported alongshore. The sea at the near side of the bar and the lagoon, are ice covered. Material has been piled up by the bulldozers for later removal.

### 6.3 Manual Restoration

Those sections which were designated for restoration but which were only lightly oiled were cleaned effectively by squads of "slick pickers". These units, local labour under the Department of Public Works supervision, removed oil and contaminated material with shovels and rakes (Photograph 9). The material was placed in plastic bags which were collected and removed to an approved dump site.

Although this method of restoration is comparatively expensive in terms of labour costs, it is most efficient and effective as only contaminated material is removed. The loss of sediment from a beach is low, so this has very little adverse effects on beach stability and except where the oil is deeply buried, it is possible to remove virtually all the contaminated material.

The sand beaches in the Point Michaud area and in the Bay of Rocks were successfully cleaned by this method. These shores were not in danger of recontamination and remained clean throughout the summer. Shingle and cobble beaches were harder to clean as the oil did not remain in cakes and pans on the surface but this method is still preferable to the use of heavy machinery on lightly oiled beaches.

Any material which was buried and not removed manually would probably be exposed in the winter or spring when the beach is combed down under vigourous wave action. This oil would be subject to wave action throughout the winter and by the following summer there should be little evidence of any contaminated material on the exposed beaches.

#### 6.4 Contract Work

This section includes detailed accounts and analyses of the contract work on five beaches and a more general report on some of the other beaches where machinery was used. Not all of the sites which were contracted or where machinery rental was used are dealt with (MacKay, 1970) but the operations reported cover a representative cross-section of the shoreline areas in Chedabucto Bay where projects were carried out. The location of the sites reported here is given in Figure 2.

# 6.4.1 Arichat

The contract for a 3,700 foot section of beach was awarded to a low bid of \$4,479 and the work was carried out over 9 days between April 30 and May 11, 1970 using a fixed-blade International TD15 bulldozer and an International TD9 skid shovel; both vehicles were tracked. A total of 422 cubic yards of material was removed and 40 cubic yards were brought in as clean replacement at Le Noir Forge.



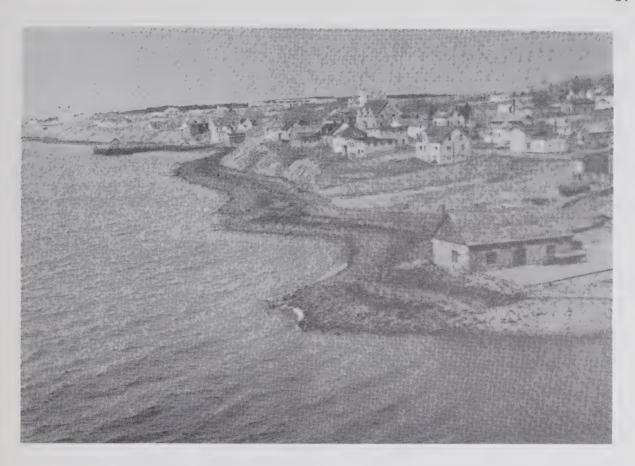
Photograph 9

SLICK PICKERS May 29, 1970

Air view of the restoration of a sand beach manually. Oil and oiled seaweed is shovelled into plastic bags which are collected for dumping. This method is very effective as spillage is negligable and only contaminated material is removed. Description of the contract area

This is a sheltered location not exposed to the direct action of waves from the Atlantic with a tidal range in the order of 4 to 6 feet (Table 1). All of the contract area was badly oiled and the beach zone was effectively paralyzed; that is, oil prevented the normal movement of sediments by wave action. In detail the section may be subdivided into four units:

- Between the Arichat wharf and Le Noir Forge the 1) shore is characterized by a 20 to 25 foot till cliff (Photograph 10). The width of the beach zone averages 40 feet at low tide and is made up of sediments derived directly from the subaerial and marine erosion of the cliff. This beach mantle was rarely more than a foot thick, mostly gravel and cobble size with a few boulders and overlay the till platform which resulted from the retreat of the cliff. Sections of the backshore are undergoing active erosion but in general the supply of material into the littoral zone is minimal. As the stone content of the till is low, most of the eroded material is removed as suspended sediments.
- 2) A small prograding beach, about 100 feet long and 60 feet wide, fronts a low backshore in the area immediately west of Le Noir Forge. This beach appears to have resulted from accretion which occured as longshore sediment transport was interrupted by the small headland on which Le Noir Forge was constructed.
- 3) The Le Noir Forge headland is a till bedrock area which has a beach of till-derived material at its seaward end. The two small beaches on either side of the headland are of gravel and are areas of more rapid accretion.
- 4) The east half of the area is made up of a series of three mid-bay bars and low active till cliffs. The bars have built out on a shallow platform which has a maximum width of about 150 feet. Sediment accretion on these small bars has deprived the intervening areas of material so that the narrow beaches have not prevented wave erosion of the backshore cliffs. The till cliffs are generally lower than the west section, being between 5 and 15 feet high. Various sections have been riprapped or protected with wood structures to prevent the erosion of property.



Photograph 10

ARICHAT

May 2, 1970

Air view of Le Noir Forge and west section at low tide. The bulldozer had made a "road" along the base of the active till cliff at the high water mark. This section was heavily contaminated and the beach zone was paralyzed. The contract area ended at the wharf.

#### Restoration

This contract was the first to be issued and work began on April 30. Contaminated material in the intertidal zone just to the west of the Le Noir Forge was heaped into piles and a "road" was bulldozed along the high water mark westwards toward the wharf. Some concern over the action of the contractor was expressed and it was decided not to remove any material beneath the cliff west of Le Noir Forge. As the "road" had already been bulldozed part of the way, the disturbed material on that section would be pushed against the base of the cliff but the contaminated material which had been piled up on the beach adjacent to Le Noir Forge would be removed.

On May 2, the "road" was bulldozed through to the west end of the contract section, near the public wharf and during May 2 and 3, the disturbed material in this section was pushed against the base of the cliff up to 4 feet above normal beach level. The piles of material west of Le Noir Forge were removed on May 5. The Department of Public Works site supervisor reported "fresh" oil on May 3 and 4 on those areas of the beach which had been subjected to machinery.

On May 6, a decision was made to remove and replace a section by Le Noir Forge which had been excavated to the till bedrock and areas were outlined in the east half of the contract where material should not be disturbed. During May 6 to 10, the areas designated for cleaning in the east section were dealt with by the contractor and on May 11 the contractor was released.

## Effectiveness and Results of Operation

a) The use of a bulldozer to drive a road through the west section thoroughly mixed contaminated and uncontaminated sediments. As the beach was scarcely wider than the resulting road the contractor should have removed the layer of contaminated material while progressively moving westward.

This action highlights an important point that the site supervisor should be aware of the implications of various actions so that he can make on the spot decisions which will be carried out by the contractor in order to prevent action which may be undesirable or endanger the stability of the beach.

b) The removal of material from the beach at the foot of the till cliffs may have accelerated basal erosion. The beach consisted of a thin cover of till-derived sediments overlying a till platform which was exposed as the "road" was bulldozed across the beach. In order to minimize the possible adverse effects of this action the disturbed sediments were pushed against the base of the cliff.

No material was removed from this area but the disturbance of the beach material led to the washing out of the fines and the release of oil into the littoral zone. The fines were transported along shore, in a band some 6 feet wide adjacent to the beach, while the oil was redeposited in thin layers near the high water line (Section 5.4). The section would have been better left, although at least this beach was no longer paralysed.

- c) In the area of beach to the west of Le Noir Forge some 62 cubic yards of material was removed by the contractor, approximately a 12 to 18 inch layer of sediments. At the completion of the contract, this area was relatively clean, with only a small amount of reoiling near high water mark. On May 26 and 27, the beach was reoiled by a slick (not from the "Arrow") 30 feet by 5 feet which was 3 inches deep in parts. Peat moss was placed over the oil on the 28th and the contaminated material was removed manually to a depth of 12 inches on June 2. Some 2 feet of beach was removed in all and even those sections re-cleaned on June 2 had been slightly reoiled by June 3.
- d) On the seaward edge of Le Noir Forge promontory, some 20 cubic yards of material were removed. This exposed the till base and the replacement of 40 cubic yards of larger, boulder sized material has increased the stability of this small section of beach. The clean material used as replacement has been contaminated.
- e) In those areas east of Le Noir Forge which should have presented no difficulty for the contractor, i.e. only a shallow surface layer of material was contaminated, the use of a bulldozer with a straight blade was demonstrated to be particularly unsuitable. As material is pushed forward by the blade, clean and contaminated material is mixed. Spillage around the edge of the blade contributes to the inefficiency of the method as this material must be removed by a subsequent pass and is often ground in by the tracks of the vehicle. The only way in which the bulldozer could operate effectively was to pile up material for removal by other equipment and this tended to mix sediments rather than remove the contaminated layer.
- f) The cleaning of selected sections of a shoreline is not effective because oil from adjacent contaminated areas as well as from offshore, can lead to reoiling. All of the beaches subjected to cleaning had been reoiled by June 4, either from alongshore or offshore.
- g) The two basic criteria that all contaminated material should be removed and that there should be no danger of recontamination were not realized on this contract section.

h) The Arichat beach was badly contaminated and much of the littoral zone was "paralyzed". From this aspect, the restoration operation at least made the beaches mobile once more, though this could have been achieved by the use of a tractor-drawn rake or hoe. The contract was the first to be awarded and many of the lessons learnt here were applied to later projects.

## 6.4.2 Blackduck Cove

The contract for this 4,600 foot section of coast in the north and north-east of the Cove (Photograph 5) was awarded for \$6,000 and the operation required 9 days between April 30 and May 12, 1970, using a fixed-blade Caterpiller D6C bulldozer and a Caterpiller 950 wheeled front-end loader. A total of 4,469 cubic yards of material was removed and 360 yards of clean boulders were brought in along the north shore.

Description of the contract area

The Cove is on the south coast of Nova Scotia but is not directly exposed to waves from the open sea. At the east end of the shallow cove a wide, medium to coarse grained sand beach is backed by a vegetated berm and brackish marsh. The beach is generally about 100 feet wide but at spring tides may be up to 300 feet wide. The abundance of sand in the offshore zone and the shallow nature of the cove indicate that there is an ample supply of sand-sized material for beach replenishment. The berm and backshore areas are stabilized by grass. Both ends of this section are areas of mud and silt accumulation presumably where the fines have been deposited by alongshore movement away from the centre of the beach.

The oil on the sand beach had "paralyzed" the sediments above high water and though some self-cleaning had taken place in the intertidal zone, oil had mixed with sand and seaweed to form large immobile cakes (Photograph 11).

The northern shore of the cove is a low silt/sand area which has a surface cover of large glacially derived boulders which protected a road and houses on the low till backshore from marine erosion. This area was heavily polluted and paralyzed.

#### Restoration

The contractor commenced work on April 30, 1970, in the sandy beach section. Contaminated material was bulldozed from the intertidal zone into large piles above high water for subsequent removal by the front-end loader (Photograph 12).

Concern was expressed over possible damage to the backshore vegetation by the machinery and the contractor thereafter maintained one track for the trucks.



Photograph 11

BLACKDUCK COVE

April 25, 1970

Air view of sand beach at the east end of the cove near high tide. Oil is present as a thick, continuous layer along the high water mark and along the edge of the vegetation. Thick patches of oil and seaweed "cakes" cover much of the intertidal zone but the beach is still mobile. The cobbles and boulders along the north shore in front of the cottages are also heavily polluted but this section is paralyzed.



Photograph 12

BLACKDUCK COVE May 2, 1970

Air view of beach restoration work at low tide. A bulldozer has pushed the contaminated sediments above the high water mark for subsequent removal by a front-end loader. The contaminated material along the edge of the vegetation was removed by the bulldozer scraping backwards down the beach in order to minimize damage to the backshore. For the same reason, access to the beach from the backshore was carefully controlled to prevent the development of blow-outs.

The sand section was cleaned by May 2 and work then began on the north shore. The contractor agreed to remove material only from the upper part of the intertidal zone. An agreement was reached to replace one section where the road on the backshore was left open to wave action by the removal of the boulders.

## Effectiveness and Results of Operation

- a) The clean-up of the sand area was a successful operation but by piling the material above water level, mixing of oiled and clean material took place. Although spillage by the bull-dozer was high it was of a much lower magnitude than on a cobble or mixed beach. Ideally the material should have been lifted off the beach rather than pushed around for removal by other equipment.
- b) Controlled access to the beach by the trucks using only one track resulted in damage to only one area of vegetation rather than wide-spread disturbance. Where contaminated material was found near the edge of the berm vegetation, the bulldozer was used to scrape down, rather than dig or push up the material. This again helped to minimize damage to the vegetated zone, (the destruction of the vegetation could have led to blow-outs).
- c) At the extremities of the sand beach the bulldozer often sank above its tracks in the silt and mud. This led to a thorough mixing of contaminated and clean material. In "soft" areas like this, very little can be done with heavy machinery.
- d) The sand areas were easily cleaned as the oil was only on the surface and this type of material is more easily handled than gravel or cobbles. An angle-blade on the bulldozer may have been an improvement, as this would have reduced spillage.
- e) Although the sand beach was clean after the contractor had completed his work, it was reoiled on various occasions subsequently from alongshore areas. By the end of July the intertidal zone was again covered with cakes of oil, sand and seaweed over large areas and these were still evident in November. The heavily polluted lagoon at the west end of the Cove released some of its oil throughout the summer and this was transported alongshore towards the beach by incoming waves (Section 5.4).

# 6.4.3 Indian Cove

#### (Fox Island Main)

This 850 foot beach was contracted for \$2,000 and involved four days work between May 15 and May 20, 1970. A Caterpiller 950, wheeled, front-end loader removed 1,368 cubic

yards of beach material during this period. The volume of sediment excavated on subsequent occasions is not available.

Description of contract area

This is a small concave pocket beach set back between two rock headlands. The beach has a maximum width of 130 feet and the material varies from coarse sand in the lower parts of the intertidal zone to cobbles at the crest of the storm ridge, the typical mixed beach of this area, with the dominant size in the cobble range.

The well developed storm ridge in the eastern half is backed by a low vegetated swale and a brackish pond. In the west half there is a continuous level area of old beach material fronted by a poorly defined beach crest ridge. The oil was confined to an area above normal high water level as a 6 to 12 inch thick caked layer approximately 10 feet wide which extended for almost the entire length of the beach (cf Photograph 6).

#### Restoration

A small experiment using a front-end loader with a 3 cubic yard capacity bucket was carried out on May 13 to test the effectiveness of this equipment for lifting off the oiled carpet. The encouraging results led to the use of the front-end loader for this operation and the beach was cleaned between May 15 and 20.

By May 25, the eastern 300 feet had been badly reoiled in the high water mark zone from the adjacent oiled rock areas (Photograph 13). On June 11 to 13, a bulldozer with a front bucket was rented to re-clean the beach. Oil was removed from the rock areas manually and limestone was spread over the rocks to stabilize the remaining oil, but recontamination occurred on several occasions throughout June and July (Section 5.4).

Beach profiles

To provide some measure of the effects of sediment removal, the beach was surveyed by 17 levelled profiles on May 14, 21, 29, June 14 and November 11. Six of these profiles, all from the east end of the beach, are presented in Figure 9.

In order to ascertain the distribution of sediment removal across the beach, profiles were taken immediately before (May 14) and after (May 21) the contract work. The limits of the thick, continuous layer of oil on the beach are shown (Figure 9b) and it is evident that the beach had been able to clean itself except for those areas above the reach of normal wave action. Most of the sediment removal was from this zone and as was noted earlier (Section 4.2), this part of the beach is not an area of accretion except during periods of storm waves or spring tides. Sediment replacement is, therefore, not rapid in this zone.



Photograph 13

INDIAN COVE

May 25, 1970

The east section of this pocket beach at low tide, four days after completion of the contract work. Before restoration, a thick continuous layer of oil, six to ten feet wide, covered the length of the beach above the high water line. (see photograph ll). The figure is standing near the high water line. Recontamination was from alongshore and the oil concentrated in the upper part of the intertidal zone.

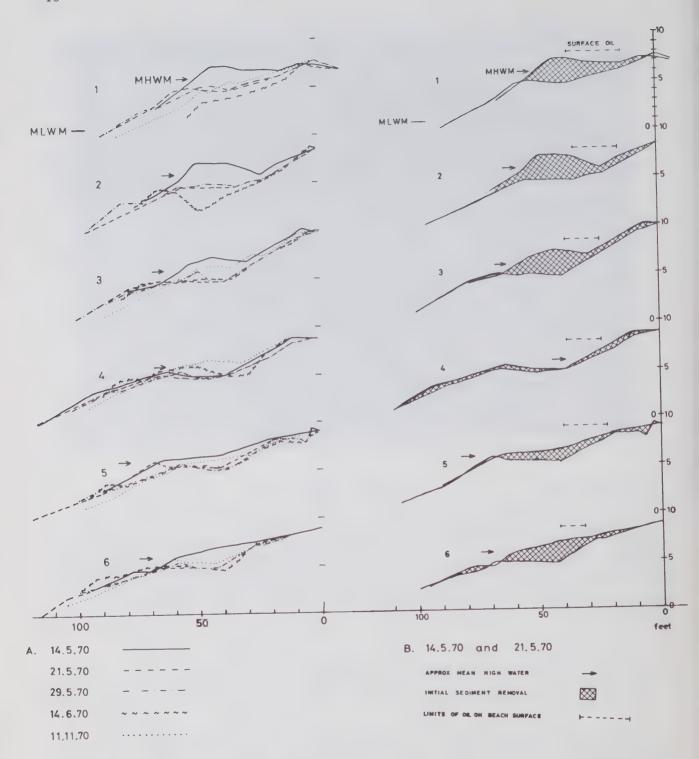
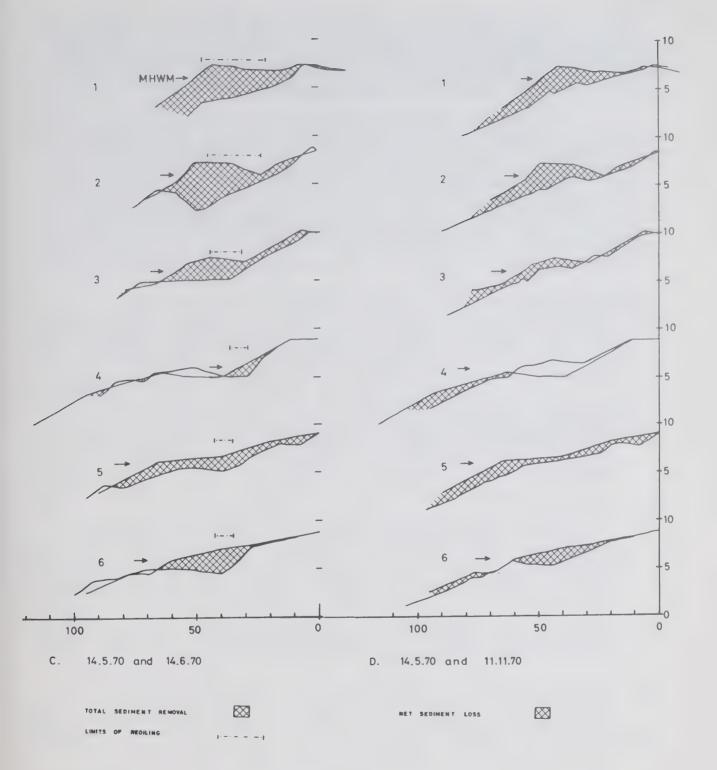


Figure 9. Beach profiles at Indian Cove.



As a result of a reoiling at the east end of the beach, further excavation took place and the effects of this are shown in Figure 9c, which compares the original situation with that following the second removal operation (June 14). Parts of this beach lost as much as five feet of sediments.

The fact that some replacement of material occurred during the summer months is indicated in Figure 9d. However, the profiles show that in the intertidal zone material had been removed by wave action. The beach had been unaffected by storm waves before November 11 and this should have been a period of net accretion in the intertidal zone. It is likely that the material was pushed up into the area around high water mark but the effect of winter storms combing down the beach may lead to a further lowering of the foreshore. Reprofiling during the spring of 1971 will give a better indication of any long term effects.

## Effectiveness and results of operation

- a) A carefully handled front-end loader was able to remove the one foot deep oiled layer efficiently. The bucket was used to cut under the oil and lift it, disturbing little of the uncontaminated sediments. This worked well as long as the operator did not attempt to fill the bucket, in which case the effect was to push rather than lift and spillage was very high. From the point of view of the contractor, this method is wasteful as the loader bucket is filled to only about one quarter of its capacity. Thus the efficiency of his operation suffers.
- b) Beaches where the oil is only on the surface are relatively easy to clean mechanically, as it is necessary to remove little sediment. However, with the equipment available, spillage is high unless the operator is particularly careful. A machine which could get under the oil carpet, in the manner of a fork lift truck, would be ideally suited to this type of situation.
- c) This beach was generally well cleaned except for occasional patches of oil contaminated material which were spilled by the loader; these were removed manually.
- d) The mixing of oil covered pebbles or cobbles with fine material leads to a surface cover of sand and silt sticking to the stone, which disquises the oil beneath. Subsequent washing and reworking by wave action tends to remove the surface material and may give the impression that the beach has been recontaminated although in fact this is a result of the mixing of material during cleaning.
- e) Reoiling of the eastern section took place (within 4 days) during a period of spring tides as oil was released from pools and cracks in the adjacent rock areas (Photograph 13). High temperatures during the preceding week had encouraged

the oil to collect in hollows and this was then easily washed out by the waves. The actual timing of the reoiling coincided with the spring tides so that the pools of oil had not been within the range of normal wave action for several weeks (Section 5.4).

f) The oil which recontaminated the beach was generally less than half an inch thick but would cause as much inconvenience as the thick layer removed earlier. Unless all sources of recontamination are dealt with, the chance of reoiling will remain. It is not practical to clean only one section of a coast, as oil from alongshore may spread to cleaned areas, which was demonstrated by the fact that this beach was reoiled several times.

## 6.4.4 Half Island Cove

This 1,500 foot section of beach was restored during a six day period between May 26 and June 15 at a cost of \$3,000. A Caterpiller 950 wheeled, front-end loader was used and 1,761 cubic yards of material were removed.

Description of the contract area

This is a wide, shingle beach partially set back from the general trend of the straight coast on the south shore of the Bay. The beach has a maximum width of about 80 feet and the lower intertidal zone is composed of fines and gravel while the storm ridge and upper beach zone are shingle. The alongshore movement of material appears to be from east to west.

The most easterly 600 feet of the contract area, commencing from a low rock platform, is backed by a 20 foot active till cliff. The next 500 feet form a ridge some 120 feet wide which has a lagoon in the rear. The central 200 feet were excluded from the contract. This area consists of a bedrock, boulder and shingle zone backed by a 5 foot active till cliff. The westerly 400 feet form a steep narrow beach with a shingle storm ridge, and a low wooded backshore. At one point, a rib of rock extends from the mid-tide zone seawards for almost 100 feet.

No oil patches were visible on the surface as all the contaminated material had been reworked by wave action and in parts had been buried to a depth of three feet. The rock platform just to the east of the contract section was badly oiled and much of this oil had collected in pools and cracks.

### Restoration

The beach was surveyed with 16 levelled profiles on May 25, prior to cleaning which commenced on the 26th. A short experiment with a road grader was conducted on May 26 to the west of the contract section on a wide, low sloping part of the beach.

The contractor left the area on the 28th but as the work was not completed satisfactorily he returned on June 13 and 14.

On June 14 and 15 material was removed from one section of the beach below high water mark. This was used as clean replacement for Phillips Harbour which has been subjected to machinery. The profiles were resurveyed on May 30 and June 15.

Results and effectiveness of operation

- a) The tests with the road grader were not productive as this machine requires a firm, flat or low angle surface. While it was able to work with some efficiency along the flat crest of the ridge, difficulties were encountered on the low beach face slope. As has been reported elsewhere (Section 5.3), a grader is useful and effective on firm, low sand beaches where the oil is on the surface. On a gravel or shingle beach it is of no practical value.
- b) The basic problem in this contract area was to remove buried, contaminated material. At the first attempt the contractor merely scraped over the surface and as this was not acceptable, he returned and removed more beach material. There still remained a great deal of contaminated sediments after the contractor was released (Photograph 14). The only way in which the method employed could have succeeded would have been to remove most of the upper sections of the beach without spillage. This is both undesirable and impractical and may have led to adverse effects, such as increased erosion of the till cliffs or breaching of the lagoonal ridge.
- c) Even if the beach were cleaned adequately, there is a danger in this area of recontamination by oil from the adjacent rock platform and other areas alongshore.
- d) This beach was not paralyzed by the oil and wave processes had been active and effective in cleaning and burying much of the contaminated material. Following the restoration programme the beach was no cleaner in terms of recreational purposes, so that there was little net gain from the operation. There was no large-scale removal of sediment and no damage to the beach itself.

# 6.4.5 Hadleyville No. 1

This 4,500 foot beach to the west of Oyster Point was contracted out for \$9,450 and work took 8 days between June 2 and 11. A Caterpiller D6C bulldozer and a Caterpillar 950, wheeled, front-end loader removed 3,980 cubic yards of material.



Photograph 14

HALF ISLAND COVE

June 3, 1970

Close-up of contaminated cobbles after restoration. This may have resulted from spillage, mixing, or the exposure of buried oiled material. Scale is graduated in decimeters.

## Description of contract area

This well developed, prograding, steep shingle beach is one of several similar formations on the north-west coast of the Bay. The longshore movement of material is to the east and large depositional ridges are evident at Oyster Point, just east of the contract area. Parts of the backshore have large well-preserved former storm ridges which attest to the gradual seaward progradation of the shoreline.

The intertidal zone has a maximum width in the order of 100 feet in the central parts of this section and the material distribution is fines in the lower intertidal zone with a shingle storm ridge above normal high water mark.

The shallow nature of the offshore area has allowed the beach to become oriented towards the dominant direction of refracted wave approach, that is, the south-east.

There was very little contaminated material visible on the surface of this beach. The oiled sediments had been reworked by wave action and were buried to a maximum depth of 4 feet. In some areas a layer of oiled material on the upper parts of the beach had been buried and the seaward edge eroded, so that a band of contaminated sediments was evident in the beach face slope (Photograph 15). The only method of cleaning with machinery involved the large-scale removal of mixed clean and oiled material.

#### Restoration

The contractor began work on June 2 and in many sections of the beach the bulldozer was used to pile up, above high water mark, material which was subsequently removed by the front-end loader. Where the oil was deeply buried, the loader made a cut down the beach and then excavated along the beach parallel to the ridge. The work was completed on June 11.

### Beach profiles

Ten profiles were surveyed across the beach on May 21, June 13, July 12 and November 11, 1970. Five of these profiles, from the central and east sections, are presented in Figure 10.

The "before and after" situation is given in Figure 10b and it is evident that material was removed largely from those areas above mean high water mark. In some places over three feet of sediments were removed in order to excavate the buried oil. Profile 6 is an example of the removal of the storm ridge crest which is well above the limits of normal wave action.

One month after completion of the contract (Figure 10c) very little sediment replacement had taken place even though the rate of accretion should be at its maximum at this time of year. This can be attributed to the fact that the intertidal beach was relatively undisturbed by the machinery and that there was no wave action on the higher parts of the beach during this period.



Photograph 15

HADLEYVILLE No. 1 May 30, 1970

This was one of the worst parts of this beach, in the west section at low tide. Contaminated material has been buried and subsequently exposed as layers in the beach face slope. This well developed beach is backed by a series of vegetated swales and former ridges.

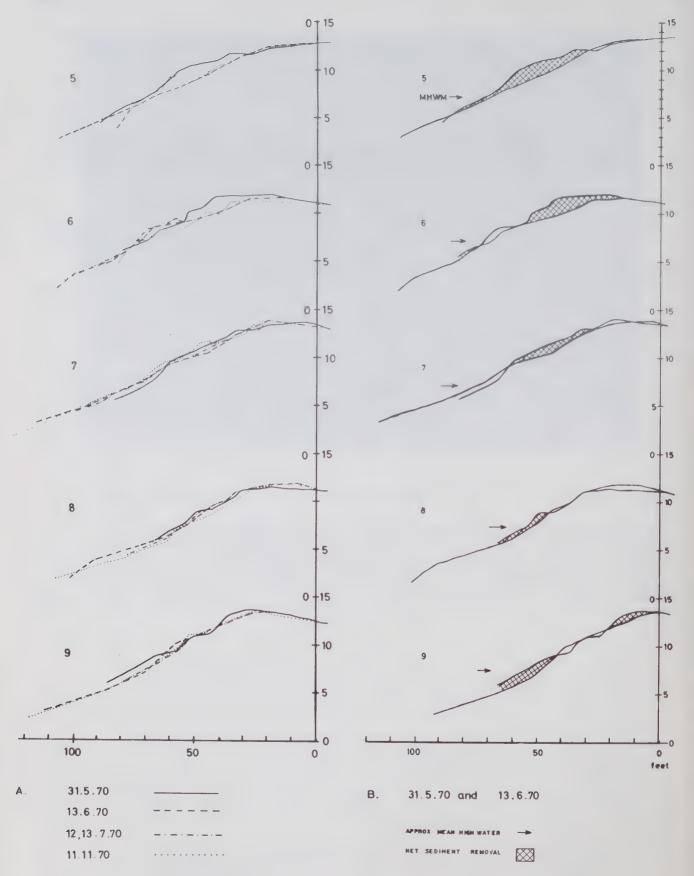
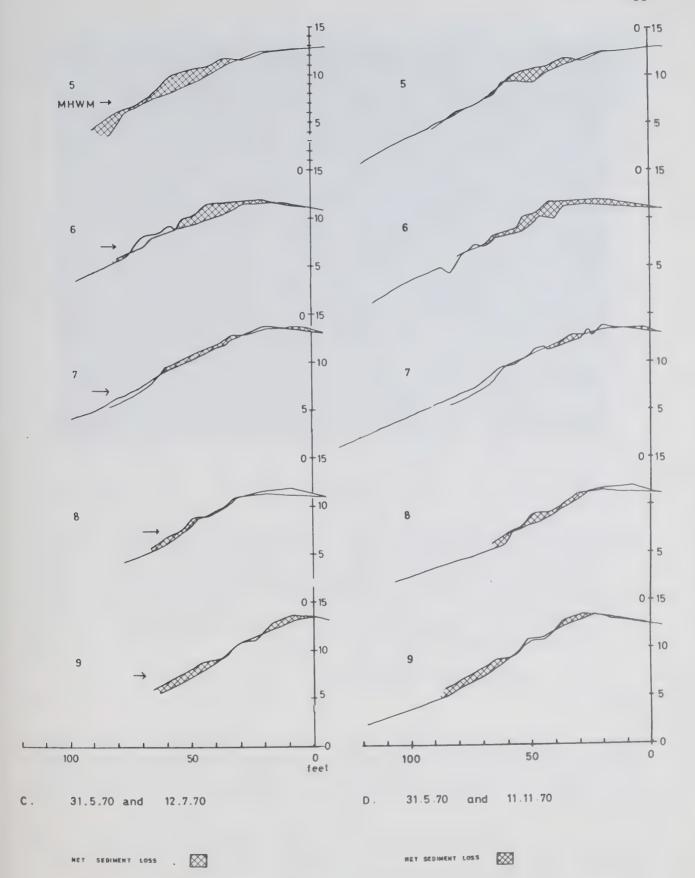


Figure 10. Beach profiles at Hadleyville No. 1.



Few changes took place in the fall and the November profiles (Figure 10d) record little significant erosion or accretion except for the presence of some small ridges near the high water mark. As most of the excavation took place above the intertidal zone, the profile changes on the lower part of the beach are probably due to normal variations in response to the littoral processes.

## Effectiveness and results of operation

- a) In certain sections where the contaminated material was buried, removal of large volumes of sediments is very inefficient. Of the total amounts excavated, only a small percentage of the material was contaminated (Section 5.2). There is no adequate method for mechanical cleaning of beaches where oil is buried.
- b) Contaminated material was still evident after the contract had been completed. The oil on the cobbles acquired a surface coating of fines which were easily knocked off or washed off by wave action.
- c) Unlike Arichat, Blackduck Cove or Indian Cove this beach was not paralyzed and normal beach processes were active. Wave action will succeed in cleaning the beach where machinery can only be used to remove large portions of the beach. (This situation is similar to Half Island Cove.)
- d) The use of a bulldozer to pile up material for removal mixes clean and oiled material. This does not contribute positively to the clean-up operation.
- e) It is unlikely that removal of sediments would have any long-term adverse effects on this beach provided large volumes are not removed from any one location. This should not detract from a critical analysis of the operation which was basically an inefficient way of dealing with the problem. Mechanical methods are not suited to this type of cleaning. Natural processes are more effective and less costly.

## 6.4.6 Other Contract Beaches

Moose Bay

This large cobble beach and spit complex on the west coast of the Bay is one of the most impressive coastal features of the area (Photograph 16) (Johnson, 1925). The west half of this beach was restored manually (Photograph 2) while the eastern 5,000 feet were put out to contract for machine work.

The problems which were faced in this area were similar to those of Half Island Cove and Hadleyville No. 1 as the oiled material was deeply buried and in order to remove it, several feet of sediments were excavated.



Photograph 16

MOOSE BAY

May 21, 1970

Air view of the east end of Moose Bay at mid tide. The bedrock island of Ragged Head, lower right, is connected to the mainland by a double tombolo. This is the western of the two arms and it gives way to a wide, long beach. The eastern limit of the area which was later restored is indicated by an arrow.

The initial attempt at manual cleaning was not continued to the end of the designated section because of the expense, which was double that for machinery on a cost per foot basis. The manual method is less harmful to the beach as little uncontaminated sediment is excavated. Although this approach does not lead to the removal of all the oiled material, mechanical methods are little better as there is no suitable equipment available.

## Hadleyville No. 2

This beach is distinguished from Hadleyville No. 1 as it is a separate coastal unit to the east of the area discussed above (Section 6.4.5). This beach area is little different from Moose Bay or Hadleyville No. 1 but is included as an example of the types of contamination which were experienced in the operation.

At the west end of this cobble beach oil had been reworked, buried, and subsequently exposed (Photograph 4). In the central area, a carpet of thick oil up to 15 feet wide was left above the limits of normal wave action (Photograph 6). At the east end of the beach cobbles gave way to coarse sand and here the oil remained as large discrete pans above the high water line (Photograph 3).

This area is a valuable example of how to deal with different situations using the methods available.

- a) The type of beach contamination in the west section represents a situation which is best left alone unless urgent requirements necessitate action. The waves had removed most of the surface oil by reworking and burial and the beach would continue to clean itself as this is an exposed beach which receives the full force of storm and swell waves from the east. To remove the contaminated material completely would require deep excavation which is harmful to the beach equilibrium and at present would be an inefficient method as no equipment is available which will remove all oiled sediments from cobble beaches.
- b) The thick cover of oil above the high water mark could be removed by a carefully operated front-end loader which has been shown to be effective in this type of situation (Section 6.4.3). No attempt should be made to remove the other sediments which have been reworked or buried by the waves for the reasons outlined above. The action taken would merely be to remove the surface layer of oil which has effectively paralyzed that section of the beach.

c) The large pans of oil on the surface of the sand can be removed manually, as these are easy to pick up with a shovel. The use of men rather than machinery is recommended, as the spillage from the equipment could be ground and mixed with clean sand. Any oil which is buried will be reworked by waves as it is exposed during the winter months. Should some of this be left on the beach the following summer, pickers could again restore the contaminated areas.

## Eddy Point

This foreland is made up of two cobble beaches which enclose a fresh-water lagoon on a shallow offshore platform (Photograph 17). According to local inhabitants the north-west shore is retreating at a rate of nearly one foot per year and as the upper part of this beach was covered with a continuous layer of oil, it was decided that this would be removed and then replaced with clean rip rap to serve as protection against any possible erosion resulting from sediment removal.

Profiles 1, 3, and 5 in Figure 11b were surveyed across this beach and show the effect of the restoration programme. Along the south-east limb of the foreland, the profiles 7, 9, 11, and 13 in Figure 11b show how the removal of buried oil 1ed to excavation in the zone above high water mark, the area where beach replenishment is least active.

Subsequent surveys in July and November (Figures 11c and 11d) indicate that no significant replacement had taken place on the north-facing beach and by November the beach had lost what little sediment had built up immediately following restoration. Profile 7 was levelled across the point of the foreland, which is the most mobile section of this feature, and the loss of material is probably the result of normal processes rather than the contract work. On the steep south-facing beach some build up had taken place along the high water line but this has not been sufficient to return the beach to its previous level.

## Deep Cove

Following the initial work on this beach (Section 6.1) further removal of contaminated sediments was carried out in June on the seaward side of the bar and in the lagoon. An angle-blade bulldozer was used in the lagoon area to remove the surface oil but the vehicle tracks acted to grind in the oil which had been buried. Although the machine was relatively efficient in scraping the surface, it did not affect the oil which was interbedded with the silt and mud except to mix it to a depth of several feet.



Photograph 17

EDDY POINT

June 12, 1970

Air view looking north. The lagoon is fresh water and rarely breached by the sea. The contract limits are given by the arrows.

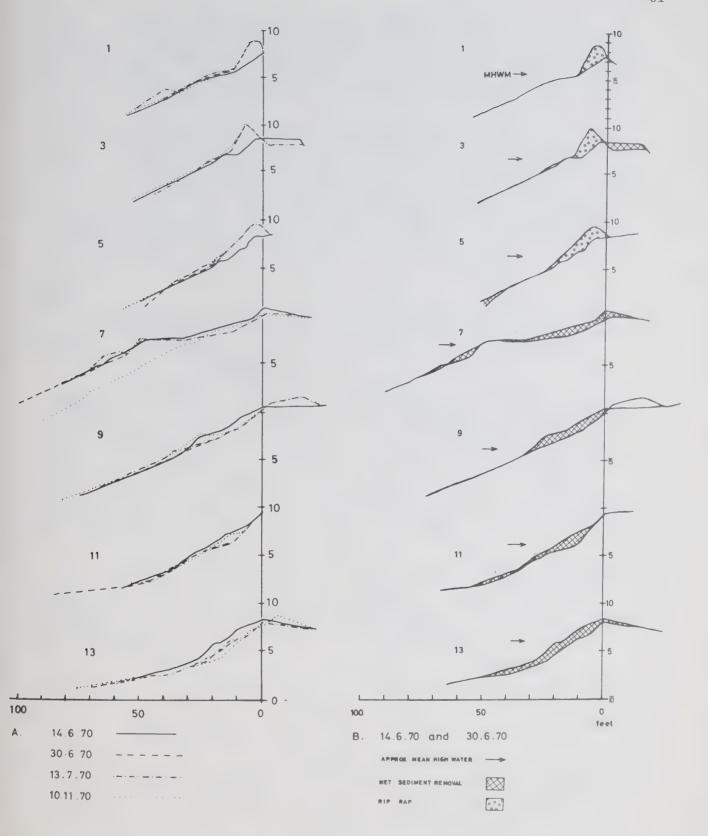
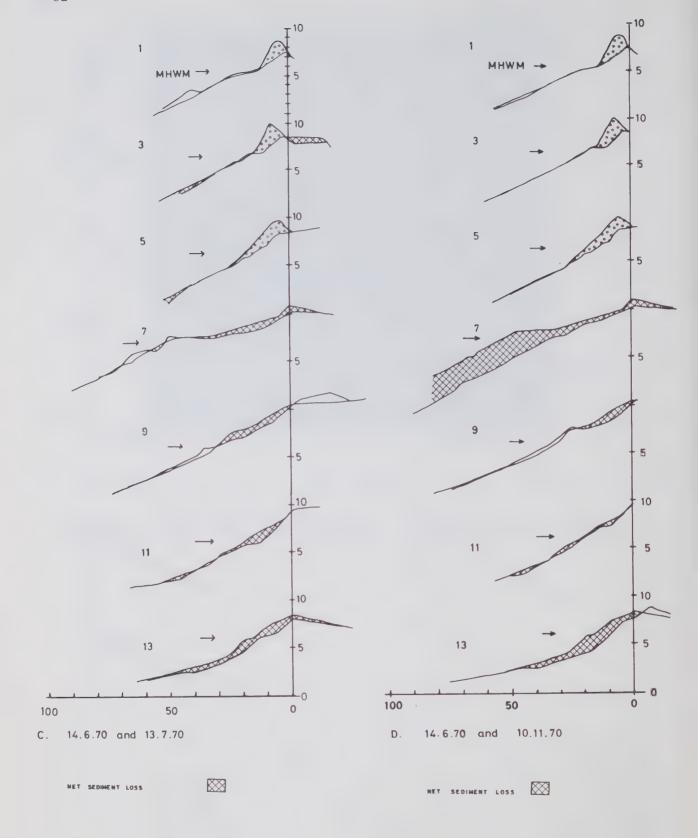


Figure 11. Beach profiles at Eddy Point.



The bar beach was restored again and material was replaced to prevent any major damage to this sensitive feature. Unfortunately the replacement sediments were taken from the spit which enclosed the lagoon and the cut into the spit which resulted will not be filled by natural processes as it is not in the active beach zone.

When this location was revisited in November it was evident that waves had washed over the road on the crest of the bar. Whether this would normally occur is not known but even if this is the case, there is clearly a serious danger that the bar could be breached. On such an unstable beach any interference with natural processes may easily accelerate erosion. Some rip rap was replaced in one section of the bar and this may help protect this part of the beach.

#### Walkerville

Several contracts were given to restore the paralyzed beaches in the sheltered areas of the north shore. One of the problems in this area was that oil was continuously moving alongshore from other contaminated areas.

This beach section near Walkerville on the north coast of Inhabitants Bay is in this sheltered area and has a very thin sediment cover which overlies a till platform and is backed by low till cliffs. The surface layer of cobbles was paralyzed by the oil and was removed by a bulldozer and a tracked front-end loader. This action exposed the fine, till-derived, sediment which overlay the bedrock and clean and contaminated material were mixed by the churning of the tracked vehicles (Photograph 18). With the removal of the surface layer of large material, the fines were removed by wave action.

In a situation like this it would have been preferable to rake over the paralyzed sections so that the beach material would be made mobile. Removal of sediment is harmful as the supply of cobble size material is very limited and erosion of the till platform and backshore may result from the loss of the protecting beach apron.

## 6.5 Summary

## 6.5.1 Sand Beaches

a) On lightly oiled sand beaches self-cleaning will take place. If restoration is necessary the manual removal of contaminated material with shovels and rakes is very efficient and effective, as this method does not involve the removal of large volumes of uncontaminated sediments (Section 6.3).



Photograph 18

WALKERVILLE

June 27, 1970

This beach on the sheltered north shore of Inhabitants Bay has had its thin surface layer of cobbles removed. This has exposed a wet layer of fines about three feet deep. These clean sediments have been thoroughly mixed with contaminated material by the tracked bulldozer and front-end loader.

- b) If the beach is heavily oiled and restoration is essential, a well-operated, wheeled, front-end loader which can cut under the oil is an efficient method of removal. The aim of the machine operator should always be to lift material in small quantities rather than push it forward and fill the bucket. There will, however, be spillage which can be removed manually. No beaches in this area were suited to the use of graders.
- c) Sand beaches can be cleaned for recreational purposes as long as there is no danger of recontamination (Section 6.4.2).
- d) When machinery is working in the beach zone it is important to prevent damage to the backshore vegetation as this could easily lead to blow-outs.

## 6.5.2 Cobble and Shingle Beaches

- a) Natural processes are effective in cleaning this type of beach unless wave action is limited and/or pollution is very heavy. Paralyzed beaches are by definition inactive and there is little self-cleaning at these locations (Section 5.2).
- b) It was found that cobble beaches could not be cleaned for recreational purposes with the methods available. Machinery alone could not remove all of the contaminated material.
- c) Where oil lies on the surface of a beach it can be removed efficiently by a well-handled, wheeled, front-end loader (Section 6.4.3). Even in these situations there is spillage from the equipment but this could be picked up manually.
- d) A lightly oiled or reworked beach with buried oil can only be restored by excavation methods. This often involves the removal of 3 to 5 feet of sediments and these are usually taken from those areas above high water mark where natural replacement is slowest (Section 6.4.5). If the storm ridge or beach crest is in any way lowered, attempts should be made to replace all the excavated material (Section 6.4.6).
- e) Local contractors have been removing beach material for construction purposes on several of the beaches in this area for a number of years. The beaches do not seem to have been damaged, as far as is known, but the long term effects are not understood, so that this activity should not be used as an argument for large-scale beach excavation. The lesson from Hallsands is an important one (Section 4.2).
- f) Where oil is buried in the beach it should only be removed if it is absolutely necessary. The low amounts of oil in these reworked sediments, often as low as 10 parts per million, mean that large volumes of material are removed for very little oil.

g) Mechanical restoration is best applied to the situations where oil is on the surface or where the beach zone is paralyzed. In the case of the latter, the object of restoration should be to remobilize the beach so that normal self-cleaning can be effective, rather than attempt to clean it by excavation (Section 6.4.1).

## 6.5.3 Recontamination

- a) No beaches should be cleaned if there is a danger of recontamination. This was evident in the instances where re-excavation was necessary (Section 6.4.3).
- b) It is not recommended that a contaminated shoreline be cleaned in sections as reoiling is likely unless the beaches can be protected or the adjacent oiled areas stabilized. Recontamination is usually light but is nevertheless undesirable.

## 6.5.4 Machinery

- a) Earth moving equipment is designed to excavate or remove large volumes of material and operators are trained to carry out this type of work. In order to use machinery for restoration purposes, time must be taken to train and supervise the operators so that they meet the requirements of this more delicate type of operation.
- b) Mechanical methods were found to be useful to remove surface oil from beaches and to remobilize paralyzed beaches. In the latter case, reworking of sediments rather than removal is more applicable. In particular, a well-handled, wheeled, front-end loader was found to be efficient to remove surface oil from cobble beaches (Section 6.4.3).
- c) Bulldozers were not satisfactory for restoration work, particularly the large 100 to 120 fly-wheel horse-power machines which were awkward to handle in confined areas. Spillage from the blade and grinding by the tracks were the major defects of this equipment, the same being true for all tracked vehicles.
- d) Material which is to be removed should be lifted directly from the beach. If it is piled up and later removed this doubles the chances of spillage and further mixing with clean sediments takes place.
- e) None of the machines were able to operate adequately on soft mud or silt areas. Even if traction was maintained this led to grinding and mixing to depths as great as four feet (Section 6.4.2).

## 6.5.5 Beach Profiles

- a) Accurately surveyed profiles across beaches are a valuable tool in establishing short and long term changes in the character of the beach as well as showing the immediate effects of excavation.
- b) The profile of a beach varies with each tidal cycle so that care must be exercised in isolating long term changes.

  Meaningful analysis requires a set of observations over a period of not less than a year. The results of the continuing survey programme on the Chedabucto Bay beaches will be reported at a later date (Owens, 1971b).

## 7. DISCUSSION

## 7.1 Selection of Beaches for Restoration After a Spill

In coastal areas affected by an oil spill, the decisions related to beach restoration should be based upon geological, wildlife and socio-economic factors. In the first instance, those shores which would require cleaning for tourist or other economic reasons should be established. The decision as to whether these beaches should be cleaned and the methods to be used would then be related to geological criteria. In certain instances, natural self-cleaning may be the acceptable solution; elsewhere the use of manual or mechanical methods may be required. Where cleaning is to be carried out, it is necessary to assess any possible adverse effects which this action could precipitate. For example:

- 1) removal of large volumes of material in an area of limited sediment supply could seriously affect the stability of the beach and backshore zones.
- 2) bars or lagoons could be breached by the disruption of the local shore environment as a result of sediment removal in particular localities.
- destruction of backshore vegetation, particularly in sand areas, could lead to blow-outs and aeolian erosion.

In all instances, the possibility of upsetting the delicate balance of the shoreline must be considered when assessing the socio-economic requirements, for in most instances, natural self-cleaning is more effective and is more desirable from the geological viewpoint. Certain shores such as rocky cliffs or shingle ridges are virtually uncleanable unless dispersents are employed, and in these areas piecemeal restoration of some beaches should be considered in the light of recontamination from the alongshore areas which would not be cleaned.

The system of rating which was used for the selection of beaches in this operation is one based solely on recreational capability (Anon., 1967). This was the only available information on the beaches of this area but was not a rating based on applicable criteria. Moose Bay has been given, for example, a "2" rating but is in fact used only by members of the small local community. The same rating was given to Point Michaud which is completely different geologically and very popular as a picnic and bathing beach. Distinction between different types of beaches for restoration is best defined by the actual recreational use, wildlife considerations, socio-economic factors, and upon

geological characteristics such as material size, sediment supply, sediment movement, and wave energy conditions. With this in mind the selection of beaches for restoration should be based on:

- a) Does the beach need to be restored?
- b) Can the work be done effectively without damaging shoreline stability?

These aspects of the restoration project could be discussed in a regional study (Section 7.2) and once the priorities have been outlined they would only require minor modification according to the actual seriousness of the contamination on given shores.

## 7.2 Regional Coastal Studies

The correct assessment of areas to be restored and the methods to be used requires an understanding of the coasts involved. To date, only a few small sections of Canada's coast, which is the longest in the world, have been investigated geologically and geomorphologically. Should future restoration projects be necessary, this means that at least reconnaissance surveys would be required before planning the operation. Such a survey should be able to provide a detailed outline of those sections of coast which would require restoration, on a priority basis, according to the character of the shore and economic, social, or wildlife requirements.

The shoreline of Chedabucto Bay had been reviewed very briefly by Johnson (1925) but until field investigations were carried out during this operation, this was the total available information on the coastal environment. Some basic research on the nature of Canada's coasts must be regarded as a necessity for future operations. This information would also greatly benefit tourist and conservation programmes.

The coasts of western Europe and the United States have been studied at least at the reconnaissance level so that restoration projects in these areas have benefitted from the existing storehouse of information. Studies of this nature could be carried out readily in Canada, at least for the areas where spills are likely, by the use of aerial photography, topographic and bathymetric maps, with additional field work for areas requiring more detailed study (Owens, 1971a). The investigation of Canada's shoreline has been neglected and this operation has brought to light the lack of even the most basic general studies in this field.

#### 7.3 Oil on Arctic Beaches

The arctic beaches are in a low energy environment where the intensity of wave and ice action varies with location and exposure (McCann and Owens, 1970). If oil were to pollute these coastal areas, the natural self-cleaning processes would act on a greatly reduced scale when compared to the exposed beaches of Chedabucto Bay. A more valid comparison would be the sheltered environments of the lagoon at Blackduck Cove or the beaches of Inhabitants Bay where wave action is very limited. In these areas it is expected that the oil will persist for several years before it is degraded by natural processes.

The lower energy conditions in the arctic mean that waves and wave-induced currents are relatively less effective in moving pebbles and cobbles when compared to lower latitude shorelines and this results in less transport and abrasion within the beach zone. Contaminated beaches would be cleaned only very slowly by this action and as there is less movement and redistribution it would be very easy for the littoral zone to become paralyzed.

If beach restoration is required in an arctic area it is important to consider that replenishment of a beach by natural processes takes place at a slower rate than in lower latitudes. Therefore, even if a beach appears to be building up, it may take a long time for any losses to be replaced. A documented example of the removal of beach sediments in an arctic location is given for the Point Barrow, Alaska area by Hume and Schalk (1964). It is reported that, as well as producing a shoreline retreat of 10 feet in 12 months, the lowering of the beach allowed ice floes to advance further inland and nearer shoreline installations. This mechanism of ice pushing up the beach is common in polar regions and such action would likely damage or destroy any fixed object in its path.

Although wave action is restricted in the arctic environment, ice may play an important role in reworking the contaminated sediments as broken ice is moved within the littoral zone (Owens and McCann, 1970). This would act to reorganize the oil and beach material but would do little to clean the beach. Biodegradation is particularly slow in this environment, so that the oil on a beach would be expected to remain much longer than in more temperate areas.

#### 7.4 Concluding Remarks

The coastal zone is an extremely complex environment, being at the interface of land, sea and air. The processes which operate within this zone are numerous and not fully understood. Any one segment of coast is the result of a series of interactions which involve a great number of variables and should any one of these variables be altered significantly, this may lead to a trend towards a different equilibrium situation which will satisfy the new process demands.

Although the coast is largely a response to the active processes of winds and waves, its form and nature closely affect these processes. Beach slope, material size, and sediment supply, for example, are critical factors in assessing longshore movement and beach nourishment. The construction of groynes, jetties and other harbour facilities has led, in the past to unexpected alterations of the adjacent coast due to the generation of new and the adjustment of existing processes, resulting in a different process-response environment. A great deal of research has now been carried out in this area and coastal engineering is regarded as a field of study in its own right. In general, work has been concerned with the construction of artifical shorelines, but outside of a few isolated examples little is known of the effects of sediment removal. In one documented case, referred to earlier (page 17 ) beach material was used for construction purposes and this was followed by a shoreline recession of 20 feet in 50 years. This example is in an area where there is no present-day supply of sediment. The beach, once deprived of sediment, could not be replaced and this "demonstrates most effectively how dangerous it is to tamper with a beach, how wrong it is to make assumptions about the drift of beach material without a full investigation and how important it is to study each part of the coast intensively and not apply general ideas too readily." (Steers, 1964).

Regional studies of Canada's coastline should be undertaken to determine the features and character of different shoreline units. A knowledge of our coasts is essential if adequate contingency plans are to be formulated. During restoration projects, geologists should be consulted to assist in the planning and operation of programmes in order to take into account local shoreline and beach conditions. In particular, the selection of beaches to be restored and the application of the various restoration methods should be carefully considered.

The role of the coastal geomorphologist or geologist is to make the planners aware of the possible effects of a beach restoration operation. Detailed analysis of a shore area requires time, particularly to determine seasonal as well as long term changes. For this reason, the preparatory work necessary for the provision of adequate information for decision making, should be carried out before a spill occurs. If this type of research is not undertaken, it will be possible to provide only a cursory assessment of the coastal environment based on a qualitative judgement. The function of the geologist in this type of operation is a very important one, but one which should be ahead of, as well as working within, the operation.

#### ACKNOWLEDGEMENTS

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Volume I - Report of the Task Force
Volume II - Report of the Scientific Coordination
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# MANUSCRIPT REPORT SERIES

No. 20

Zero padding as a means of improving definition of computed spectra

R.F. Henry and P.W.U. Graefe



Marine Sciences Branch
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# Manuscript Report Series No. 20

# ZERO PADDING AS A MEANS OF IMPROVING DEFINITION OF COMPUTED SPECTRA

R.F. Henry and P.W.U. Graefe

## Forward

The concept of zero-padding was first introduced in computation of correlation functions via the Fast Fourier Transform. However, its use in the contexts described in this report has not been discussed in any publication known to the authors.

## Abstract

A simple method of data modification is described which in conjunction with the Fast Fourier Transform provides economically a detailed picture of the frequency content of a sampled time signal and thus improves the accuracy with which the positions of spectral peaks can be determined. Also discussed is the use of zero-padding in the accurate computation of Fourier transforms of functions which are non-zero over a finite interval.



# Improved Spectral Peak Determination

This note describes a method for facilitating the accurate determination of the frequencies of periodic components in oceanographic records. In theory, such frequencies can be found to any specified degree of accuracy by identification of peaks in spectra computed from sufficiently long records, but in practice, the available records may be relatively short. Record length may be limited by simple logistic considerations, such as availability of ship time. Other constraints on record length which the writers have encountered are

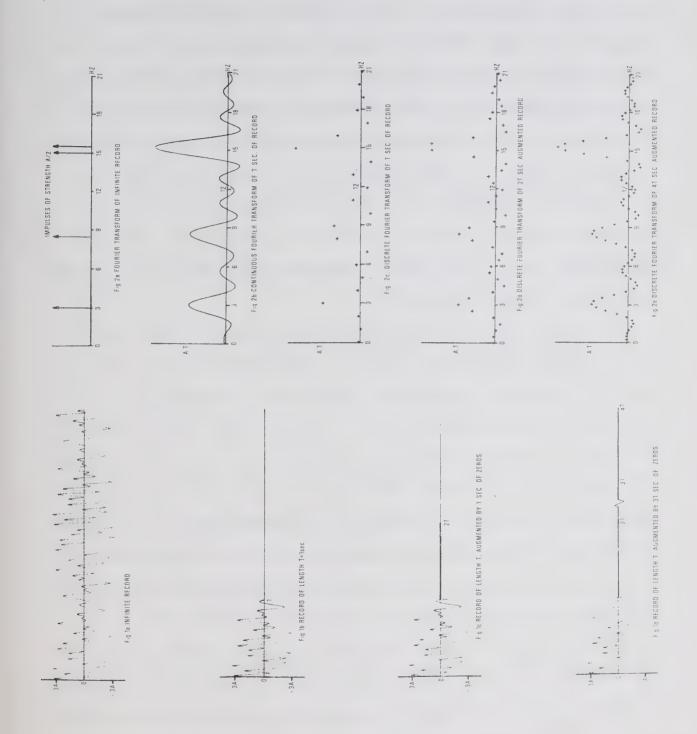
- a) in computer simulations of certain problems in fluids (for example, resonance studies of coastal bays and inlets), accumulation of round-off error acts as 'noise' which swamps the computed record if the computation is continued for too long;
- b) in determining components which have periods of the order of 10 years or more, practically all the reliable quantitative records available are relatively short;
- c) in calculating local values of period for phenomena with slowly-varying frequency, records are divided into short lengths in each of which the period is assumed to be approximately constant.

When discrete spectra are calculated from relatively short time records, accurate determination of peaks presents two difficulties. Firstly, the frequency resolution may not be sufficient to separate components whose frequencies differ

only slightly; nothing but use of longer records can rectify this problem. The second difficulty, the solution of which is discussed here, is that points in the computed spectrum are normally too widely spaced to outline peaks adequately.

As an illustration, consider the time record shown in Fig. 1(a), which consists of four periodic components of equal amplitudes and of frequencies 3.0, 8.5, 15.0 and 15.5 Hz: the spectrum of this signal is shown in Fig. 2(a). Suppose that in practice only a portion of this record, such as that of length T in Fig. 1(b), is available for analysis. The best approximation to the true spectrum (Fig. 2(a)) which can be obtained from this truncated record is shown in Fig. 2(b). This represents the convolution of the true spectrum with the Fourier transform of a rectangular function of length T through which, in effect, we view the original record (Jenkins and Watts, 1969). It will be noted that the frequencies at which peaks occur in the 'smeared' spectrum, Fig. 2(b), agree well with those of the periodic components shown in Fig. 2(a) except for the 15.0 and 15.5 Hz components which lie too close together to be separated with the resolution ( $\Delta f = 1/T = 1 \text{ Hz.}$ ) attainable with record Fig. 1(b).

Since  $\Delta f = 1/T$  is the finest spacing at which independent spectral estimates can be formed from a record of length T, practically all conventional algorithms available for computing spectra from sampled data are so arranged as to give values of the spectrum at points  $\Delta f$  apart. Consequently, when the record Fig. 1(b) is sampled, the discrete spectrum normally computed is that shown in Fig. 2(c).



Since the component periods are unknown, the coincidence of a sample point in Fig. 2(c) with a peak in Fig. 2(b), as happens at 3 Hz, is accidental and in general, peaks in Fig. 2(b) are often poorly defined by the widely spaced sample points in Fig. 2(c). For instance, it can be seen from the neighbouring sample points that the second peak in Fig. 2(b) lies in the interval 8.25 to 8.75 Hz but its precise location is hard to judge. However, the situation would obviously be improved if the number of samples could be increased, say for instance doubled as in Fig. 2(d), where the points are  $\Delta f/2 = 1/2T$  apart. Here the position of the second peak is far better defined.

The values needed for the improved plot, Fig. 2(d), can be found in two ways. Firstly, instead of using a standard Fourier transform algorithm to compute the N values in Fig. 2(c) according to the formula

(1) 
$$S(n.\Delta f) = \Delta t. \sum_{j=0}^{N-1} x(j.\Delta t) e^{-2\pi i.jn/N}, n=0,...,N-1$$

where  $i = \sqrt{-1}$  and  $x(j.\Delta t)$ , j = 0,...,N-1 are the sampled values taken from the record, Fig. 1(b), it is possible to write a special-purpose algorithm to calculate the 2N points in Fig. 2(d), which are given by

(2) 
$$S(m.\frac{\Delta f}{2}) = \Delta t. \sum_{j=0}^{N-1} x(j.\Delta t).e^{-2\pi i.jm/2N}, m=0,...,2N-1.$$

A simpler method of achieving the same result without special programming is to augment the record in Fig. 1(b) to that shown in Fig. 1(c) by adding an equal length of null

record, then sampling at intervals  $\Delta t$  as before and performing the <u>standard</u> discrete Fourier transformation on the 2N data points thus obtained. The spacing of points in the computed spectrum is governed by the augmented record length 2T and so, in this case, is equal to  $1/2T = \Delta f/2$  as required for Fig. 2(d). Of course, the samples  $x(j.\Delta t)$ ,  $j=N,\ldots,2N-1$  are all zero and consequently the standard Fourier algorithm in this case gives

$$S(m.\frac{\Delta f}{2}) = \Delta t. \sum_{j=0}^{2N-1} x(j.\Delta t) e^{-2\pi i j m/2N}$$
 (by analogy with equation (1))

(3) 
$$= \Delta t \cdot \sum_{j=0}^{N-1} x(j \cdot \Delta t) e^{-2\pi i j m/2N} + \Delta t \cdot \sum_{j=N}^{2N-1} x(j \cdot \Delta t) e^{-2\pi i j m/2N}$$

$$= \Delta t. \sum_{j=0}^{N-1} x(j.\Delta t) e^{-2\pi i j m/2N} + 0, m=0,...,2N-1$$

Thus, the values computed in this manner are identical with those given by (2). Essentially this second method of obtaining the points in Fig. 2(d) involves data modification, whereas the first method involves modifying the Fourier transform algorithm. Without question, modifying the data by simply adding zeroes is the easier technique in practice.

The process of augmentation can be continued indefinitely. For instance if the augmented record in Fig. 1(b) is augmented to double its length by the addition of a further length 2T of null record as shown in Fig. 1(d) and then sampled, the corresponding computed spectrum is as shown in Fig. 2(e). This has double the number of points in Fig. 2(c) and gives an even better representation of Fig. 2(b). Doubling the record

length at each stage is the most natural choice, if, as is normally the case, a Fast Fourier Transform with radix 2 is used to compute the spectrum.

When records are augmented by zeroes, the use of weighting, for example by means of a Hamming data window, is to be recommended. This does have the disadvantage of increasing the effective filter bandwidth, that is of broadening each peak in Fig. 2(b), but its main effect, reduction of the noticeable side-bands associated with the peaks, is particularly desirable with augmented records, since otherwise the increased number of spectral points shows up every side-band to full effect. The data window should be based on the original record length rather than on the length of the augmented record.

# Application to Finite-Length Functions

It is sometimes necessary to compute accurately the Fourier transforms of functions which are known to be zero outside a certain specified range. If the Fourier transform is computed in the most obvious manner by using samples of the function taken only in that part of the range where it is non-zero, the result may be seriously misleading. A good illustration of this difficulty occurs when it is necessary to obtain the Fourier transforms of so-called 'data-windows' (Jenkins and Watts, 1969), for instance, the rectangular or boxcar data window:

$$f(t) = 1,$$
  $|t| \le \frac{T}{2}$   
= 0  $|t| > \frac{T}{2}$ 

shown in Fig. 3(a). The continuous Fourier transform of this function is

$$F(\omega) = 2 \sin (\frac{1}{2} \omega T) / \omega$$

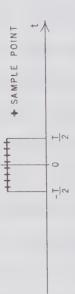
which is shown in Fig. 4(a).

If the Fourier transform of this window or some more complicated finite-range function has to be computed from sampled values rather than analytically, the simplistic approach would be to use only samples from within the range where the function is non-zero. Suppose for instance that the Fourier transform of the rectangular data window is computed with a conventional algorithm using the equi-distant samples shown in Fig. 3(b). The resulting discrete Fourier transform is shown in Fig. 4(b). There are two different but equivalent explanations of why Fig. 4(b) is such a poor representation of the true Fourier transform, Fig. 4(a). These are

a) the interval (1/T) between the computed points shown in Fig. 4(b) is determined by the length of record in the time domain, in this case T, on which the computed transform is based. Unfortunately, apart from the point at  $\omega=0$ , the computed spectral points coincide with zeroes



FIS. 3a : RECTANGULAR DATA WINDOW OF LENGTH T



3b: RECTANGULAR WINDOW 3a WITH EQUI-SPACED SAMPLES

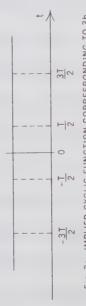


Fig. 3c : IMPLIED CYCLIC FUNCTION CORRESPONDING TO 3b

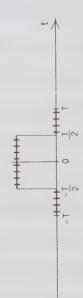


Fig. 3d : RECORD 3b AUGMENTED WITH EQUAL NUMBER OF ZEROES

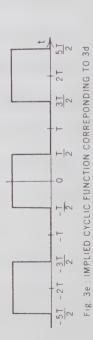




Fig. 4b : DISCRETE TRANSFORM OF UNAUGMENTED RECORD 3b



FIG. 4c : DISCRETE TRANSFORM OF AUGMENTED RECORD 3d



FIR. 4d : DISCRETE TRANSFORM OF RECORD AUGMENTED TO LENGTH 4T

- of the true spectrum  $F(\omega) = 2 \sin(\frac{1}{2} \omega T)/\omega$ .
- b) The fact that only a zero-frequency component is found in the computed transform (Fig. 4(b)) implies that the corresponding sampled function in the time domain is constant over -∞ < t > ∞. Thus it appears as if the original rectangular function of Fig. 3(a) is in effect repeated at intervals T all along the time axis as shown in Fig. 3(c). In fact this cyclic repetition of the original sampled function is an inherent property of discrete Fourier transforms. It may be noted that the repetition interval is equal to T, the length of record on which the computed transform is based.

A better representation of the true Fourier transform, Fig. 4(a), can be obtained by zero-padding. For instance, in Fig. 3(d), where the original samples are augmented by an equal number of zeroes, i.e. increasing the effective record length from T to 2T, the resulting computed spectrum is that shown in Fig. 4(c). Since the computed spectral values are now spaced 1/2T apart, a much better representation of the true spectrum is achieved. Fig. 4(d) shows additional improvement obtainable by further zero-padding, the case illustrated corresponding to an effective record length of 4T. This shows, in terms of interpretation a), how zero-padding can be used to improve the poor definition in Fig. 4(b).

In terms of interpretation b), the effect of zeropadding is to increase the interval at which the original function is repeated along the time axis. For instance, when the sampled rectangular window of Fig. 3(b) is padded with an equal number of zeroes as shown in Fig. 3(d), the corresponding cyclic time function implied under interpretation b) is that shown in Fig. 3(e). Since in the limit, when the number of zeroes added is increased indefinitely, the spacing between repetitions of the original function becomes infinite and hence the implied cyclic function is identical to the original function. It is reasonable to expect that even a lesser degree of zero-padding will result in some improvement over the unpadded case.

In conclusion, it can be seen that zero-padding essentially results in a legitimate form of interpolation between the coarsely-spaced spectral points obtained by straightforward computation of the discrete Fourier transform of the original sample points.

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# MANUSCRIPT REPORT SERIES

No. 21

On the water of the Canadian Arctic archipelago, an atlas presentation of 1961 and 1962 data

F.G. Barber and A. Huyer



Marine Sciences Branch

published for Environment Canada by

Department of Energy, Mines and Resources, Ottawa

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# Manuscript Report Series No. 21

# ON THE WATER OF THE CANADIAN ARCTIC ARCHIPELAGO, AN ATLAS PRESENTATION OF 1961 AND 1962 DATA

F.G. Barber and A. Huyer

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### ' 1. Introduction

This work represents a step in the development of those oceanographic data observed in and near the Canadian Arctic archipelago (Fig. 1) during the years 1960, 1961 and 1962 (Table 1); a presentation of part of the 1960 observations is in preparation. In the latter work it is emphasized that the data exhibit considerable variability, most of which is believed to be real. No change in this consideration has occurred and for this reason a separation by year for the presentation here has been maintained. It may be that with increased understanding a significant integration of such data will become possible.

In 1961 the distribution of stations occupied north of Foxe Basin and Hudson Strait\* was uneven (Fig. 2), but was much better than usual from Barrow Strait and Smith Sound to Davis Strait. On this account it was decided to limit the presentation here for the 1961 data to the latter area. Similar consideration of the 1962 station distribution (Fig. 3) led to a more widespread presentation which includes most of the data (Table 1) north of Foxe Basin and Hudson Strait. Other presentations of portions of the data of Table 1 include a heat budget for Barrow Strait (Huyer and Barber, 1970) and a fresh water budget for the Gulf of Boothia (Huyer and Barber, 1971). Ford and Hattersley—Smith (1965) utilized the data observed in Nansen Sound in 1962 in CCGS "John A. MacDonald", while Milne, Herlinveaux and Langleben (1962) and Milne (1963) reported on the 1961 ICE PACK observations and on that portion of the 1961 CCGS "Labrador" data in M'Clure Strait respectively.

Data from three of the 1962 "Labrador" stations in Smith Sound (consec numbers 1 to 3) were utilized by Collin (1965 p8) in a determination of the volume transport through Smith Sound and data at "John A. MacDonald" 1962 station (consec number 57) were presented by Palfrey (1968 p57).

More recently Muench (1971) in a detailed examination of northern Baffin Bay utilized most available data including those for 1960, 1961 and 1962.

#### 2. Method

The plotting sheet used (CHS 750) was derived from chart 7000 titled "Arctic Islands" and was provided in quantity by the Canadian Hydrographic Service. In the preparation of

<sup>\*</sup>Data obtained in 1961 in the Hudson Bay system have been reported (Anon., 1964a; b).

Table 1 A listing of the relevant data for 1960, 1961 and 1962 indicating the year obtained, name of ship, Cruise Reference Number and a literature reference to the data if available:

Year	Ship	CRN	Reference
1960	"Westwind"	00670	Anon., 1964c
1960	"Theta"	329	Anon., 1964d
1960	"Labrador"	340	do
1961	"Edisto"	00688	-
1961	(ICE PACK)	763	Herlinveaux, 1961
1961	(PCSP) *	339	out.
1961	"John A. MacDonald"	344	Anon., 1966b
1961	"Labrador"	341	do
1962	"Atka"	31966	-
1962	"John A. MacDonald"	359	Anon., 1966a
1962	"Labrador"	362	Anon., 1967a
1962	(ICE CAMP)	784	Herlinveaux, 1963
1962	"Salvelinus"	377	Anon., 1963a
1962	(land-based)	363,4,5,6,7	Anon., 1963b

<sup>\*</sup>Polar Continental Shelf Project

the sheets the data were placed beside the appropriate station in blue ink\* and the sheet subsequently became the final line drawing. The values were derived from individual station graphs of each variable against depth using available bathythermograph data and T-S diagrams where appropriate; interpolated data of the data records were not used. Occasionally, it occurred that some of the tabulated values of the data records did not fit the overall interpretation. When these appeared anomalous, an attempt was made to determine whether they were real or due to blunders. In only the two instances below was it possible to find definite errors:

CRN	Consec	Depth	Sal Temp Correction	Justification
362	<u>A</u>	20	0.36	bathythermogram shows positive temperatures above 50 m.
359	7	10	31.376 1.89	inadvertantly omitted from data record.

In some instances it seemed almost certain that a blunder had occurred and here the values were "corrected" or deleted. These are shown in Table 2, which includes a statement concerning the basis for the suspicion. It is realized that these alterations are largely a matter of our interpretation.

The contour interval is indicated on each distribution. A dashed contour represents one additional to the interval indicated there, while a dotted contour indicates that the configuration is in doubt, because either the interpretation or the observation is in doubt. Many of the 1962 stations were located within the archipelago and in contouring in this area it was assumed that there was continuity in the particular distribution in spite of the barriers posed by the islands. For this reason, many of the contours initially extended across these land areas and subsequently had to be removed.

### 3. Data

It appears that the 1961 salinity determinations for each of the three ships were carried out with conductivity bridges after completion of the field work. The "Edisto"

<sup>\*</sup>In many of the figures the values so placed are still visible.

Table 2. Suspected errors in the tabulated values of salinity and temperature in two of the data records for 1962.

CRN	Consec	Depth (m)	Correct Sal (0/00)	tion Temp (°C)	Justification
359	4	200 250	33.308 33.690		strong density inversion.
	13	0 10	29.690 30.130		strong density inversion.
	18	0	delete		strong density inversion; correction not obvious.
	33	350	delete		strong density inversion; correction not obvious.
	37	500 600	delete delete	delete delete	density inversion and somewhat low temperatures.
	39	200	delete	delete	slight inversion where strong gradient expected.
	41	10 20 30 50	29424 31020 31170 31411		a number of density inversions.
	58	0 10	26260 26700		strong density inversion.
362	6 13 14 17	600 450 378 10 30 223	31146 31176 delete	delete delete delete	
	46	145 194	33480 33553		density inversion.
	52	10	32600 32839		density inversion.
	53	300	32039	0.54	bathythermogram shows temp. at 275 m positive and increasing; "Atka" station shows 0.54 at almost the same depth and position:

samples were analysed "approximately two months after they had been collected", while those of "Labrador" and "John A. MacDonald" were not analysed completely until February and May of 1962 respectively. It appears likely that the determinations for the 1962 samples would have been carried out over a similar period of time.

The vertical distribution of salinity at about 1150~m and deeper (Fig. 4) observed in "Labrador" in 1961 and 1962 suggests the existence of a small systematic error of about 0.01~O/oo between the data of the two years. Consideration of the temperature data (Fig. 4) does not suggest that systematic differences occur in the data of the two surveys.

#### 4. Remarks

A feature of most of the shallower presentations is the strong gradient in each distribution across Baffin Bay. The feature appears to be associated with a movement of a relatively low-salinity cold water out of the archipelago and along the east coast of Baffin Island. This is a view of long-standing of course, which recently received some additional support in a note (Huyer and Barber, 1971) about the distribution of freshwater within the archipelago (after Tully, 1958). The interpretation was based on the oceanographic data for 1961 and 1962 and on the distribution of freshwater in relation to precipitation (Fig. 7) and runoff (Sanderson and Philips, 1967), particularly around the Gulf of Boothia.

Data on observed ice conditions during 1961 are available (Archibald et al., 1962) and Dunbar (1962 p5) remarked that, "In general the 1961 season was unfavourable for navigation in the Queen Elizabeth Islands". The 1962 season within the archipelago was relatively light (Anon., 1963c; Markham and Hill, 1963; Black, 1965; Lindsay, 1968) and may have contributed to higher than usual near-surface temperatures and greater storage of seasonal heat (Huyer and Barber, 1970). This situation appears to be reflected in those near-surface temperature data obtained in the "John A. MacDonald" (Fig. 13) from the Gulf of Boothia north to western Lancaster Sound, Barrow Strait, western Jones Sound and north to the vicinity of Tanquary Fiord.

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- Figure 3. The approximate positions of stations occupied in 1962 north of the latitude of Hudson Strait.

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  "Labrador" (Anon., 1967a) and CCGS "John A. MacDonald" (Anon., 1966a). (b) Stations occupied by five land-based parties (Anon., 1963b), from the pack ice (Herlinveaux, 1963) and on a survey by USS "Atka".
- Figure 4. Temperature and salinity values at depth below about 1150 metres in Baffin Bay from a number of sources. The circled salinity values at "Labrador" stations 70 and 71 in 1960 are believed to reflect sampling error.
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#### Part A

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- Figure 10. The distribution of depth and temperature on a  $\sigma_t$  surface during 1961 from data observed in "Labrador". (a) 27.50. (b) 27.60. (c) 27.65. (d) 27.70. (e) 27.72. (f) 27.74.
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#### Part B

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- Figure 13. The distribution of salinity and temperature during 1962 from data observed on "John A. MacDonald", "Labrador", "Atka", and from the Ice Camps and by the land-based parties. On the distributions at 200 m and greater the hatched lines show the approximate limits of areas where the greatest depth in any section of the channel is less than the depth of the distribution. (a) 30 m. (b) 50 m. (c) 75 m. (d) 100 m. (e) 150 m. (f) 200 m. (g) 250 m. (h) 300 m. (i) 400 m. (j) 500 m. (k) 600 m. (l) 800 m. (m) 1000 m.
- Figure 14. The distribution of depth and temperature in 1962 on the 34.40/00 surface. (a) Depth. (b) Temperature.

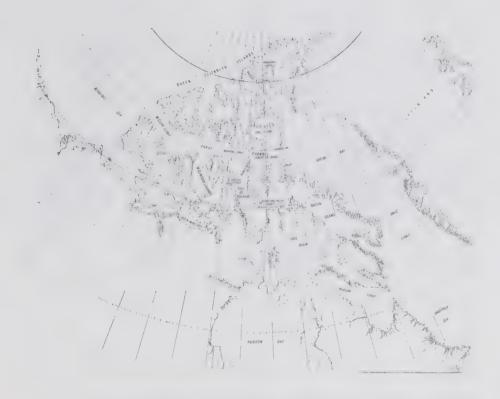


Figure 1. Some place names in the Arctic archipelago.



Figure 2. The approximate positions of stations occupied in 1961 north of Foxe Basin and Hudson Strait.

# Figure 3

The approximate positions of stations occupied in 1962 north of the latitude of Hudson Strait. (a) Stations occupied on surveys by CCGS "Labrador" (Anon., 1967a) and CCGS "John A. MacDonald" (Anon., 1966a). (b) Stations occupied by five land-based parties (Anon., 1963b), from the pack ice (Herlinveaux, 1963) and on a survey by USS "Atka".



3(a)



3 (b)

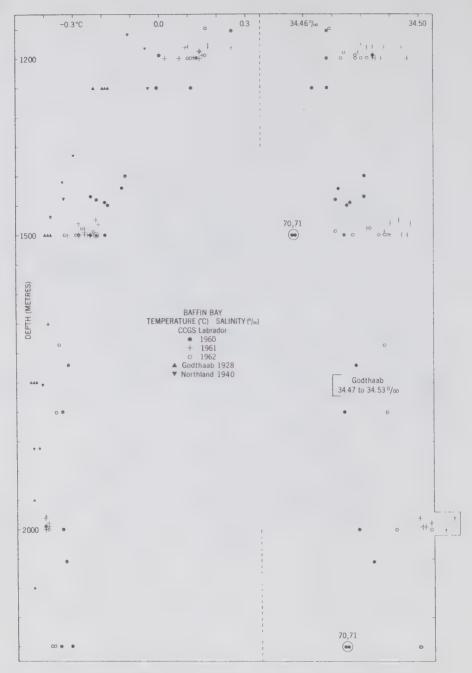


Figure 4. Temperature and salinity values at depth below about 1150 metres in Baffin Bay from a number of sources. The circled salinity values at "Labrador" stations 70 and 71 in 1960 are believed to reflect sampling error.

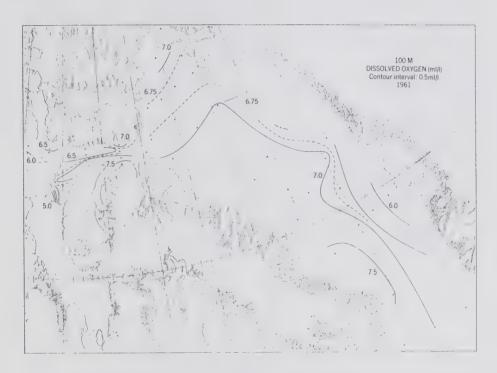
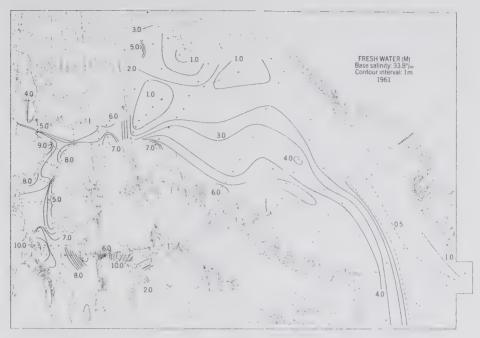
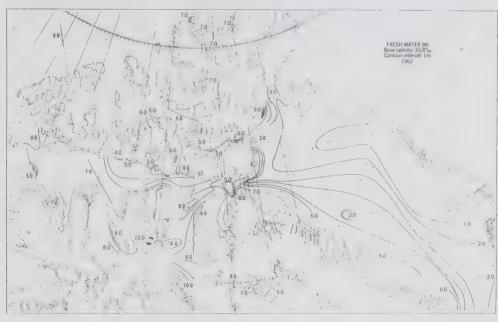


Figure 5. The distribution of dissolved oxygen at 100 m observed in "Labrador" in 1961. The relatively low value in the Gulf of Boothia (consec number 39) could be due to error.

The distribution of the amount of fresh water. (a) 1961. (b) 1962.



6(a)



6(b)

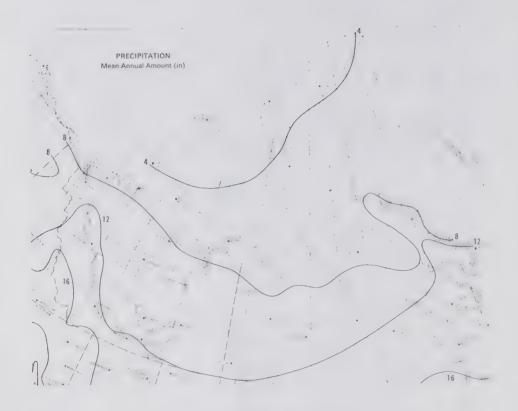


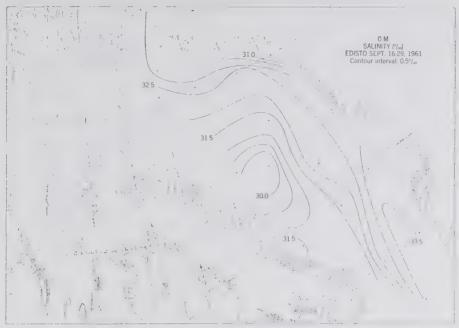
Figure 7. Precipitation in the Northwest Territories (Anon., 1967b).

PART A

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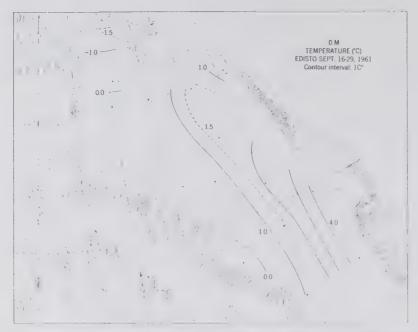
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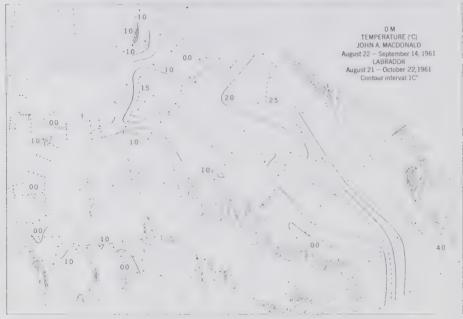
The distribution of salinity and temperature at the surface during 1961 from data observed in the vessels "Edisto", "John A. MacDonald" and "Labrador". (a) Salinity. (b) Temperature.





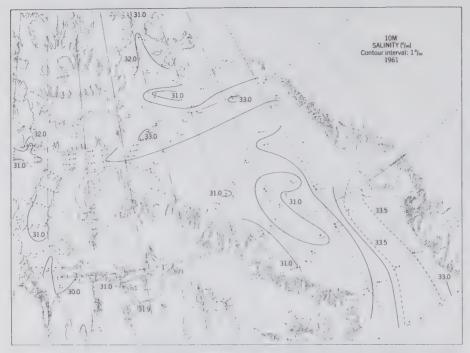
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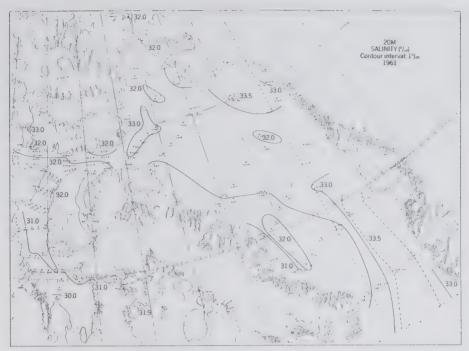
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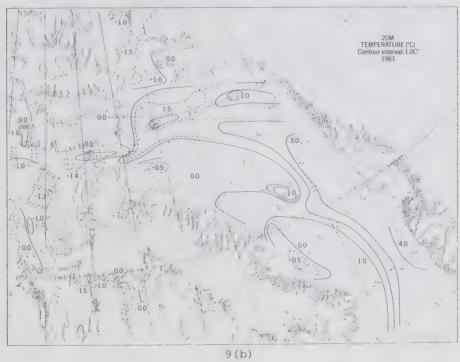
The distribution of salinity and temperature at standard depths during 1961 from data observed in the "Edisto", "John A. MacDonald" and the "Labrador". (a) 10 m. (b) 20 m. (c) 30 m. (d) 50 m. (e) 75 m. (f) 100 m. (g) 150 m. (h) 200 m. (i) 250 m. (j) 300 m. (k) 400 m. (l) 500 m. (m) 600 m. (n) 800 m. (o) 1000 m.





9(a)

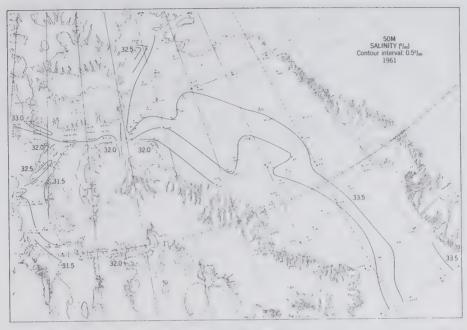


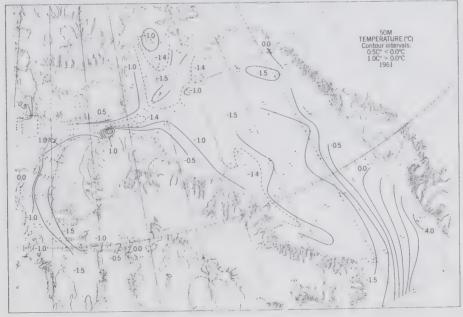




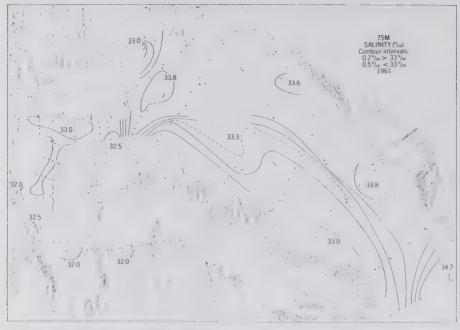


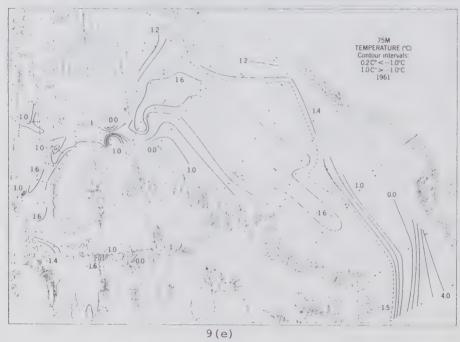
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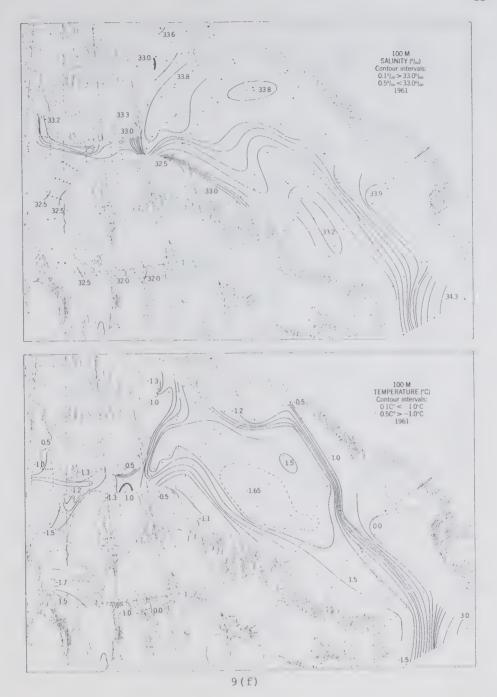




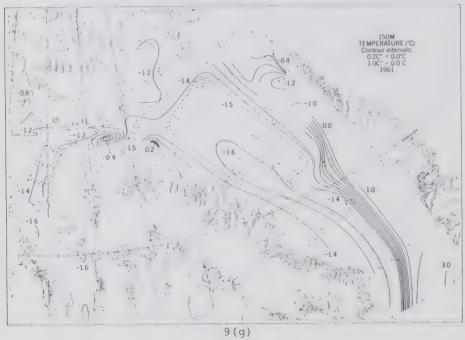
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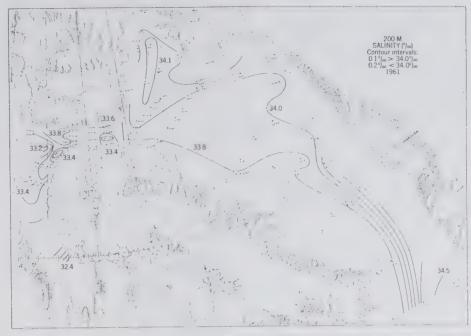


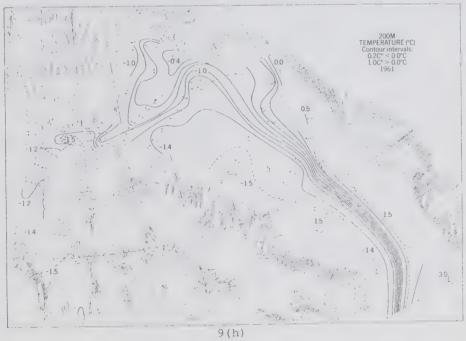


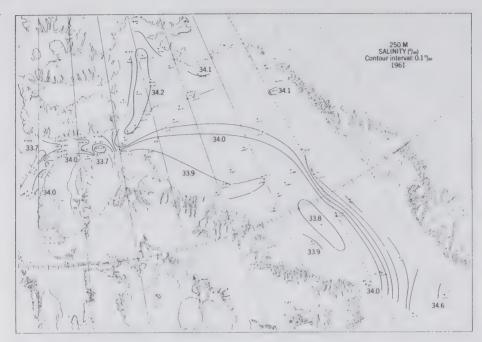


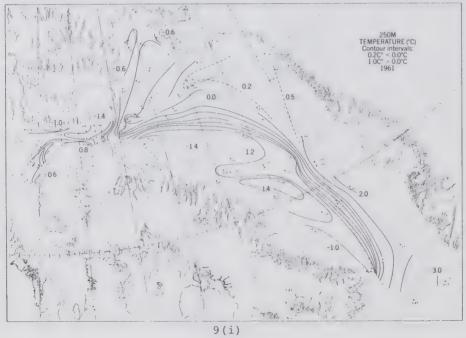


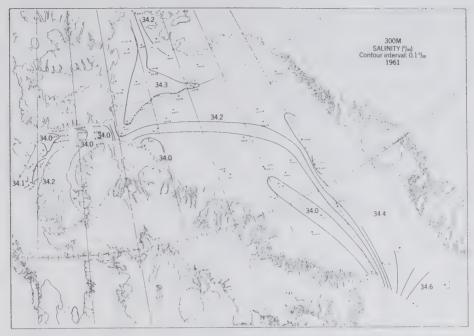


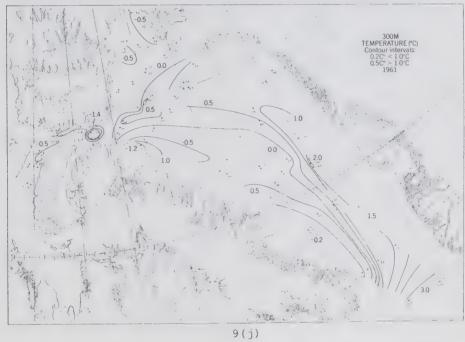


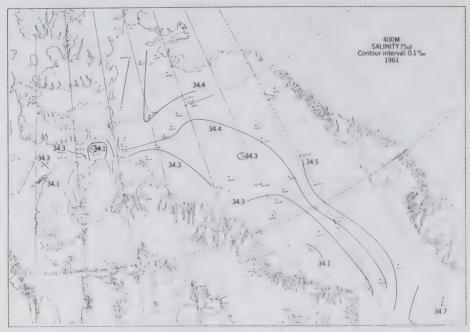






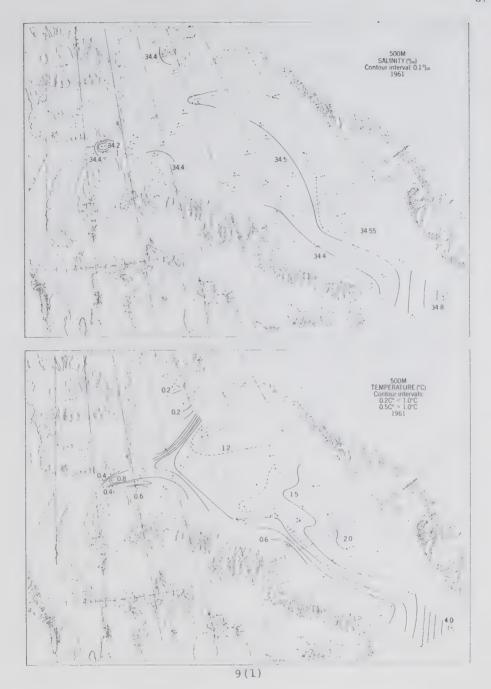


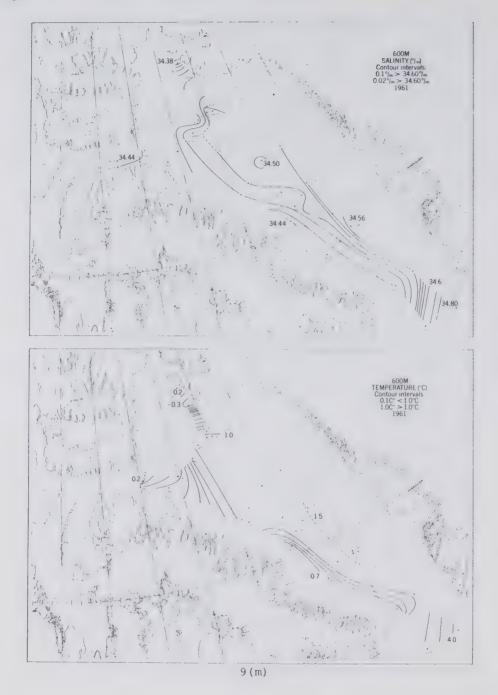


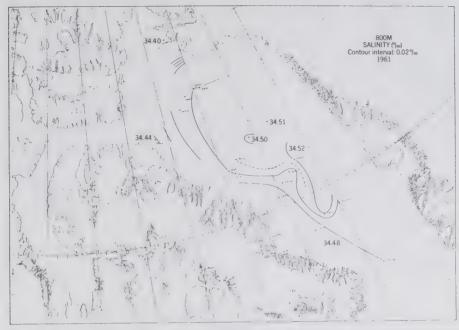




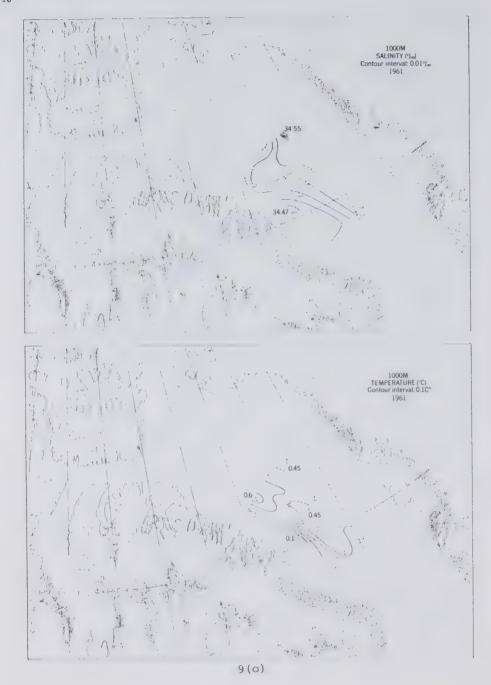
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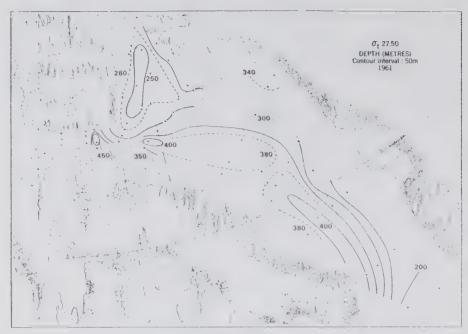


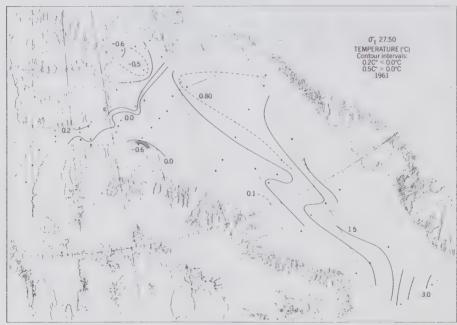




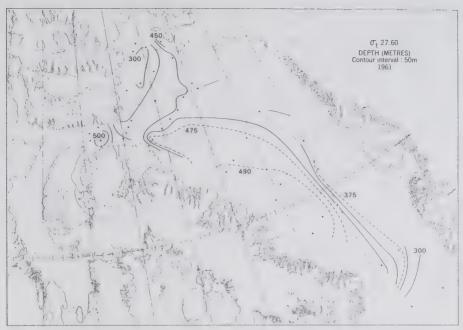


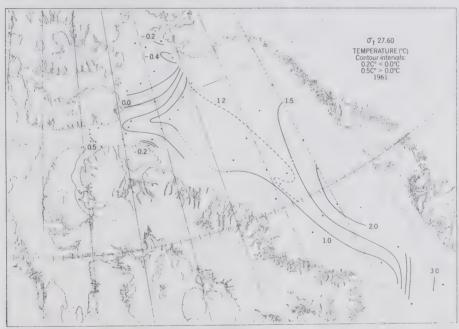
The distribution of depth and temperature on a  $\sigma_t$  surface during 1961 from data observed in "Labrador". (a) 27.50. (b) 27.60. (c) 27.65. (d) 27.70. (e) 27.72. (f) 27.74.



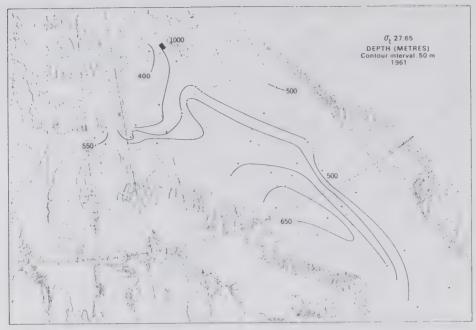


10(a)



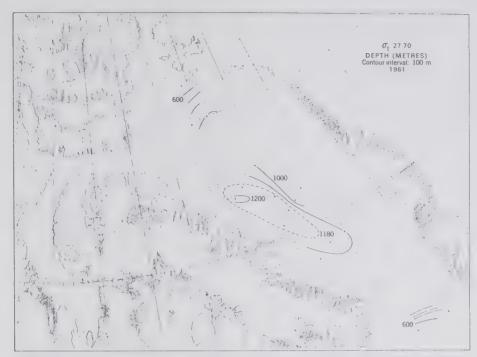


10(b)





10(c)









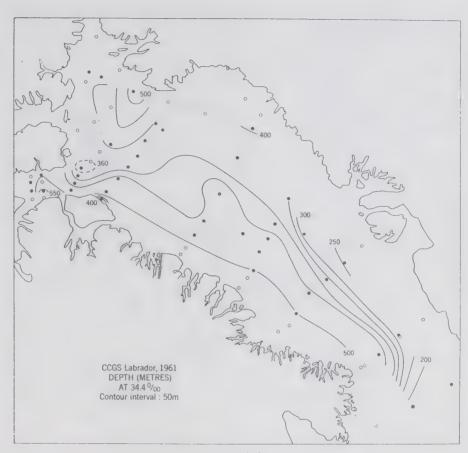
10(e)



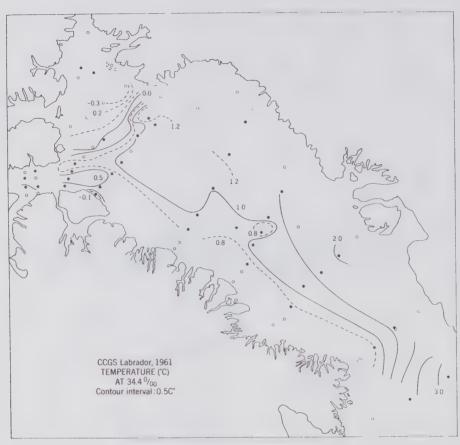


10(f)

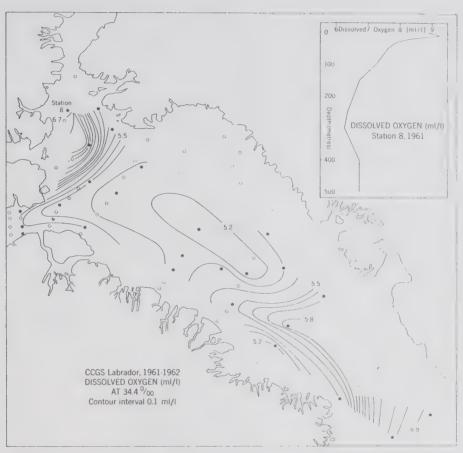
The distribution of depth, temperature and dissolved oxygen on the  $34.4^{\circ}/\circ$  os surface. (a) Depth. (b) Temperature. (c) Dissolved oxygen (includes 1962 data).



11(a)



11(b)



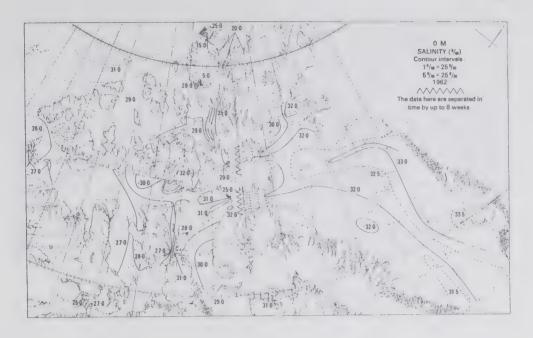
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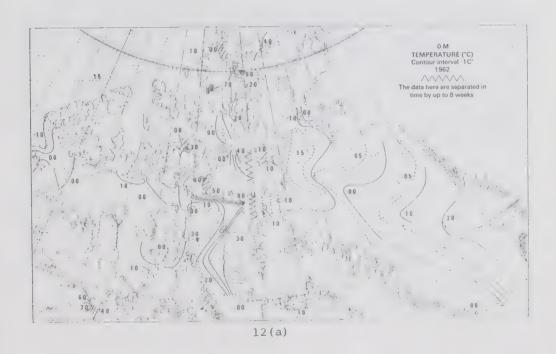


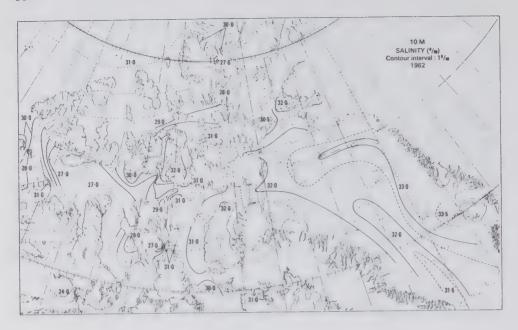
PART B

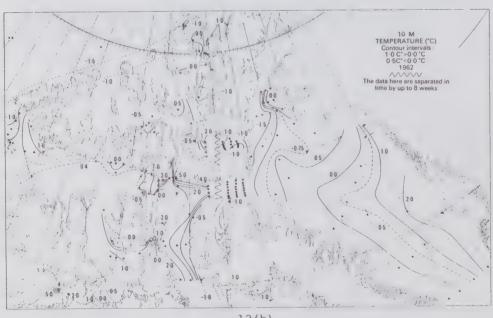
#### Figure 12

The distribution of salinity and temperature near the surface during 1962 from data observed from "John A. MacDonald", "Labrador", by the land-based parties and from the Ice Camps. (a) 0 m. (b) 10 m. (c) 20 m.

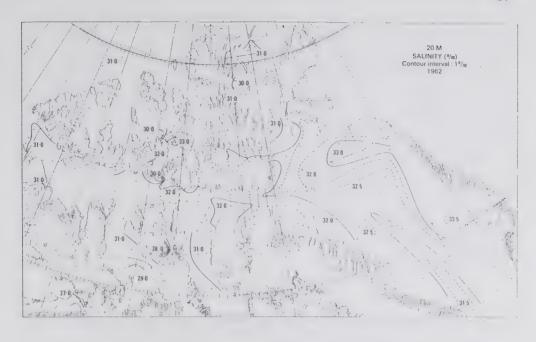


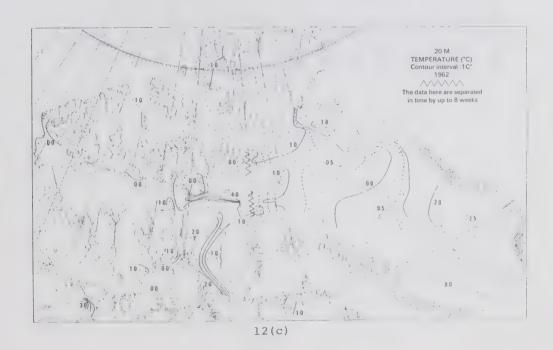






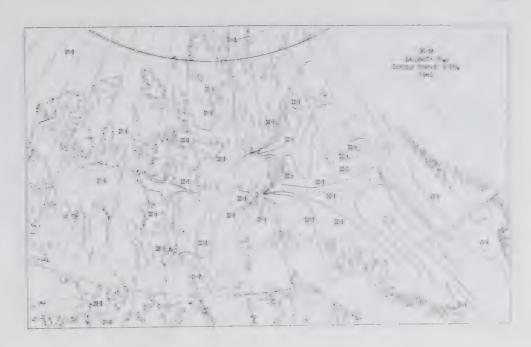
12(b)



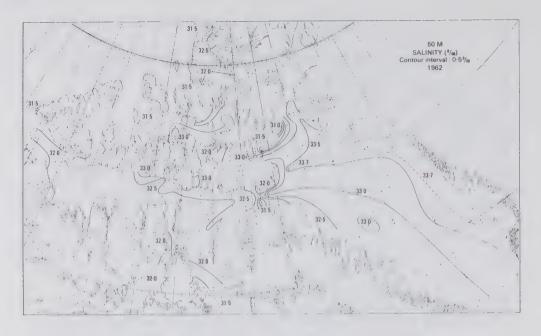


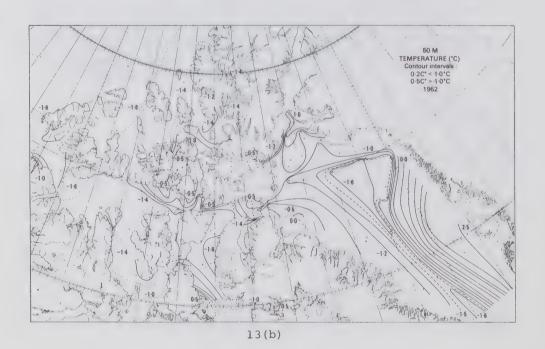
#### Figure 13

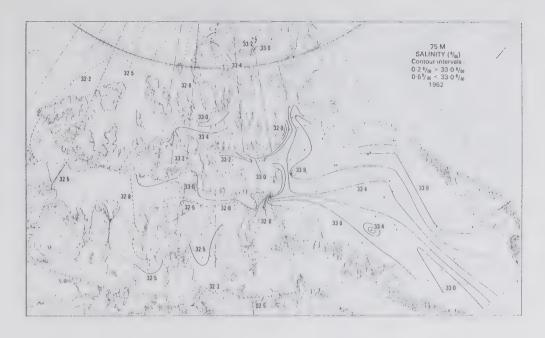
The distribution of salinity and temperature during 1962 from data observed on "John A. MacDonald", "Labrador", "Atka", and from the Ice Camps and by the land-based parties. On the distributions at 200 m and greater the hatched lines show the approximate limits of areas where the greatest depth in any section of the channel is less than the depth of the distribution. (a) 30 m. (b) 50 m. (c) 75 m. (d) 100 m. (e) 150 m. (f) 200 m. (g) 250 m. (h) 300 m. (i) 400 m. (j) 500 m. (k) 600 m. (l) 800 m. (m) 1000 m.

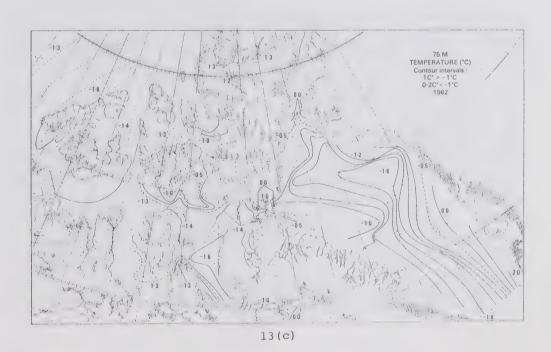


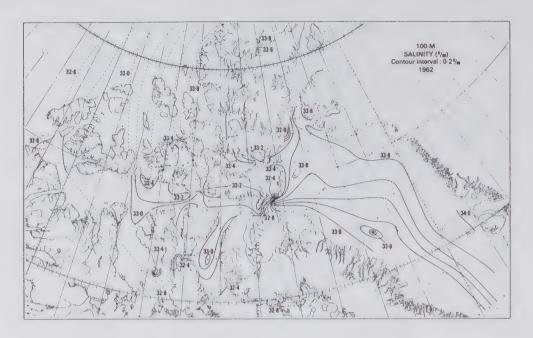


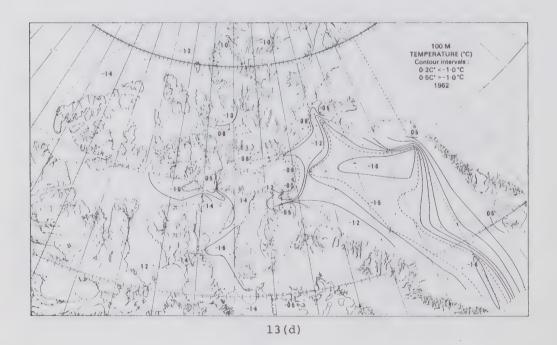


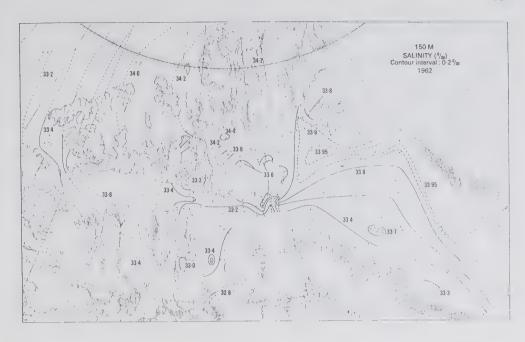


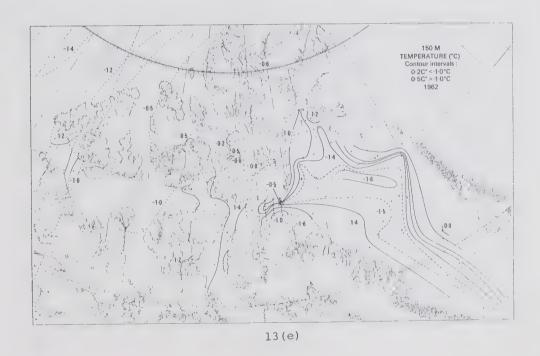


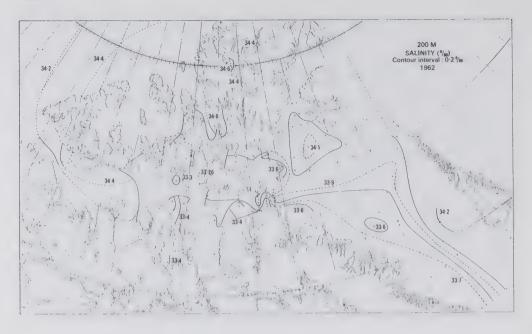


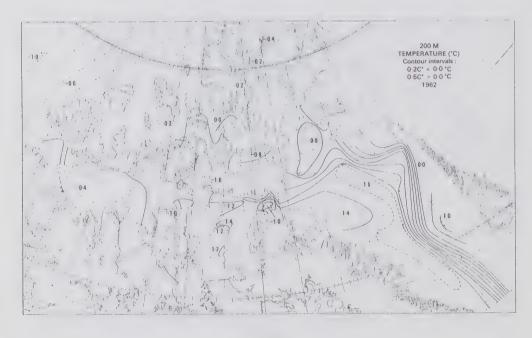




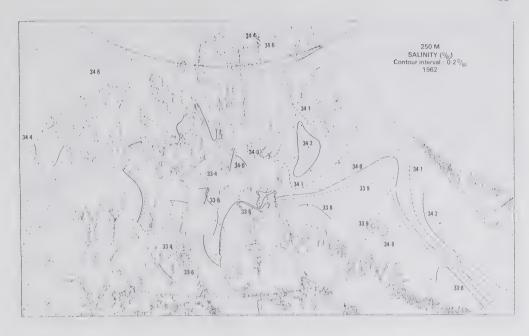


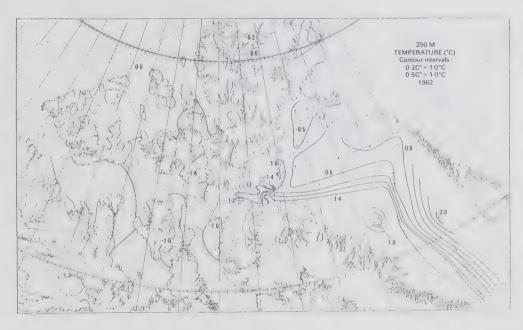




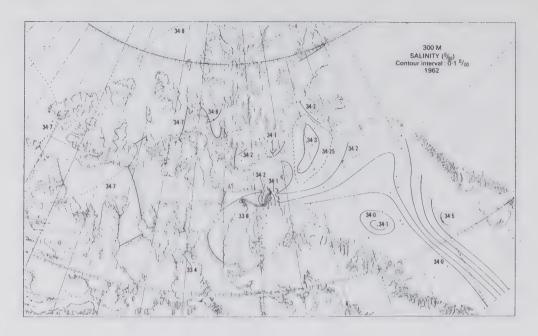


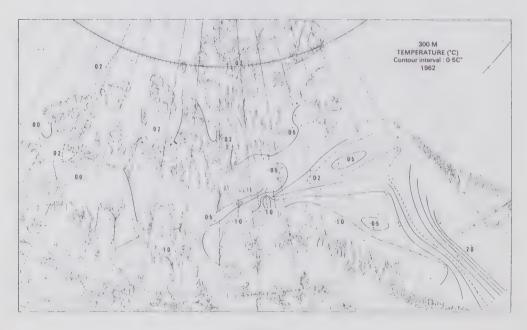
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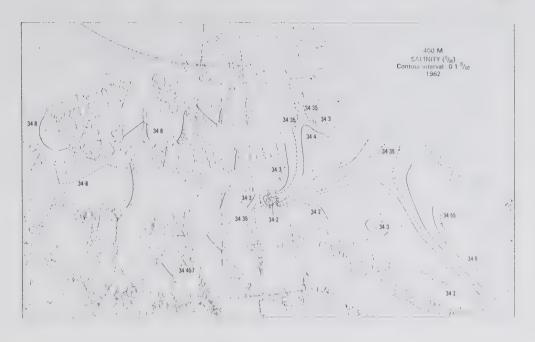


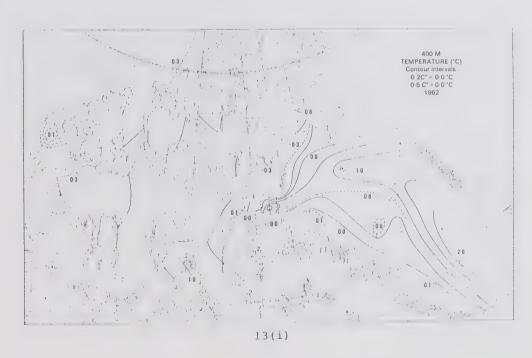
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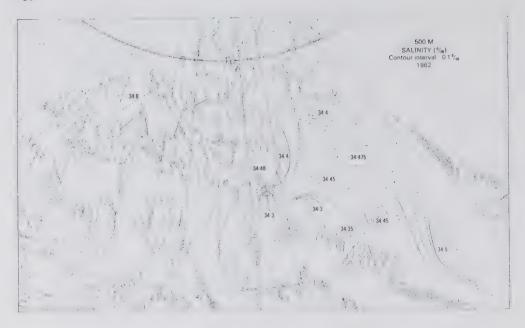


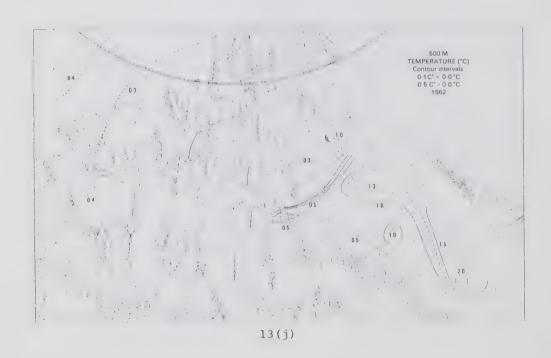


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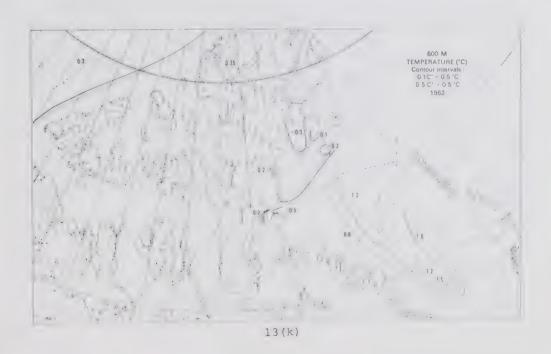


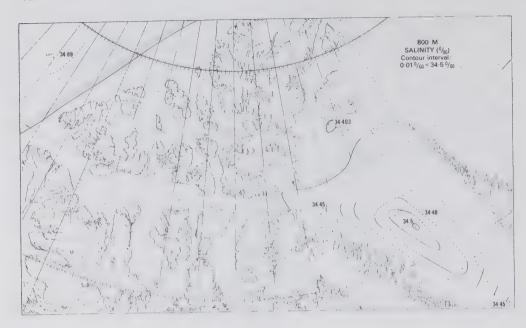


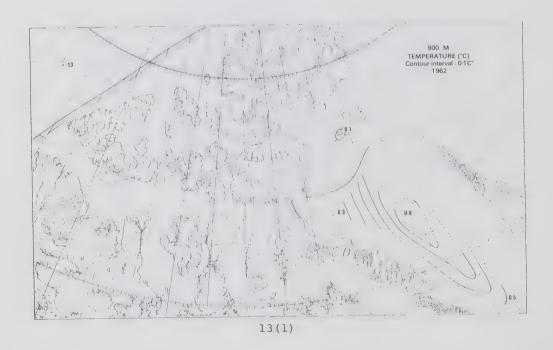


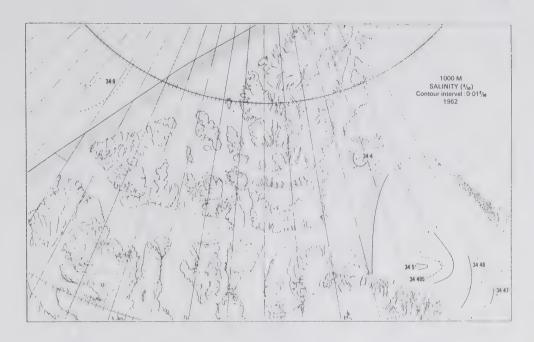


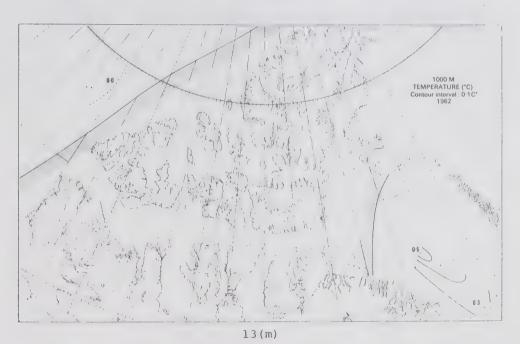






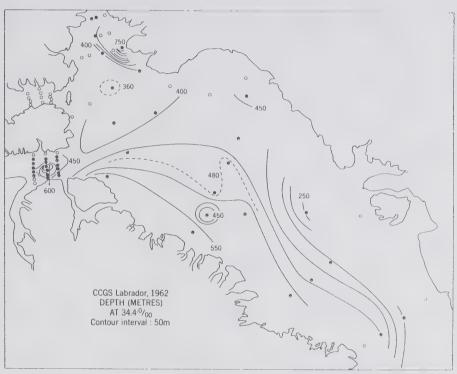




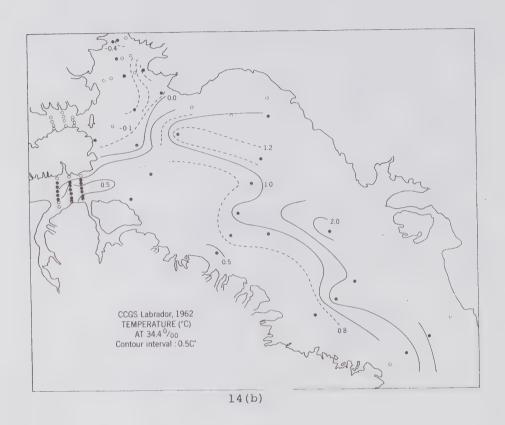


### Figure 14

The distribution of depth and temperature on the  $34.4^{\rm O}/{\rm oo}$  surface. (a) Depth. (b) Temperature.



14(a)



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No.22

Numerical values of the conical function  $K_p(x)$  for a range of values of order p and argument x and their zeros and bend points

T.S. Murty, P.J. Richards and J.D. Taylor



# Manuscript Report Series No.22

NUMERICAL VALUES OF THE CONICAL FUNCTION  $K_P(X)$  FOR A RANGE OF VALUES OF ORDER P AND ARGUMENT X

AND THEIR ZEROS AND BEND POINTS

T.S. Murty, P.J. Richards and J.D. Taylor



#### 0. Abstract

The conical functions  $K_p(\mathbf{x})$  arise in many problems, e.g., in calculations of the effect of rotation on gravity mode frequencies in laboratory basins. Since the numerical values of these functions are not tabulated to cover all the ranges of interest, these were computed and tabulated here.

#### 1. Introduction

The Legendre function  $P_{\nu}(x)$  becomes a conical function  $K_p(x)$  for  $\nu = \frac{-1}{2} + ip$  where p is real. Usually these are referred to as Mehler's conical functions with the argument  $x \ge 1$  (Hobson, 1931). Hobson gives the following infinite integral:

$$K_{p}(x) = \sqrt{\frac{2}{\pi}} \cdot \frac{\cosh(p \pi)}{\pi} \int_{0}^{\infty} \frac{\cos(pt) \cdot dt}{(x+\cosh t)^{1/2}}$$
 (1)

However this formula was found to suffer severely from rounding errors for p>6 even when double precision is used. Hence the procedure described below is used.

#### 2. Mathematical problem

For the range  $1 < x \le 10$  for the argument and  $0 \le p \le 10$  for the order we used the following formula (Erdelyi, 1953).

$$P_{v}(x) = \frac{1}{\pi} \int_{0}^{\pi} (x + \sqrt{x^{2} - 1} \cos t)^{v} dt$$
 (2)

But 
$$K_p(x) = P_{\frac{-1}{2}+ip}$$
 (3)

Define 
$$x + \sqrt{x^2-1} \cos t \equiv U$$
 (4)

Substituting (4) and (2) into (3) one gets

$$K_{p}(x) = \frac{1}{\pi} \int_{0}^{\pi} u^{-\frac{1}{2} + ip} dt$$

$$= \frac{1}{\pi} \int_{0}^{\pi} \frac{u^{ip}}{u^{1/2}} dt$$

$$= \frac{1}{\pi} \int_{0}^{\pi} \frac{\cos(p \cdot \log u)}{u^{1/2}} dt$$
(5)

The derivative of  $K_{p}(x)$  is given by

$$K_p'(x) = \frac{1}{\pi} \int_0^{\pi} \frac{d}{dx} \frac{\cos(p \cdot \log u)}{u^{1/2}} dt$$

After some algebra this becomes

$$K_{p}'(x) = \frac{-1}{\pi} \int_{0}^{\pi} \frac{1}{u^{3/2}} \left( \frac{\cos(p \cdot \log u)}{2} + p \cdot \sin(p \cdot \log u) \right) \left( 1 + \frac{x \cdot \cos t}{\sqrt{x^{2} - 1}} \right) dt$$
 (6)

#### 3. Numerical procedure

The integrands of the expressions for  $K_p(x)$  and  $K_p(x)$  are even functions of t at both t=0 and t= $\pi$ . Hence the end corrections for the trapezoidal integration rule will be identically zero, as these corrections consist of odd differences only. The simple trapezoidal rule was used here to evaluate both  $K_p(x)$  and  $K_p'(x)$ .

We found it necessary to use a different number of points for different ranges of the argument to obtain twelve

digits after the decimal point.

For x=1 the formula (6) for  $K_p'(x)$  has a pole and hence cannot be used. In this case for  $x \le 1+10^{-10}$  the following formulae which are derived from the hypergeometric formula for  $P_{\nu}(x)$  are used

$$K_{p}(x) = 1.0$$
 (7)  
 $K_{p}'(x) = -\left(\frac{p^{2}+0.25}{2}\right)$ 

#### 4. Results

The results are arranged in two sets of tables. The first table lists the value of the conical function  $K_p(x)$  for x = 1.0(0.4)9.8 and for p = 0.1(0.1)10.0 to 8 decimal places. The second table lists the zeros, value of the derivative at the zeros, bend points and value of the function at the bend points, again to 8 decimal places.

## Table 1

The value of the conical function  $K_p(x)$  for x = 1.0(0.4)9.8 and for p = 0.1(0.1)10.0 to 8 decimal places.

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X KP(X) 1.52 0.901520790 00 2.0 0.616096950 00 2.4 0.515859700 00 3.2 0.365849710 00 4.4 0.258815060 00 4.4 0.258815000 6.0 0.258815000 6.1 0.179730790 00 6.0 0.179730790 00 6.0 0.179730790 00 6.0 0.179730790 00 6.0 0.3658520 00 6.0 0.36580000 6.0 0.36580000 6.0 0.36580000 6.0 0.3680000 6.0 0.3680000 6.0 0.3680000 6.0 0.3680000 6.0 0.3680000 6.0 0.38400388000 8.0 0.38400380000 8.0 0.38400380000 8.0 0.38400380000000000000000000000000000000	X 1.2 0.88440056D 00 1.6 0.69869324D 00 2.0 0.55641355D 00 2.4 0.44434535D 00 2.8 0.35417321D 00 3.2 0.28038397D 00 3.6 0.21916934D 00 4.8 0.87175704D-01 5.6 0.27660570D-01 5.6 0.35640008D-02 6.9 -0.17237560D-01 7.6 -0.3560580D-01 7.6 -0.3560580D-01 8.0 -0.3560580D-01 8.0 -0.3560580D-01 9.2 -0.10793324D 00 10.0 -0.124918D-01
X	1.1 1.5 1.5 1.5 1.5 1.5 1.5 1.5
X KP(X) 1.0 0.100000000001 1.8 0.67499126D 00 2.2 0.56335412D 00 2.5 0.47288219D 00 3.8 0.33587267D 00 4.2 0.28244173D 00 4.5 0.1977861D 00 5.0 0.1977861D 00 5.0 0.1977861D 00 5.0 0.1977861D 00 5.0 0.1641466D 00 5.0 0.1021732D-01 7.0 0.43054771D-01 7.0 0.43054771D-01 7.0 0.26135801D-01 8.2 -0.26794566D-01 8.2 -0.26794566D-01 9.4 -0.26335801D-01	X

P= 1.10

X KP(X)  1.3 0.80627734D 00  2.1 0.45971057D 00  2.5 0.34376304D 00  2.9 0.25188452D 00  3.3 0.117791838D 00  4.1 0.67934794D-01  4.5 0.2666734794D-01  4.5 0.2666734794D-01  5.3 -0.36882191D-01  6.9 -0.1789793D 00  8.1 -0.17897924D 00  8.5 -0.15888435D 00  8.7 -0.16823192D 00  9.3 -0.16823192D 00		X
000000000000000000000000000000000000000		000000=0=000000000000000000000000000000
310 160 040 040 040 050 3320 780 770 7770 7770 7770 7770 8800 8800		1040 1300 1300 1300 1300 1300 1300 1300
56 (X) 33 (56 (X) 34 (56 (X)) 37 (25 (X)) 37 (X) 37 (X)		100 100 100 100 100 100 100 100 100 100
00000000000000000000000000000000000000		00000000000000000000000000000000000000
000000000000000000000000000000000000000		000000000000000000000000000000000000000
5880 220 2320 2320 2230 2230 2330 2300 230		848820000000000000000000000000000000000
00000000000000000000000000000000000000		995 X X X X X X X X X X X X X X X X X X
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KP(X)
21451622D
25101910D
113952147D
13952147D
13414793D
15443263D
15443263D
221068620D
22100208D
22100208D
22150217371D
224146326D
225628133D
2256288B
225628B
KP(X)
44696272D
49967680D
32131872D
18940694D
14188507D
44259347D
89564525D
11333361D
11333361D
122624703D
22624703D
226642341D
226642341D
226642341D
226642341D
226642341D
226642341D
226642341D
226642341D
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      00000000000000000000
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      Mr-wamr-wamr-wamr-wamr
                                                                                                                                                                                               MP - 400 MP - 400 MP - 400 MP
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KP(X)
0.82355820D
0.55388551D
0.21889397D
0.311238665D
0.3131928965D
0.3131928997D
0.313192899D
0.313192899D
0.11238665D
0.11238665D
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0.11238665D
0.11238665D
0.11238665D
0.11238669D
0.112389319D
0.1123839309
       0000000000000000000000
                                                                                                                                                                                        KP(X)

80031148D

29287822D

14418719D

36127356D

17748483D

17748483D

17748483D

2207866D

221739398D

221739398D

221739398D

222881700D

224179260D

224207756D

225823399D

222863399D
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90757531D

40355051D

40355051D

25057133D

149610428D

16910428D

16910428D

16927845D

13933037D

16342597D

16342597D

16342597D

16342597D

16342597D

16342597D

16342597D

163468032D

2264681D

226681D

226681D

22687D

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89506111D
33840095D
17696021D
625640134D
17539251D
17039251D
173851D
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        KP(X)
0-100000000
0-67694894D
0-67694894D
0-28463854D
0-16193132D
0-616193132D
0-69152124D-0
0-99215457D-0
0-99215457D-0
0-13782009D
0-0-17782009D
0-0-2259143D
0-0-2259143D
0-0-22531775D
0-0-22531775D
                                                                                                                                                                                         KP(X)

•63662683D

•63662683D

•85948101D-0

•65958035D-0

•74862781D-0

•12539203D

•12539203D

•20982140D

•23854392D

•23854392D
                                                                                                                                                                                                04000040000400004000040
         0400400400400400400
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a

P= 1.30

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KP(X)
0.64472796D
0.32343217D
0.11282361D
-0.27417942D--0.12119102D
-0.22365917D
-0.26801891D
-0.26809177D
-0.26402781D
-0.26402781D
-0.26402781D
-0.26402781D
-0.26402781D
-0.26402781D
-0.26402781D
-0.26609177D
-0.26609179380
       KP(X)
0.68039016D 0
0.38311835D 0
0.18121876D 0
0.41253112D-0
-0.55904051D-0
-0.2077841D 0
-0.23016353D 0
-0.2435200D 0
-0.25540370D 0
-0.255403550D 0
-0.255403550D 0
-0.255403550D 0
-0.25683109D 0
-0.25683109D 0
-0.25683109D 0
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                -- NNNWW444NN000VF000000
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                 00000000000000000000000
                                                                                                                                                                                                     KP(X)
0.74976313D
0.39127283D
0.15766479D
0.25459192D--0.11017926D
-0.21525243D
-0.264343993D
-0.264343993D
-0.26536167D
-0.26536167D
-0.26636167D
-0.2663618D
-0.2663618D
-0.26626403D
-0.17881953D
          KP(X)
0.77568960D
0.44676106D
0.22487257D
0.3546475D-0.3540475D-0.22543163D
0.22543163D
0.22634146D
0.22634146D
0.22634160D
                  U0040U0040U0040U0040U00
                  N4048N4048N4048N4048N40
                                                                                                                                                                                                        KP(X)
0.86755240D
0.46674408D
0.20732322D
0.35673645D-0
-0.20580266D
0.23767839D
-0.23767839D
-0.256495295D
-0.266649542D
-0.26649542D
-0.26649542D
-0.26649542D
-0.26649542D
-0.26649542D
-0.26649542D
-0.26649542D
-0.26649542D
-0.26649530D
-0.26649530D
-0.26649530D
-0.26649542D
-0.26649542D
-0.26649530D
-0.26649530D
-0.26649530D
            KP(X)
88171362D 0
51695409D 0
10497498D 0
12127026D 0
94472494D 0
15236669D 0
15236669D 0
22615291D 0
2261733D 0
2261730D 0
                     11111111
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                     ---NWWWW44NNNOOFFF
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0.10000000
0.252363480
0.723049230--0.195215700-0.252792140
-0.264308890-0.264308890-0.264308890-0.264308890-0.26430890-0.264881730-0.264881730-0.2648900-0.2628949500-0.185144940-0.185148940-0.172521750-0.160057020
              KP(X)
0.100000000
0.32531924D
0.13187970D-0
0.13187970D-0
0.13187970D-0
0.13187970D-0
0.13187970D-0
0.13187970D-0
0.13187970D-0
0.23438720D
0.23438720D
0.23438720D
0.23438720D
0.24687331D
0.255610656D
0.256610656D
0.256610650D
0.256610656D
0.25661000D
0.2566100D
0.256610D
0.25661
                                                                                                                                                                                                            O4000040000400040040
                        040040040040040040040
                        ---NVWWW4400000FFF000000
                                                                                                                                                                                                                   ---NUMMM4400000----
                  X . . . . . .
110
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	10
X 11.3 10.3 11.7 0.26351686D 2.5 10.46706827D-01 2.5 -0.46706827D-01 2.5 -0.17775801D 3.3 -0.23077800D 3.7 -0.25104424D 00 4.5 -0.27576485D 00 4.5 -0.27576485D 00 4.5 -0.27576485D 00 6.1 -0.27576485D 00 6.2 -0.27576485D 00 6.3 -0.27576485D 00 6.3 -0.27576485D 00 6.3 -0.27576485D 00 6.1 -0.27576485D 00 6.2 -0.27576485D 00 6.3 -0.27576485D 00 6.3 -0.27576485D 00 6.3 -0.27576485D 00 6.3 -0.27576485D 00 6.3 -0.27576485D 00 6.3 -0.27576485D 00 6.3 -0.27576485D 00 6.3 -0.27576485D 00 6.3 -0.27576485D 00 6.3 -0.27576485D 00 6.3 -0.27576485D 00 6.3 -0.27576489D 00 6.3 -0.2767643D 00 6.3 -0.276763065D 00 6.3 -0.276763065D 00 6.3 -0.276763065D 00 6.3 -0.276763065D 00 6.3 -0.1777673060D 00 8.3 -0.1177673090D 00 9.3 -0.117600902D 00 9.3 -0.11500902D 00 9.3 -0.11500902D 00 9.3 -0.11500902D 00	1.3 0.56939837D 00 2.1 -0.16304123D-01 2.5 -0.14855053D 00 2.9 -0.22565859D 00 3.7 -0.28600313D 00 4.5 -0.28308273D 00 4.9 -0.27025070D 00 5.3 -0.28308273D 00 6.1 -0.28338448D 00 5.3 -0.25338448D 00 5.3 -0.23411675D 00 6.5 -0.19253400D 00 6.5 -0.19253400D 00 7.7 -0.13098802D 00 8.5 -0.93612670D-01 8.5 -0.60041791D-01
X KP(X) 1.5 0.72260621D 00 2.0 0.92053513D-01 2.4 -0.62221047D-01 2.8 -0.15977339D 00 3.2 -0.22000633D 00 3.6 -0.25516812D 00 4.4 -0.27324579D 00 4.8 -0.27324579D 00 5.2 -0.27783864D 00 5.4 -0.27783864D 00 5.5 -0.27783864D 00 6.4 -0.27783864D 00 5.6 -0.27783864D 00 6.8 -0.27783864D 00 6.9 -0.277837267D 00 8.9 -0.11868961D 00 9.6 -0.11868961D 00 9.6 -0.10419387D 00	X 1.6 0.69429638D 00 2.0 0.27861041D 00 2.0 0.288053990-01 2.4 -0.12165653D 00 3.2 -0.25949182D 00 3.6 -0.28301698D 00 4.0 -0.28301698D 00 5.2 -0.28371698D 00 6.0 -0.21878176D 00 6.4 -0.197877910 00 6.4 -0.1978010 00 6.8 -0.17675720 00 6.8 -0.17675720 00 7.6 -0.13590429D 00 8.9 -0.17675720 00 7.6 -0.13590429D 00 8.0 -0.3805639970-01 8.8 -0.805639970-01 9.6 -0.484656565-01
X KP(X) 1.51 0.41545224D 00 2.3 0.41548386D 00 2.3 0.213948339D-01 2.7 -0.23948339D-01 2.7 -0.23948339D-01 3.5 -0.2768776D 00 3.9 -0.2768776D 00 4.3 -0.27887824D 00 4.3 -0.27887824D 00 5.1 -0.27887824D 00 6.3 -0.27887824D 00 7.5 -0.27887824D 00 6.3 -0.27878495D 00 6.3 -0.27878465D 00 6.3 -0.27878782D 00 6.3 -0.27878782D 00 6.3 -0.278782D 00 6.3 -0.278782D 00 6.3 -0.278782D 00 6.3 -0.278782D 00 6.3 -0.278782D 00 6.3 -0.2787782D 00	X 1.1 1.5 1.5 0.363418600 1.9 0.799917520-01 2.3 -0.909699850-01 3.5 -0.278994140 3.5 -0.278994140 0.278994140 0.0 3.9 -0.2878994140 0.0 3.9 -0.2878994140 0.0 3.9 -0.287280440 0.0 5.1 -0.287280440 0.0 5.1 -0.28728040 0.0 5.3 -0.28728040 0.0 5.3 -0.28728040 0.0 5.3 -0.28728040 0.0 5.3 -0.28728040 0.0 5.3 -0.28728040 0.0 5.3 -0.28728040 0.0 5.3 -0.28728040 0.0 5.3 -0.28728040 0.0 5.3 -0.28728040 0.0 5.3 -0.2873940 0.0 5.3 -0.2848527180-01 0.0 0.0 0.0 0.0 0.0 0.0 0.0 0
X 1.0 1.0 1.0 1.0 1.0 1.0 1.0 1.0	X KP(X) 1.0 0.10000000001 1.4 0.45977867D 00 2.2 -0.56028291D-01 2.6 -0.17205898D 00 3.4 -0.27381985D 00 4.2 -0.27381985D 00 4.2 -0.28805589D 00 4.2 -0.28805589D 00 5.4 -0.28866532D 00 5.4 -0.28866532D 00 5.4 -0.2883607D 00 5.8 -0.28036730 00 5.8 -0.18726336D 00 7.0 -0.16533450D 00 7.0 -0.16533450D 00 7.0 -0.16533450D 00 7.0 -0.16533450D 00 8.2 -0.10720137D 00 8.2 -0.10720137D 00 8.6 -0.1653450D 00 7.9 -0.16533450D 00 7.9 -0.16533450D 00 7.9 -0.16533450D 00 9.9 -0.16532526D 00 9.9 -0.16532520 00 9.9 -0.1653

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KP(X)
0.48978266D 00
0.87640751D-01
-0.12999946D 00
-0.24231460D 00
-0.29281597D 00
-0.29852774D 00
-0.29852774D 00
-0.27783784D 00
-0.27783784D 00
-0.27783784D 00
-0.2783784D 00
-0.2783784D 00
-0.2783784D 00
-0.2783784D 00
-0.2783784D 00
-0.288497869D 00
-0.95528797D-01
-0.88475253D-01
-0.955287970D-01
-0.4349708D-01
            KP(X)

0.53004515D 00

1.4510721D 00

1.02439700D-01

1.02926417175D 00

1.029840672D 00

1.029840672D 00

1.027280834D 00

1.025058901D 00

1.03728450D-01

1.03728450D-01
                                                                                                                                                                                                                                                                                         11000mm44400000rr888000
                      0.66491428D 00
0.22230681D 00
-0.31355350D-01
-0.25203689D 00
-0.29857660D 00
-0.29857660D 00
-0.29304941D 00
-0.2764415D 00
-0.25647484D 00
-0.25626568D-01
                                                                                                                                                                                                                                                                                           00000
                                                                                                                                                                                                                                                                                                                                                                                                                                                                  31896583D-0
17497701D-0
51415744D-0
73770461D-0
                                                                                                                                                                                                                                                                                           0.63454340D

0.16659537D

-0.22132606D

-0.228439699D

-0.30563540D

-0.263540D

-0.2635420D

-0.26753420D

-0.26753420D

-0.19509536D

-0.19509536D

-0.16308540D

-0.16408540D

-0.164
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534D-0
                                                                                                                                                                                                                                                                                                                                                                                                                                                          67D-(
                                                                                                                                                                                                                                                                                   KP(X)
63454340D
16659537D
                KP(X)
6649142
2223068
3135535
                                                                                                                                                                                                                                                                                             V0040V0040V0040V0040V00
                                                                                                                                             4000040000
                           U0040U0040U00
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                           --NNNWW444NN000VV0000000
                   KP(X)
0.82039948D 00
0.31098580D 00
0.19449620D-01
-0.14657009D 00
-0.28197678D 00
-0.29554918D 00
-0.2862957D 00
-0.13510773D 00
-0.13510773D 00
-0.1463950-01
                                                                                                                                                                                                                                                                                              00
                                                                                                                                                                                                                                                                                                                                                                                                                0.10979985D
0.817475940-0
0.55723253D-0
0.31797062D-0
0.99596108D-0
0.98489715D-0
0.27719378D-0
                                                                                                                                                                                                                                                                                    KP(X)
0.80320198D
0.25849651D
0.25849651D
0.38157829D
0.3031779D
0.30459948D
0.28934612D
0.286462330D
0.28494294D
0.23494294D
0.23494294D
0.23494294D
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13759885D-0
58087462D-0
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0.10000000 01
0.36540867D 00
0.19885037D-01
-0.29898527D 00
-0.29898527D 00
-0.29427242D 00
-0.27141045D 00
-0.27141063D 00
-0.177901738D 00
-0.177901738D 00
-0.177901738D 00
-0.23427763D-01
0.23427763D-01
0.23427763D-01
                      KP(X)

0.100000000 01

0.412881270 00

0.77909152D-01

-0.21989626D 00

-0.29584401D 00

-0.29736456D 00

-0.29736456D 00

-0.29736450D 00

-0.29787312D-01

-0.32701385D-01

-0.32701385D-01

-0.32701385D-01
06
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                                   ---NNMMM4400000NNN00000
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0-000000000
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0.40719734D

-0.22299391D

-0.31980887D

-0.31980887D

-0.26490332D

0.17499647D

0.13002537D

0.13002537D

0.14797409D

0.14797409D

0.15700784D

0.1570084D
                                                                                                                                       Mr-400mr-400mr-400mr
                                                                                                                                      KP(X)

0. 60326983D

0. 13978807D

0. 26028844D

0. 26028830D

0. 26098837D

0. 26098837D

0. 26098327D

0. 260983200

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0. 26098300

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-0.29991283D 00

-0.27156850D 00

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-0.25650236D-01

-0.5484653D-01

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0.66325699D-03

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0.786353548D-01

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-0.19437181D 00

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14227866D 00
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       KP(X)

-0.12412220D 001

-0.35543984D 000

-0.29150075D 000

-0.29156287D 000

-0.82166287D 000

-0.82166287D 000

-0.820618382D 000

-0.820437250D 000

-0.820430D 000

-0.82044669D 000

-0.828600247D 010

-0.85983166D 010

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-0.8600247D 010

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-0.8660247D 010
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-0.15977194D

-0.25931089D

-0.26424422D

-0.2689331D-00

-0.2689331D-00

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-0.26889331D-00

-0.276765B

-0.17658D

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-0.17670805D-0

-0.4030513D-0

-0.94030513D-0
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-0.39930969D-01
-0.36043616D 00
-0.27331407D 00
0.42403598D-01
0.14207535D 00
0.20854611D 00
0.18279003D 00
0.1551447D 00
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0.1527903D 00
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-0.35384168D 0
-0.24335627D 0
-0.559710614D-0
0.17474848D 0
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0.21693417D 0
0.159595105D 0
0.159595105D 0
0.34164767D-0
-0.40697547D-0
-0.37627670D-0
-0.12980051D 0
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-0.32109986D 00

-0.15006812D 00

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0.1797107D 00

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0.17971007D 00

0.14532152D-01

-0.1353163D-01

-0.1353163D-01
KP(X)

-0.28814082D

-0.33945777D

-0.18975844D

0.98334294D-01

0.256834294D-01

0.27439338D

0.21441795D

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KP(X)

-0.22401097D 01

-0.35166793D 00

-0.19707794D 00

-0.18377563D-01

0.1837337D 00

0.21349263D 00

0.21349263D 00

0.18819647D 00

0.18819647D 00

0.1885331D-01

-0.45218819D-01

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-0.45218819D-01

-0.45218819D-01

-0.12333310D 00

-0.12333310D 00

-0.12333310D 00

-0.13594662D 00
     KP(X)
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19311636D
23252089D
633252089D
713815768D
20401868D
1077747651D
10591650D
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KP(X)

14205111D 0

25685320D 0

17315461D 0

29076006D-0

1857269D 0

1857269D 0

19287606D 0

19287606D 0

14542033D 0

163469D 0

1633469D 0

16448728D 0

14430624D 0

144468728D 0

144468728D 0

144468728D 0

14456880 0

14456880 0
      KP(X)
36259704D
20973426D
12859704D
20973426D
125046690D
20930737D
212904807D
212904807D
213627230D
213627230D
213627230D
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13502035D 0
22549780D 0
225549780D 0
225549780D 0
15825756D 0
11644609D-0
12643274D 0
135432374D 0
14351871D 0
14351871D 0
144853374D 0
144859376D 0
         KP(X)
362679600 0
362679600 0
6002477500
9602477500
18454340 0
221181630 0
221181630 0
2251181630 0
2251181630 0
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0.197470511D

0.25600433D

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0.12875816D

0.12875816D

0.11796478D
          KP(X)

-0.33052694D

-0.30423476D

-0.10830872D

0.63847192D-01

0.217063894D

0.21711656D

0.21952299D

0.19381275D

0.1092369D-01

-0.289934005D-01

-0.89930189D-01

-0.89930189D-01

-0.89930189D-01

-0.13528232D

-0.13528232D

-0.1352832D

-0.1352833000

-0.13531936D

-0.13857874D

-0.13857874D
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-0.25232756D 00

-0.34056819D 00

-0.15870903D 00

0.26551899D-01

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0.22213146D 00

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0.20212355D 00

0.16954705D 00

0.68279915D-01

-0.83733920D-01

-0.83733920D-01

-0.13335207D 00

-0.13948302D 00

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- 27795482D 0

- 32498177D 0

- 11823942D 0

- 181230998D 0

- 2380660D 0

- 22380660D 0

- 21762489D 0

- 13219422D 0

- 13219422D 0

- 25412670D 0

- 25412690D 0

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-0.172829369

-0.346805550

-0.13435513D

0.18961789D

0.228674729

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0.168584960

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0.10298840

0.50191740D-

0.50181224D-

-0.52040891D-

-0.14530465D

-0.14614199D

-0.14614199D
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-0.20179605D

-0.33269124D

-0.11024365D

0.21240172D

0.21264172D

0.13767941D

0.7505273D

-0.46814359D

-0.15010114D

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-0.36532043D 00

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0.68518971D-01

0.226538971D-01

0.22653897D-01

-0.34263382D-01

-0.34265382D-01

-0.34265382D-01

-0.15204787D 00

-0.13624160D 00

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-0.13624161D 00

-0.15204787D 00
  0.559743120-01
-0.36850327D 00
-0.19390398D 00
0.27144349D-01
0.2510830671D 00
0.22118301D 00
0.22118301D 00
0.64960726D-01
0.64960726D-01
-0.41181995D-01
-0.41181995D-01
-0.4123759D 00
-0.13132759D 00
-0.14834386D 00
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KP(X)

-0.32078392D

-0.28173248D

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0.22357520D

0.22591704D

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-0.54458638D-0

-0.54458638D-0

-0.58635181D-0

-0.58635181D-0

-0.12642601D

-0.126426001D

-0.126420001D
     KP(X)

-0.30079890D

-0.30524377D

-0.76510893D-0

0.20899495D

0.22899495D

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-0.12858972D

-0.1383689D
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XP(X)

1 -0.25392509D

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5 0.17669567D

0.23710118D

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0.6215780D-0

-0.1228873D-0

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-0.22885637D 0
-0.31478616D 0
-0.53128332D-0
0.22832271D 0
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0.29234273D-0
-0.34446939D-0
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-0.15394110D 0
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        KP(X)

38677412D-01

34746739D 00

14250694D 00

22310085D 00

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16183645D 00

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16183645D 00

152255414D-01

48509423D-02

59924341D-01

13922725D 00

13922725D 00

15828404D 00

15141993D 00

15177754D 00

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151714800D 00

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229600033D-01

46753021D-01

46029960D-03
        XP(X)
35825051D 0
10723736D 0
21732622D 0
231742622D 0
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12010528D 0
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113202096D 0
11323251472D 0
115231528D-0
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11523251472D 0
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11521472D 0
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-0.37216764D

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0.61060891D-01

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0.64813701D-01

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   KP(X)

-0.33785221D 00

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0.17657936D-02

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0.83152259D-01

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-0.351964570 00

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KP(X) -0.27692695D 00 -226887934D 00 -226887934D 00 -2263194D 00 0.23866546D 00 0.19766744D 00 -19766744D 00 -19766744D 00 -19766744D 00 -19766744D 00 -19766744D 00 -19766744D 00 -0.15961254D 00 -0.15961254D 00 -0.15961254D 00 -0.15961254D 00 -0.15961254D 00 -0.15961254D 00 -0.15961330 00 -0.26509777D-02	C C C C C C C C C C C C C C C C C C C
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X 1.61	X 10.1 10.2 10.30093951D 10.36660913D 10.76366456D 20.300176668889D 20.1 20.2 30.1 30.2 30.
X 1=0 1=0 1=0 1=0 1=0 1=0 1=0 1=0	x 1-0 1-0 1-0 1-0 1-0 1-0 1-0 1-0

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XP(X)

7 -0.18029675D

1 0.13668564D

0 0.24585562D

0 0.20238258D

-0.11070891D-0

-0.14704140D

-0.16572230D

-0.16572230D

-0.145659D

-0.145669D

-0.145669D

-0.145669D

-0.145669D

-0.16572230D

-0.16572220D
         -0.31648044D 0
-0.21192532D 0
0.23748871D 0
0.22088346D 0
0.22088346D 0
0.13260504D 0
-0.13260504D 0
-0.15651900D 0
-0.1565190D 0
-0.15651900D 0
(X)
48044
2532
                                                                                                                                                                                                                                                                             KP(X)

-0.15176020D 00

0.25171984D 00

0.25578106D 00

0.22213453D 00

0.14392617D-01

-0.14392617D-01

-0.15710265D 00

-0.13710265D 00

-0.16571225D 00

-0.173139842D-01

0.49128250D-01

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0.49128250D-01

0.49128250D-01

0.49128250D-01
  KP(K)
-0.12511011D 00
-0.29503041D 00
0.23374834D-01
0.23476224D 00
0.15778834D 00
-0.39478470D-01
-0.10882191D 00
-0.15026821D 00
-0.1562838D 00
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-0.35758655D 00

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-0.34587808D 00
-0.25012085D 00
-0.2369885D 00
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-0.236983187D-01
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-0.16972799D 00
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00.559362793D-01
-0.156049310D-01
-0.16481698D
00.17169094D
00.12366946D
00.123698536D-01
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-0.35902742D
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0.26533726D-01
-0.15841348D 00
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-0.15912896D 00

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-0.16921461D 00

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-0.13721657D-01

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-0.29508751D-01

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-0.31788763D 00

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-0.30211504D 0
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0.24909074D 0
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-0.128310681D 00
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0.99265451D-01

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0.65584250-
0.6582815800-
0.65882800-
0.943865750-
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KP(X)
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10013722D 00

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-0.16522813D 00

-0.97219045D-01

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X 16.3 -0.364959500 00 2.1 0.224608880 00 2.5 0.104971500-01 2.9 -0.152589610 00 3.3 -0.191791590 00 3.1 -0.5291791590 00 4.5 0.34922590 00 5.7 0.143533700-01 5.7 0.143533700-01 5.7 0.143533700-01 5.8 0.102707650 00 6.9 0.668293460-01 7.7 -0.863293460-01 7.8 -0.668293460-01 7.9 -0.876382470-01 8.9 -0.678536380-01 8.9 -0.678536380-01		X KP(X) 1.3 -0.35512445D 00 2.1 0.20548466D 00 2.5 -0.24528120D-01 2.9 -0.17262408D 00 3.7 -0.18627306D 00 4.5 -0.1521026D 00 4.5 0.67114581D-01 4.9 0.12231137D 00 5.7 0.13919140D 00 6.1 0.11377399D 00 6.5 0.76744913D-01 6.5 0.76744913D-01 6.5 0.76744913D-01 8.1 -0.52819311D-02 7.7 -0.411053960-01 8.5 -0.10397473D 00 9.3 -0.10397473D 00
X KP(X) 1.52 -0.35597772D 00 1.66 0.69811407D-01 2.0 0.25549588D 00 2.0 0.66474923D-01 2.8 -0.12341360D 00 3.2 -0.19247984D 00 3.6 -0.15925357D 00 4.8 0.85358806D-01 4.8 0.85358806D-01 5.6 0.12865540D 00 6.9 0.12865540D 00 6.0 0.13539575D 00 6.0 0.13539575D 00 6.1 0.13539575D 00 6.2 0.33618660D-01 8.8 -0.33618660D-01 8.8 -0.3361869D-01 9.2 -0.37926111D-01 9.6 -0.10583869D 00		X 1.2 -0.36554701D 00 1.6 0.99602895D-01 2.0 0.24600864D 00 2.4 0.31894841D-01 2.8 -0.14900430D 00 3.2 -0.19331004D 00 4.0 -0.42179531D-01 4.4 0.48278013D-01 4.8 0.42179531D-01 6.0 0.12152818D 00 6.4 0.8667332D-01 6.8 0.45260589D-02 7.2 0.45260589D-02 7.6 -0.32734385D-01 8.9 -0.6550531D-01 8.8 -0.6550531D-01 8.9 -0.10872126D 00 9.2 -0.10872126D 00 9.6 -0.10872126D 00
KP(X)  1.5  1.65		X 10.1 -0.38121767D-01 10.9 0.2646833D-01 20.3 0.92022972D-01 20.7 -0.11623132D 00 30.9 -0.15733876D 00 30.9 -0.15733876D 00 30.9 -0.15733876D 00 30.9 -0.15733876D 00 30.9 -0.15733876D 00 50.9 0.15733870-01 50.9 0.12843656D 00 50.9 0.10758935D 00 50.9 0.10758935D 00 50.9 0.1055055TD 00
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-0.3434608D 00

0.21905038D 00

0.18285865D 00

-0.18720244D 00

-0.1744714D 00

-0.17270976D-01

0.96419918D-01

0.12683131D 00

-0.90371440D-01

-0.90371440D-01

-0.90371440D-01

-0.92918824D-01

-0.92918824D-01

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-0.10656335D 00
                                                                                                                                                       KP(X)

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0.23608007D

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-0.90129746D-01

-0.19599502D

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0.14006804D

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-0.61831050D-0

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-0.93199759D-0

-0.99374962D-0

-0.99374962D-0
KP(X)
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21905038D
18285865D
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-0.36831274D-01
-0.18567761D 00
-0.17596628D 00
-0.27597860D-01
0.27597860D-01
0.172978D 00
KP(X)

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0.23245700D 00

0.12778076D 00

0.18770613D 00

0.18770613D 00

0.79863370D-01

0.79863370D-01

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0.13188173D 00

0.137612D 00

0.137652D 00

0.12662769D-01

0.28626707D-01

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0.587566060-01

-0.19504105D 00

-0.32714501D-01

0.612495990-01

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-0.110055050-01

-0.55038905D-01

-0.550389050-01

-0.1007325D 00

-0.11007325D 00
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0.21898323D

0.21966556D-0

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0.16002537D 0
0.91269203D-0
0.13310212D 0
0.13373182D 0
0.24687387D-0
0.72310777D-0
0.13143079D 0
0.13165D-0
0.13165D-0
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0.14922075D 0
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74804430D-0
97977167D-0
11050114D 0
11328917D 0
   KP(X)
100000000
18656932D
26114285D
12249142D
10540747D
19525071D
15646451D
58958022D-0
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9218638D-
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97280744D
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-0.31494824D 00

0.24986063D 00

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-0.19884526D 00

-0.13411816D 00

0.134174D-01

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0.18249777D-01

0.18249775D-01

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0.182617908D 00

0.29956900D-01

-0.91460450D-01

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-0.916660D 00

-0.11521833D 00

-0.11521833D 00

-0.99273666D-01
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0.98971917D-01
-0.19577126D 00
-0.10717787D 00
0.16122802D-01
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0.15005632D 00
0.1562304D 00
0.1562304D 00
0.52119132D-01
-0.32117913D-01
-0.87014360D-0
-0.11683235D 0
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29834088D
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-0.38479774D 00

0.17824250D 00

-0.19442152D 00

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-0.15858349D 00

-0.47652585D-01

0.60572308D-01

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0.13459743D 00

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-0.41660D 01

-0.85812397D-01

-0.86425020D-01

-0.86425020D-01

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0.20006135D

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-0.12179515D

0.23059281D

-0.42240139D-01

-0.19443959D

-0.19443959D

-0.46182709D

-0.485621059D-01

-0.485621059D-01

-0.485621059D-01

-0.486621059D-01

-0.2460811D-01

-0.28544666D-01

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-0.28544666D-01

-0.2963508D-01

-0.3963508D-01

-0.39633017D

-0.99633017D

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-0.99633017D
                  KP(X)
-0.10167292D
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0.17827913D
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0.14151095D
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0.43124119D-0
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-0.13274250D 00

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-0.11507620BD-01

-0.370889153D 00

-0.95304272D-01
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0989846D-0
5569706D-0
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KP(X)  1.51 -0.1034509D 00 1.55 00.10341390D 00 20.3 -0.21298059D 00 20.3 -0.20113440D 00 30.1 -0.104115D-01 30.5 -0.10721658D 00 50.5 -0.10721658D 00 50.5 -0.90686072D-02 60.7 -0.90686072D-01 7.5 -0.90686072D-01 7.5 -0.90686072D-01 80.7 -0.90686072D-01	X 1. 1 -0.16029932D 00 1. 5 0.12958821D 00 2. 3 -0.10348000D 00 2. 7 -0.20228920D 00 3. 5 0.23033047D-01 3. 5 0.12168928D 00 4. 7 0.12168928D 00 4. 7 0.12168928D 00 5. 1 0.81469614D-01 5. 5 -0.41063238D-01 6. 3 -0.41063238D-01 6. 3 -0.41063238D-01 6. 3 -0.41063238D-01 8. 3 -0.41957103D-01 8. 3 -0.96159527D-01 8. 3 -0.96159527D-01 8. 3 -0.41937103D-01 9. 5 0.15973348D-01
X	X

P= 6.10

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-0.88965842D-0

-0.20388126D

-0.84129672D-0

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-0.157869638D-0

-0.746275D-0

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-0.72395985D-0

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-0.94885767D-0

-0.994838917D-0

0.994838917D-0
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0.26827021D 00

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-0.46851763D-01

0.96817172D-02

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-0.27303300D 00

-0.2766561D-02

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0.49774D 00

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0.93819479D-01

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-0.56945877D-01

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37609146D
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0.86964737D-01
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0.11275428D-01
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0.1221569D 00
0.32934530D-01
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0.20159937D-01

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•23927146D-02
•37750298D-01
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0.1000000D
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0.18001905D
0.18525370D
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-0.14068573D 00
-0.18418962D 00
-0.18418962D 00
-0.13186338D 00
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KP(X)
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-0.38818470-01
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-0.67512528D-01

-0.20560872D 00

-0.58396906D-01

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-0.36910388D 00

0.27475097D 00

-0.20851857D 00

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-0.15937416D 00

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              CO.26034068D 00

CO.24665727D 00

CO.26891047D-01

CO.20829162D 00

CO.73201749D-01

CO.132576785D 00

CO.132576785D 00

CO.13957995D 00

CO.123787D-01

CO.12377918D 00

CO.10157944D 00

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CO.10157944D 00
              -0.24539107D 00

0.23267787D 00

0.57154894D-01

-0.20127798D 00

-0.13189782D 00

0.43166824D-01

0.87139203D-03

0.87258253D-01

0.87258253D-01

-0.72401210D-01

-0.12778845D 00

-0.12778845D 00

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-0.38746804D-01

0.34999795D-02

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-0.16929182D-00

0.13318749D-00

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0.122258489D-01

-0.56489013D-01

-0.56489013D-01

-0.10830790D-00

-0.10830790D-01

-0.49433156D-01

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0.659831046D-02

0.659831046D-01

0.6598311D-01

0.65983711D-01

0.95281936D-01
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-0.16226564D 00

0.20946665D-01

0.14911966D 00

0.15324704D 00

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-0.19212936D-01

-0.12779566D 00

-0.12779566D 00

-0.12779566D 00

-0.12779560D-01

-0.4231458D-01

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-0.12287700D 00

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-0.27167472D 0
0.10045238D 0
0.12098356D 0
0.27944721D 0
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-0.115451D 0
-0.12098356D 0
-0.25059715D 0
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0.12022151D 00
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-0.69392708D-01
-0.13180342D 00
-0.13180342D 00
-0.130462940D 00
-0.13046343D-01
-0.72260714D-02
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-0.74379439D-03

-0.21257166D 00

-0.69076099D-01

0.12903690D-01

-0.45428109D-01

-0.13446517D 00

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-0.15535023D-01

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C. 85321883D-

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C. 15108665D-

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C. 12790365D-

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C. 28166009D-

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0.92632530D-01

0.99182768D-01

-0.1061678D-01

-0.1384956D-01

-0.13584484D 00

-0.13584956D-01

0.9182768D-01

0.85651317D-03

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0.33429203C

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1.1062399D

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1.5635374D-0

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             KP(X)

-0.32389944D

0.27843338D

-0.11500377D

0.455870044D

0.12519058D

0.10370048D-0

-0.89318648D-0

-0.12304451D

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-0.298562110-01

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0.2988672590-01

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470 920 920 930 930 930 930 930 930 930 93	760 760 760 7500 8800 8800 8800 8800 8800 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 97000 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 9700 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 97000 970000 97000 97000 97000 97000 97000 97000 97000 97000 970
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KP(X)  1.0  1.0  0.244370960  1.0  2.2  0.674837180-01  2.2  0.16737890-01  3.0  0.171296990  3.4  0.107865270  4.2  0.107865270  4.2  0.1178680  0.1078680  6.2  0.10737345890-01  6.2  0.10737345890-01  6.2  0.107395880  7.0  0.107395880  7.0  0.107395880  7.0  0.295743390-01  6.2  0.804891970-01  6.3  0.295743390-01  6.6  0.107395880  7.0  0.107395880  7.0  0.107395880  7.0  0.273362230-01  8.2  0.656580560-01  9.8  -0.656580560-01	X

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49867320D-0
90515880D-0
48978866D-0
13133665D 0
72804879D-0
35810342D-0
68356686D-0
10591713D 0
11329830D 0
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-0.23607840D 00

-0.15484447D 00

-0.15320514D-01

0.15624973D-01

-0.16654830D-01

-0.11876392D 00

-0.13765220D 00

-0.1376522D 00

-0.1376522D-01

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14234925D 00
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851930240-01
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-0.159887970 -0.159887970 -0.192723580 -0.960804930-01 -0.173641550-01 -0.1779250 -0.1779250 -0.17779250 -0.17779250 -0.17779250 -0.17779250 -0.17779250 -0.17779250 -0.17779250 -0.17779250 -0.17779290-01 -0.194605700-01 -0.194605700-01 -0.194605700-01 -0.194605700-01 -0.353937390-01 -0.353937390-01 -0.353937390-01		PKP(X)  -0.13971292D  0.1834893D-01  -0.17803981D  0.15251173D  -0.13256020D  0.4676649D-01  0.4676649D-01  0.4676649D-01  0.4676649D-01  0.4676649D-01  0.4676649D-01  -0.4676649D-01  -0.4676649D-01  -0.4676649D-01  -0.46635180D-01  -0.46635180D-01  -0.46635180D-01  -0.46635180D-01
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X 1.61 -0.38453908D 00 1.95 -0.22026365D 00 2.7 0.22026365D 00 2.7 0.17699852D 00 3.1 0.60617709D-01 3.5 -0.92948257D-01 3.9 -0.14438041D 00 4.7 -0.94692311D-01 5.5 0.11532133D 00 5.9 0.11532133D 00 5.7 0.26449562D-01 7.5 -0.64467613D-01 8.7 -0.9745864D-01 8.7 -0.9745864D-01 9.1 -0.72196081D-01 9.5 -0.46537403D-01		X KP(X) 1.55 -0.38847359D 00 1.69 -0.21842265D 00 2.7 0.17357317D 00 3.1 0.32236028D-01 3.5 -0.11381542D 00 3.9 -0.14074139D 00 4.7 0.25783975D-01 5.5 1 0.963848D-01 5.9 0.99785375D-01 6.7 -0.1961849D 00 5.9 0.96335984D-01 6.7 -0.49871971D-01 7.5 -0.97687107D-01 8.3 -0.97687107D-01 8.3 -0.97687107D-01 8.3 -0.97687107D-01 8.3 -0.97687107D-01 9.5 -0.21644180D-01
X 1.00 0.10000000000000000000000000000000	3.20	X KP(X) 11.0 0.10000000 01 11.8 -0.20533616D 00 2.2 -0.21775293D-01 2.6 0.17720945D 00 3.0 0.77113869D-01 3.8 -0.14547123D 00 4.2 -0.93740826D-01 5.4 0.25040943D-02 5.8 0.10798450D 00 6.2 0.65356065D-01 7.8 -0.3504091D-01 8.2 -0.97020201D-01 8.6 -0.83273836D-01 9.4 -0.29390865D-01

40			
	x KP(x) -3 0.20847691D 00 -7 -0.14565986D 00 -5 0.179970966D-01 -3 0.179970966D-01 -3 -0.78518339D-01 -7 -0.14652224D 00 -1 -0.93816299D-01 -0.93816299D-01 -0.0036482D 00 -0.0036482D 00 -0.0036682D 00 -0.003682D		X KP(X) 1.3 0.22326922D 00 2.1 -0.53212614D-01 2.5 0.17796226D 00 2.9 0.7796226D 00 3.7 -0.10157718D 00 3.7 -0.14430134D 00 4.5 0.35255398D-01 4.5 0.35255398D-01 6.1 0.29478208D-01 6.1 0.29478208D-01 6.5 -0.29456013D-01 7.7 -0.98635733D-01 8.9 -0.74243594D-01 8.9 -0.5291957D-01 9.3 0.15125304D-01 9.3 0.15125304D-01
	(X) 24465D 00 258208D-02 1 46621D 00 05892D 00 02892D 00 028843D-01 00 028843D-01 00 028843D-01 00 03860D-01 00 03860D-01 00 03860D-01 00 03890D-01 00 03891D-01 00 03891D-01 00 03891D-01 00 03891D-01 00 03891D-01		1x) 52561D-01 19217D-01 33804D 00 38495D 00 38495D 00 385385D-01 37862D 00 37862D 00 37862D 00 37862D 00 37862D 00 37862D 00 31211D-01 74227D-01 74227D-01 74227D-01 74227D-01 74227D-01 74527D-01
	X 1 1 1 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2		11.
	KP(X)  0.39159978D 00  0.16281592D 00  0.82852210D-01  0.316549707D 00  0.316549707D 00  0.3469187D-01  0.53018691D-01  0.547965065D-01		KP(X)  -0.39392177D 00  0.14147008D 00  0.20463856D 00  0.15304365D 00  -0.25657770D-01  -0.14169384D 00  -0.14169384D 00  -0.16195931D-01  0.77316242D-01  0.77316242D-01  0.77316242D-01  0.1079257D 00  0.1079257D 00  0.10793642D-01  -0.54268105D-01  -0.54268105D-01
	KP(X)  4 0.27795894D 00  8 -0.21429398D 00  9 0.65463488D-02  9 0.17873419D 00  10 0.49874298D-01  4 -0.10874757D 00  8 -0.10874757D 00  8 -0.10874757D 00  9 0.1201749D-01  9 0.31060323D-01  9 0.3573368D-01  9 0.93573368D-01  9 0.93573368D-01  9 0.93573160000  9 0.93573368D-01	B + 40	KP(X)  0 0.100000000  8 0.27270182D  0 34396023D-0  0 17564373D  0 0.21423481D-0  0 0.22993739D-0  0 0.22993739D-0  0 0.22993739D-0
		O.	UNWWW44WW0006666

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KP(X)

0.23661528D 00

-0.18128697D 00

-0.18054156D 00

0.18054156D 00

-0.42827546D-01

-0.42827546D-01

-0.42827546D-01

-0.42827546D-01

-0.42827546D-01

-0.428476D-01

-0.58682251D-01

-0.28682251D-01

-0.28682251D-01

0.39960207D-01
                                                                                                                                                                                                                                                                                                                                KP(X)
-0.19548653D
-0.18236892D-
0.17868910D-
0.12343049D-
0.12343049D-
0.15580552D-
0.15580552D-
0.10001247D-
-0.2561647D-
-0.2561000-
-0.2561000-
-0.2561000-
-0.2561000-
-0.2561000-
-0.2561000-
-0.2561000-
-0.2561000-
-0.2561000-
-0.2561
             KP(X)
-0.5843808D-01
-0.5843808D-01
-0.16836342D 00
0.92699946D-01
-0.9182445D-01
-0.9182445D-01
-0.9182445D-01
0.1927997D 00
                                                                                                                                                                                                                                                                                                                                 // 00 + B // 00 
                                                                                                                                                                                                                                                                                                                                   396176070 00

178447330 00

146673200 00

116655930 00

148397840 00

742788650-01

402468560-01

112682530 00

118205470 00

118205470 00

118205470 00

118205470 00

118205470 00

118205470 00

118205470 00

118205470 00

11820547340 00

118205800 00

118205800 00

118205800 00

118205800 00

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118205800 00
                 KP(X)
0.39617607D t
0.95844167D-0
17844733P
               -0.39544498D 0
-0.19299052D 0
-0.19299052D 0
0.12827993D 0
-0.53316704D-0
-0.97367624D-0
0.97367624D-0
0.975187774D 0
-0.97518775D-0
0.97518775D-0
0.97518775D-0
0.97518775D-0
0.97518775D-0
0.97562454BD-0
-0.98590996D-0
-0.7562454BD-0
       KP(X)
3954449
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        ---NNMMM44NNN001---
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0.26546820D 01

0.61221212D-01

0.61221212D-01

0.1367230D 00

-0.13952257D 00

-0.13952257D 00

-0.15203054D-01

0.81383228D-01

0.81383228D-01

0.12197593D 00

0.12197593D 00

-0.156851512D-01

-0.55861512D-01

-0.55861512D-01
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0.26744981D 00

0.55457238D-01

0.16217888D 00

-0.41489578D-01

-0.40333302D-01

0.40333302D-01

0.40333302D-01

0.40333302D-01

0.40333302D-01

0.40333302D-01

0.11567315D 00

0.11699761D-01

0.106699761D-01

0.106699761D-01

0.106699761D-01

0.106699761D-01
    0.25874934D 00
0.20706954D 00
0.17250152D 00
-0.13839226D-01
-0.10571857D 00
0.10571857D 00
0.10332665D 00
0.12268881D 00
0.13175814D-01
-0.51226240D-01
-0.92202640D-01
-0.92202640D-01
-0.92202654D-01
                                                                                                            153D-01
                                                                                                                        310-01
                                                                                                                               919465D-0
                                                                                                                                     27910-0
                                                                                                      5606365
1572535
2365470
5574363
KP(X)
2587493
                                                                                                                                76919
                                                                                                000000
8757
8757
8757
8757
8757
                                                                                                                                                                                       0.10841956D-01
0.10516864D 00
0.76232321D-01
0.35694758D-02
-0.61085445D-01
-0.98108241D-01
-0.80914039D-01
0.24214945D-01
0.24214945D-01
0.39258415D-01
 KP(X)
-0.36024187D-01
-0.10542450D 00
-0.68527299D-01
-0.42982644D 00
-0.12811085D 00
-0.12811085D 00
-0.17421892D-01
-0.86835130D-01
-0.86835130D-01
-0.95865148D-01
-0.95865148D-01
-0.95865148D-01
-0.95865148D-01
-0.95865148D-01
-0.958651435D-01
-0.86835130D-01
-0.86835130D-01
-0.86835130D-01
                                                                                                                                                                                                         18199826D
12044853D-0
14231531D
11164300D
10841956D-0
                                                                                                                                                                                  KP(X)
15274557D-0
12691920D
                                                                                                                                                                                                                                                                                                                               8271014D-
                                                                                                                                                                                                    22156730-
                                                                                                                                                                                         0000000
                                                                                                                                                                                        -0.39529483D 00

0.47766076D-01

0.17302664D 00

0.68646709D-01

-0.13345447D 00

-0.20825414D-01

0.88219946D-01

0.91852510D-01

0.91852510D-01

0.91852510D-01

-0.4702519D-01

-0.4702519D-01

-0.4702519D-01

-0.53254007D-01

-0.53254007D-01

-0.53254007D-01

0.87415054D-01

0.87415054D-01
  KP(X)

-0.39612297D 00

0.72000240D-01

-0.16126925D 00

0.93788424D 00

-0.10173935D-01

-0.48400152D-01

0.65872513D-01

0.1078262D-01

-0.48645562D-01

-0.48645562D-01

-0.4864562D-01

-0.86174901D-01
                                                                                                                                                                                           KP(X)
0.100000000 01
0.23273597D 00
0.13048753D 00
0.12145251D 00
-0.86905468D-01
-0.53098078D-01
0.66420988D-01
0.66420988D-01
0.10521453D 00
0.10521853D-01
-0.836605253D-01
   KP(X)
0.100000000 01
0.24539002D 00
-0.21852328D 00
0.10973254D 00
0.14057780D 00
-0.62348176D-01
-0.78463341D-01
0.40939562D-01
0.40939562D-01
0.65709695D 00
0.65709695D-01
-0.64316D-01
-0.97809384D-01
-0.64316D-01
-0.97809386D-01
-0.97809380D-01
-0.66333514D-01
-0.54970349D-01
                                                                                                                                                                                                                                                                                                  0.212993230-01
0.207135250-01
0.552283930-01
0.778462560-01
0.871219040-01
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           --- N/WWW.44000000----
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0.27452923D 00
0.80508450D-01
0.14801868D 00
-0.67598903D-01
-0.15037870D 00
0.152821373D-01
0.12497873D 00
0.12497873D 00
0.12497873D 00
0.12497873D-01
-0.40810417D-01
-0.40810417D-01
-0.40810417D-01
-0.5134915D-01
0.35950752D-01
0.35950752D-01
0.87807312D-01
                                                                                                              0.22515821D 00
0.10377133D 00
0.13040768D 00
0.327133D 00
0.32714424D-01
0.8485237D 00
0.8485237D 00
0.62223602D-01
-0.10193037D-01
0.268279539D 00
-0.26827954D-01
0.26827954D-01
0.26827954D-01
0.26827954D-01
0.26827954D-01
                                                                                                               E.
                                                                                                               0.25621798D-01
0.11234758D-01
0.11234758D-01
0.17021893D 00
0.15246690D 00
0.15246690D 00
0.12604441D 00
0.126018976D-01
0.26615995D-01
0.26615995D-01
0.26615995D-01
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0.26615995D-01
0.49871439D-01
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0.49871439D-01
0.49871439D-01
   0.52990018D-02
-0.14678709D 00
-0.17813686D 00
-0.15949145D-01
-0.91131779D-01
-0.91312779D-01
0.11842342D 00
0.11842242D 00
0.11842242D 00
0.52713132D-01
-0.23727475D-01
-0.95365734D-01
-0.95365734D-01
-0.6194677D-01
                                                                    261946770-01
608606360-01
819559430-01
883712950-01
815363000-01
KP(X)
                                                                      000000
                                                                                                                @U004@U004@U004@U00
     -0.39135605D 00
-0.97575050D-03
-0.97575050D-03
-0.18399083D 00
0.18384710D-01
-0.18384710D-01
0.3479959D 00
0.11363978D 00
0.11363978D 00
0.46991140D-01
-0.8986752D-01
-0.8986752D-01
-0.8986752D-01
-0.8986752D-01
-0.8986752D-01
-0.8986752D-01
-0.8986752D-01
    -0.39370201D 00

0.23366228D-01

-0.12023315D 00

0.41933224D-01

-0.13618826D 00

-0.1368885D 00

0.12631687D 00

0.12534687D 00

0.12534685D 00

0.12534685D 00

-0.1255231D 00

-0.10049395D 00

-0.10049395D 00

-0.15853892D-01

0.53353581D-01

0.53353581D-01

0.69865973D-01

0.69865973D-01
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$3750-01
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1880237
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                                                                                                                  0400400400400400400400
      ---NNWWW4400000VVV000000
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P = 9.10

KP(X)  0.28374964D  0.1226483297D  0.10981160D  0.10281160D  0.1226145D  0.1226145D  0.12290834D  0.62746523D  0.62746523D  0.62746523D  0.62746523D  0.62746523D  0.62746523D  0.62746523D  0.62746523D  0.62746523D  0.668004189D  0.75750399D  0.86393864D  0.86393864D  0.69644712D  0.69644712D	KP(X)  0.28587583D  0.14332570D  0.86783103D-01  0.12955631D  0.12955631D  0.12955631D  0.1308778D  0.11985573D  0.11985573D  0.11985573D  0.22553126D-01  0.37670856D-01  0.37670856D-01  0.37670856D-01  0.37670856D-01  0.37670856D-01  0.37670856D-01  0.55229980D-01  0.55229980D-01  0.5531577D-01  0.5531577D-01  0.5531577D-01
WUNUWW44400000	
KP(X) -0.45620827D-01 -0.18088113D -0.15847093D -0.68882033D-01 -0.41371779D-01 -0.41371779D-01 -0.86468139D-01 -0.86468139D-01 -0.862700019D-01 -0.62700019D-01 -0.62700019D-01 -0.62463889D-01 -0.62463889D-01 -0.62463889D-01 -0.62463889D-01 -0.62463889D-01 -0.62463889D-01 -0.62463889D-01	KP(X)  0.65225007D-01  0.1678217D 00  0.16741295D-01  0.92352902D-01  -0.14238239D 00  -0.1614922D 00  0.12350400D 00  0.59574839D-01  -0.27538102D-01  -0.39366311D-01  0.83656989D-01  0.83656989D-01  0.83656989D-01  0.83656989D-01  0.83656989D-01  0.83656989D-01  0.83656989D-01  0.83656989D-01
KP(X)  1.1  1.5  -0.250392690-01  1.5  -0.726259160-01  2.7  -0.132479440-01  3.1  -0.152637030  3.5  -0.152637030  0.126913640-01  4.3  0.126913640-01  4.3  0.126913640-01  5.5  -0.152637030  0.26577220-01  5.5  -0.103757580  6.3  -0.103757580  6.3  -0.103757580  7.5  0.231857250-01  7.5  0.24837250-01  8.3  0.848830170-01  8.7  0.899251530-01  9.5  0.576209760-01	X
1.0	X

	51	
KP(X)  1 0.28635131D 00  1 0.15893331D 00  2 0.61848271D-01  3 -0.14282839D 00  1 0.48955262D-01  1 0.98022589D-01  2 -0.59781354D-01  3 -0.55051815D-01  5 -0.55051815D-01  5 -0.55051815D-01  5 -0.55051815D-01  5 -0.55051815D-01  6 -0.97813534D-01  7 -0.1073646D 00  1 0.920770360-01  2 0.83465810D-01  3 0.29023116D-01	KP(X)  1 00.28519124D 00  1 00.17139637D 00  5 00.35691953D-01  9 -0.15172968D 00  1 00.73165243D-01  1 00.73165243D-01  1 00.73165243D-01  2 -0.78353777D-01  1 00.83829356D-01  2 -0.30973428D-01  3 -0.44387199D-01  3 0.44387199D-02  7 28427598D-01	
KP(K) -0.84365414D-01 -0.86672692D-01 -0.12480633D 00 -12480633D 00 -13482113D-01 -0.13482113D-01 -0.34261335D-01 -0.52673521D-01	-0.215651050 -0.215651050 -0.215651050 -0.103732810 -0.103732810 -0.103732810 -0.103732810 -0.747003810-01 -0.747003810-01 -0.747003810-01 -0.747003810-01 -0.747003810-01 -0.74720-01 -0.651763280-01 -0.6527385360-01 -0.6527385360-01 -0.6527385360-01	
NN NW W444 NN O O O O C C C C C C C C C C C C C C		•
1.   -0.37993203D   00   1.   1.   -0.37993203D   00   1.   2.   3.   0.   17066872D   00   2.   7.   0.   17066872D   00   2.   7.   0.   17066872D   00   1706872D   00   17068	1.51 -0.37471175D 00 1.55 -0.93437468D-01 2.3 -0.93437468D-01 2.3 -0.15875970D-02 3.5 -0.15875970D-01 3.5 -0.1468325D-01 5.1 -0.57765537D-01 5.2 -0.10591652D-01 5.3 -0.59247661D-01 6.3 -0.59247661D-01 7.5 -0.82932652D-01 8.3 -0.593247661D-01 9.5 -0.92933432D-01 9.5 -0.929334304D-01 9.5 -0.929334304D-01	
X	X	

P= 9.30

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KP(X)

0.27806874D 00

-0.18639062D 00

0.18615881D 00

-0.15569337D 00

0.31046644D 00

0.31046444D 00

0.31046444D 00

0.31046444D 00

0.31046644D 00

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0.31046644D 00

0.31046444D 00

0.310464444D 00

0.31046
                 KP(X)

0.28241955D 00

0.18927579D 00

0.89423348D-02

-0.53708699D-01

0.94264914D-01

0.94264914D-01

0.94264914D-01

0.94264914D-01

0.95431554D-01

-0.10229122D 00

-0.41960818D-01

-0.55431554D-01

-0.51421849D-02

0.93684925D-01

0.93684925D-01

0.93684925D-01

0.93684925D-01

0.93684925D-01

0.93684920-01
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.22676907D 00
.14655759D 00
.55647387D-01
.15254274D 00
.70918194D-01
.87355578D-01
.12916289D 00
.56111141D-01
.4959439D-01
.10394190D 00
.57382849D-01
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.89822870D-01
.92746292D-01
                                     0.12099131D 00

0.12892219D 00

0.80494760D-01

0.80494760D-01

0.13139479D 00

0.55153970D-01

0.13139479D 00

0.79044481D-01

0.79044481D-01

-0.19387233D-01

-0.1938462D 00

-0.77868161D-01

-0.20400440D-01

0.37322123D-01

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0.24863528D-01

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0.24863528D-01
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.35346233D-01
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0.12676907D 0
0.156759D 0
0.152547387D-0
0.87355574D 0
0.1291618194D-0
0.56111141D-0
-0.10394190D 0
-0.10394190D 0
0.56309626D-0
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0.38866458D-0
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                            KP(X)
2099131D
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0.54576787D-01

0.12538475D 00

-0.12789354D 00

-0.12789354D 00

0.54855115D-01

0.54855115D-01

0.54855115D-01

-0.21654085D-01

-0.21654085D-01

-0.21654085D-01

-0.77766508D-01

0.47966508D-01

0.47966508D-01

0.48170873D-01

0.94427994D-01

0.94427994D-01
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0.27137161D-01
0.56361430D-01
0.74416739D-01
                                KP(X)

-0.36881290D 00

-0.11430630D 00

0.298112140-01

0.12733544D 00

0.292094970-01

0.292094970-01

0.998320050-01

0.998320050-01

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0.998320050-01

0.9983178840-01

0.9205684840-01

0.239953440-01

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0.239953440-01

0.239953440-01

0.239953400-01
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                                                0.85361841D-01

0.85361841D-01

0.17472689D 00

0.17472689D 00

0.15679197D-01

0.12638378D 00

0.36725816D-02

0.36725816D-02

0.36725816D-01

0.36725816D-01

0.36725810D-01

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0.36725810D-01

0.3673230D-01

0.965176066D-01

0.965176066D-01

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00000000
                                          KP(X)
100000000
10690795D
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9.50
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                                             --NNWWW44NNN00FFF00000
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lid

## Table 2

The zeros, value of the derivative at the zeros, bend points and value of the function at the bend points, to 8 decimal places.

ZEROS AND DERIVATIVES AT THE ZERO POINTS, BEND POINTS AND FUNCTION VALUES AT THE BEND POINTS OF THE CONICAL FUNCTION KP(X)

FUNCTION	0.0	0 * 0	0 • 0	0 • 0	-0.22648034D 00	-0.24217967D 00
BEND POINT	0 • 0	0	0.00	0	0.836484620 01	0.693225290 01
DERIVATIVE	-0.33233290D-01	-0.54660979D-01	-0.82381328D-01	-0.11639415D 00	-0.15658010D 00	-0.20276474D 00
P* 0.90 ZERO	0.81185516D 01	0.60666379D 01	0.48026857D 01	P# 1.20 0.39684929D 01	P# 1.30 0.33880547D 01	P# 1.40 0.29670618D 01

FUNCTION -0.25617995D 00	-0.26866050D 00	-0.27979230D 00	-0.28973276D 00	-0.29862366D 00	-0.30659097D 00 0.0	-0.31374566D 00 0.0
BEND POINT 0.58869363D 01	0.51005880D 01	0.449381720 01	0.40154535D 01	0.36313471D 01 0.0	0.33180137D 01 0.0	0.30588768D 01
DERIVATIVE -0.25475703D 00	-0.31237035D 00	-0.37543280D 00	-0.44379121D 00	-0.51731181D 00 0.38623858D-01	-0.59587916D 00 0.49560327D-01	-0.67939428D 00
P# 1.50 ZERO 0.26513731D 01	0	P# 1.70 0.22163913D 01 P# 1.80	0.20623778D 01 P# 1.90	0.19366371D 01 0.93792963D 01 P# 2.00	0.18325356D 01 0.80961006D 01 P# 2.10	0.17452963D 01 0.70902858D 01

FUNCTION -0.32018488D 00	-0.32599347D 00 0.0	-0.33124535D 00 0.16904983D 00	-0.33600489D 00 0.17560593D 00	-0.34032816D 00 0.18176394D 00	-0.34426403D 00 0.18754626D 00
BEND POINT 0.28419653D 01 0.0	0.26584580D 01	0.25017371D 01 0.89456896D 01	0.23667577D 01 0.79964115D 01	0.22496181D 01 0.72098682D 01	0.21472614D 01 0.65510057D 01
DERIVATIVE -0.76777267D 00 0.76540975D-01	-0.86094228D 00 0.92644161D-01	-0.95884168D 00 0.11050758D 00	-0.10614184D 01 0.13013285D 00	-0.11686277D 01 0.15151436D 00	-0.12804311D 01 0.17464106D 00
P# 2.20 ZERO 0.16714047D 01 0.62875583D 01	D# 2.30 0.16082264D 01 0.56367396D 01	D# 2.40 0.15537533D 01 0.51017156D 01	P# 2.50 0.15064316D 01 0.46564513D 01	P# 2.60 0.14650434D 01 0.42818157D 01	P# 2.70 0.14286221D 01 0.39635079D 01

FUNCTION -0.34785514D 00 0.19297519D 00	-0.35113877D 00 0.19807259D 00	-0.35414755D 00 0.20285957D 00 0.0	-0.35691010D 00 0.20735634D 00 0.0	-0.35945159D 00 0.21158211D 00 0.0	-0.36179415D 00 0.21555501D 00 0.0
D.20572641D 01 0.59936455D 01	0.19776844D 01 0.55179429D 01	0.19069506D 01 0.51086489D 01 0.0	0.18437798D 01 0.47538992D 01 0.0	0.17871154D 01 0.44443573D 01 0.0	0.17360810D 01 0.41725983D 01 0.0
-0.13967956D 01 0.19949785D 00	-0.15176928D 01	-0.16430982D 01 0.25432809D 00 -0.51158784D-01	-0.17729906D 01 0.28426073D 00 -0.59916706D-01	-0.19073514D 01 0.31584320D 00 -0.69528811D-01	-0.20461647D 01 0.34905388D 00 -0.80008615D-01
P# 2.80 ZERD 0.13963924D 01 0.36906706D 01	0.13677261D 01 0.34549445D 01	000	0.13191200D 01 0.30701263D 01 0.82531355D 01	3.2	0.12796716D 01 0.27714981D 01 0.69703527D 01

3.40 ZERO	DERIVATIVE	BEND POINT	FUNCTION
0.126267110 01 0.26465672D 01 0.64551353D 01	-0.21894163D 01 0.38387137D 00 -0.91366447D-01	0.16899445D 01 0.39326603D 01 0.97788474D 01	-0.36395735D 00 0.21929208D 00 -0.13724093D 00
3.50 0.12471942D 01 0.25347976D 01 0.60052002D 01	-0.23370939D 01 0.42027479D 00 -0.10360980D 00	0.16480901D 01 0.37197126D 01 0.89950237D 01	-0.36595848D 00 0.22280929D 00 -0.14118591D 00
12330622D 01 24343726D 01 56099563D 01	-0.24891868D 01 0.45824390D 00 -0.11674368D 00	0.16099967D 01 0.35298092D 01 0.83127951D 01	-0.36781285D 00 0.22612156D 00 -0.14496511D 00
3.70 12201221D 01 23437813D 01 52608629D 01	-0.26456854D 01 0.49775927D 00 -0.13077091D 00	0.15752209D 01 0.33597010D 01 0.77154542D 01	-0.36953410D 00 0.22924278D 00 -0.14858452D 00
.80 12082422D 01 22617582D 01 49509754D 01	-0.28065812D 01 0.53880232D 00 -0.14569243D 00	0.15433836D 01 0.32066940D 01 0.71895519D 01	-0.37113432D 00 0.23218590D 00 -0.15205027D 00

FUNCTION	-0.37262433D 00 0.23496293D 00 -0.15536850D 00	-0.37401378D 00 0.23758504D 00 -0.15854533D 00 0.0		-0.37531133D 00 0.24006258D 00 -0.16158679D 00 0.0		-0.37652472D 00 0.24240514D 00 -0.16449881D 00 0.0		-0.37766094D 00 0.24462164D 00 -0.16728715D 00 0.0
BEND POINT	0.15141586D 01 0.30685392D 01 0.67241727D 01	0.14872643D 01 0.29433471D 01 0.63103892D 01 0.0		0.14624564D 01 0.28295213D 01 0.59408478D 01 0.0		0.14395220D 01 0.27257050D 01 0.56094505D 01 0.0		0.14182751D 01 0.26307396D 01 0.53111082D 01 0.0
DERIVATIVE	-0.29718669D 01 0.58135542D 00 -0.16150760D 00	-0.31415358D 01 0.62540182D 00 -0.17821438D 00 0.54016339D-01		-0.33155820D 01 0.67092577D 00 -0.19580964D 00 0.60982540D-01	`	-0.34940003D 01 0.71791237D 00 -0.21428926D 00 0.68475865D-01		-0.36767860D 01 0.76634767D 00 -0.23364839D 00 0.76504177D-01
P# 3.90 ZERO	0.11973088D 01 0.21872370D 01 0.46746028D 01	0000	P* 4.10	0.11778984D 01 0.20572184D 01 0.42044052D 01 0.89432611D 01	P# 4.20	0.11692599D 01 0.20002906D 01 0.40034111D 01 0.83528320D 01	p* 4 .30	0.11612412D 01 0.19479624D 01 0.38213180D 01 0.78267631D 01

FUNCTION -0.37872627D 00 0.24672032D 00 -0.16995741D 00	-0.37972637D 00	-0.38066638D 00	-0.38155092D 00	-0.38238421D 00
	0.24870882D 00	0.25059425D 00	0.25238316D 00	0.25408166D 00
	-0.17251501D 00	-0.17496516D 00	-0.17731287D 00	-0.17956294D 00
	0.12119042D 00	0.12380106D 00	0.12632739D 00	0.12877175D 00
Dend Point 0.13985521D 01 0.25436306D 01 0.50415486D 01 0.0	0.13802094D 01 0.24635202D 01 0.47971652D 01 0.95654168D 01	0.13631195D 01 0.23896658D 01 0.45748974D 01 0.89782353D 01	0.23214216D 01 0.43721347D 01 0.84501074D 01	0.13322605D 01 0.22582235D 01 0.41866404D 01 0.79734305D 01
DERIVATIVE -0.38639348D 01 0.81621854D 00 -0.25388153D 00 0.85074128D-01	-0.40554431D 01	-0.42513073D 01	-0.44515244D 01	-0.46560917D 01
	0.86751269D 00	0.92021863D 00	0.97432557D 00	0.10298235D 01
	-0.27498274D 00	-0.29694568D 00	-0.31976374D 00	-0.34343016D 00
	0.94191243D-01	0.10386001D 00	0.11408397D 00	0.12486580D 00
P* 4.40 ZERO 0.11537840D 01 0.18997425D 01 0.36558013D 01 0.73560800D 01	P# 4.50 0.11468366D 01 0.18552042D 01 0.35048840D 01 0.69333026D 01	0.11403532D 01 0.18139755D 01 0.33668759D 01 0.65521592D 01	P# 4.70 0.11342932D 01 0.17757307D 01 0.32403254D 01 0.62073625D 01	P# 4.80 0.11286202D 01 0.17401836D 01 0.31239803D 01 0.58944318D 01

FUNCTION -0.38317008D 00 0.25569541D 00 -0.18171*00D 00	-0.38391199D 00	-0.38461312D 00	-0.38527637D 00	-0.38590438D 00
	0.25722965D 00	0.25868926D 00	0.26007877D 00	0.26140239D 00
	-0.18378841D 00	-0.18577240D 00	-0.18767595D 00	-0.18950287D 00
	0.13342408D 00	0.13563684D 00	0.13777720D 00	0.13984750D 00
Dela point 0.13183016D 01 0.21995776D 01 0.40164889D 01 0.75417721D 01	0.13052136D 01 0.21450493D 01 0.38600162D 01 0.71496622D 01 0.0	0.12929247D 01 0.20942554D 01 0.37157778D 01 0.67924285D 01	0.12813707D 01 0.20468564D 01 0.35825154D 01 0.64660616D 01	0.12704935D 01 0.20025512D 01 0.34591288D 01 0.61671074D 01
DERIVATIVE -0.48650066D 01 0.10867029D 01 -0.36793803D 00 0.13620740D 00	-0.50782669D 01	-0.52958705D 01	-0.55178155D 01	-0.57441003D 01
	0.11449550D 01	0.12045717D 01	0.12655452D 01	0.13278684D 01
	-0.39328043D 00	-0.41945041D 00	-0.44644109D 00	-0.47424564D 00
	0.14810994D 00	0.16057399D 00	0.17359952D 00	0.18718603D 00
	-0.57163228D-01	-0.63077080D-01	-0.69356265D-01	-0.76005714D-01
P# 4.90 ZERO 0.11233017D 01 0.17070821D 01 0.30167554D 01 0.56095512D 01	P* 5.00 0.11183086D 01 0.16762030D 01 0.29177071D 01 0.53494559D 01 0.99634819D 01	D# 5.10 0.11136147D 01 0.16473484D 01 0.28260113D 01 0.51113399D 01 0.93982968D 01	0.11091965D 01 0.16203422D 01 0.27409460D 01 0.48927818D 01 0.88854669D 01	0.11050327D 01 0.15950273D 01 0.26618762D 01 0.46916833D 01 0.84187767D 01

FUNCTION	-0.38905484D 00	-0.38949391D 00	-0.38991242D 00	-0.39031162D 00
	0.26816181D 00	0.26912000D 00	0.27003708D 00	0.27091529D 00
	-0.19904623D 00	-0.20042881D 00	-0.20175918D 00	-0.20303974D 00
	0.15092570D 00	0.15256831D 00	0.15415815D 00	0.15569710D 00
	-0.11523547D 00	-0.11697893D 00	-0.11867622D 00	-0.12032852D 00
BEND POINT	0.12168366D 01	0.12094772D 01	0.12024890D 01	0.11958473D 01
	0.17884953D 01	0.17597204D 01	0.17325275D 01	0.17068010D 01
	0.28793948D 01	0.28035804D 01	0.27324054D 01	0.26654926D 01
	0.48061632D 01	0.46337871D 01	0.44731977D 01	0.43233353D 01
	0.81281164D 01	0.77636643D 01	0.74269234D 01	0.71151655D 01
DERIVATIVE	-0.71928641D 01 0.17298043D 01 -0.65779917D 00 0.28040723D 00 -0.12388666D 00	-0.74494898D 01 0.18014100D 01 -0.69112162D 00 0.29787516D 00 -0.13322236D 00	-0.77104458D 01 0.18743225D 01 -0.72520877D 00 0.31588801D 00 -0.14294996D 00 0.65546842D-01	-0.79757313D 01 0.19485378D 01 -0.76005528D 00 0.33444288D 00 -0.15307040D 00 0.71030089D-01
P# 5.90	0.10844277D 01 0.14717022D 01 0.22860539D 01 0.37640102D 01 0.63346205D 01	0.10815931D 01 0.14549894D 01 0.22363364D 01 0.36449088D 01 0.60757904D 01	0.10788996D 01 0.14391650D 01 0.21895319D 01 0.35335877D 01 0.58357925D 01 0.97204038D 01	0.10763379D 01 0.14241662D 01 0.21454129D 01 0.34293747D 01 0.56128393D 01 0.92685189D 01

FUNCTION -0.39069267D 00 0.27175672D 00 -0.20427276D 00 0.15718700D 00 -0.12193694D 00	-0.39105665D 00 0.27256334D 00 -0.20546040D 00 0.15862964D 00 -0.12350264D 00	-0.39140454D 00 0.27333698D 00 -0.20660472D 00 0.16002672D 00 -0.12502674D 00	-0.39173728D 00 0.27407937D 00 -0.20770766D 00 0.16137992D 00 -0.12651036D 00 0.99508369D-01
Dend Point 0.11895294D 01 0.16824354D 01 0.41832564D 01 0.68259864D 01 0.0	0.11835145D 01 0.16593350D 01 0.25431288D 01 0.40521190D 01 0.65572614D 01	0.11777834D 01 0.16374123D 01 0.24870953D 01 0.39291698D 01 0.63071075D 01	0.11723184D 01 0.16165872D 01 0.24341515D 01 0.38137332D 01 0.60738517D 01 0.97363343D 01
DERIVATIVE -0.82453456D 01 0.20240522D 01 -0.79565602D 00 0.35353677D 00 -0.16358430D 00 0.76788564D-01	-0.85192881D 01 0.21008623D 01 -0.83200603D 00 0.37316661D 00 -0.17449201D 00 0.82825268D-01	-0.87975580D 01 0.21789648D 01 -0.86910054D 00 0.39332928D 00 -0.18579361D 00 0.89142915D-01	-0.90801548D 01 0.22583566D 01 -0.90693498D 00 0.41402161D 00 -0.19748898D 00 0.95743943D-01
P# 6.30 ZERD 0.10738994D 01 0.14099357D 01 0.33316687D 01 0.54053489D 01 0.88512474D 01	D# 6.40 0.10715765D 01 0.20644292D 01 0.32399313D 01 0.52119179D 01 0.84651733D 01	0.10693619D 01 0.13835748D 01 0.20272100D 01 0.31536787D 01 0.50312978D 01 0.81072863D 01	0.10672488D 01 0.13713527D 01 0.19919632D 01 0.30724754D 01 0.48623750D 01 0.77749254D 01

FUNCTION	-0.39205572D 00 0.27479211D 00 -0.20877106D 00 0.16269081D 00 -0.12795460D 00 0.10098816D 00		-0.39236067D 00 0.27547673D 00 -0.20979668D 00 0.16396095D 00 -0.12936055D 00		-0.39265287D 00 0.27613464D 00 -0.21078618D 00 0.16519182D 00 -0.13072927D 00		-0.39293302D 00 0.276751BD 00 -0.21174114D 00 0.16638484D 00 -0.13206182D 00
BEND POINT	0.11671032D 01 0.15967866D 01 0.23840704D 01 0.37052021D 01 0.58560039D 01 0.93180818D 01		0.11621227D 01 0.15779434D 01 0.23366454D 01 0.36030296D 01 0.56522333D 01 0.89294790D 01		0.11573629D 01 0.15599957D 01 0.22916886D 01 0.35067218D 01 0.54613490D 01 0.85678136D 01		0.11528109D 01 0.15428871D 01 0.22490285D 01 0.34158324D 01 0.52822826D 01 0.82306721D 01
DERIVATIVE	-0.93670780D 01 0.23390349D 01 -0.94550492D 00 0.43524043D 00 -0.20957776D 00 0.10263053D 00		-0.96583271D 01 0.24209969D 01 -0.98480616D 00 0.45698255D 00 -0.22205942D 00 0.10980460D 00		-0.99539016D 01 0.25042403D 01 -0.10248346D 01 0.47924477D 00 -0.23493325D 00 0.11726783D 00		-0.10253801D 02 0.25887624D 01 -0.10655864D 01 0.50202392D 00 -0.24819837D 00 0.12502170D 00
P# 6.70 ZERO	0.10652313D 01 0.13597145D 01 0.19585491D 01 0.29959288D 01 0.47041540D 01 0.74657324D 01	P* 6.80	0.10633036D 01 0.13486231D 01 0.19268407D 01 0.29236836D 01 0.45557428D 01	06.90	0.10614604D 01 0.13380443D 01 0.18967216D 01 0.28554181D 01 0.44163400D 01	7 * 00	0.10596968D 01 0.13279468D 01 0.18680853D 01 0.27908404D 01 0.42852244D 01

FUNCTION -0.39320176D 00 0.27737560D 00 -0.21266307D 00 0.16754140D 00 0.10658779D 00 0.000000000000000000000000000000000	-0.39345971D 00 0.27796110D 00 -0.21355338D 00 0.16866280D 00 -0.13462248D 00 0.10791085D 00	-0.39370742D 00 0.27852476D 00 -0.21441343D 00 0.16975033D 00 -0.13585257D 00 0.10920440D 00	-0.39394544D 00 0.27906765D 00 -0.21524451D 00 0.17080521D 00 -0.13705047D 00 0.11046908D 00
Dend Point 0.11484548D 01 0.22085085D 01 0.33299565D 01 0.51140736D 01 0.79159016D 01	0.11442834D 01 0.15109815D 01 0.32487267D 01 0.49558568D 01 0.76215773D 01	0.11402863D 01 0.21333271D 01 0.31718088D 01 0.48068511D 01 0.73459745D 01	0.11364539D 01 0.14818551D 01 0.20984134D 01 0.30988982D 01 0.46663500D 01 0.70875444D 01
DERIVATIVE -0.10558025D 02 0.26745612D 01 -0.11070578D 01 0.52531686D 00 -0.26185380D 00 0.13306746D 00	-0.10866573D 02 0.27616345D 01 -0.11492452D 01 0.54912045D 00 -0.27589840D 00 0.14140617D 00	-0.11179445D 02 0.28499804D 01 -0.11921452D 01 0.57343163D 00 -0.29033095D 00 0.15003868D 00	-0.11496641D 02 0.29395968D 01 0.12357545D 01 0.30515011D 00 0.15896569D 00
P# 7.10 ZERO 0.10580083D 01 0.13183013D 01 0.18408344D 01 0.27296850D 01 0.41617456D 01 0.64220011D 01 0.99602471D 01	7.20 0.10563907D 01 0.13090809D 01 0.18148793D 01 0.26717100D 01 0.62013887D 01 0.65013887D 01	0.10548401D 01 0.13002609D 01 0.17901377D 01 0.39354020D 01 0.59942857D 01 0.59942857D 01	0.10533527D 01 0.12918179D 01 0.17665337D 01 0.25644378D 01 0.38315218D 01 0.57996097D 01

				69
FUNCTION	-0.39417425D 00 0.27959076D 00 -0.21604784D 00 0.17182861D 00 0.11170550D 00	-0.39439432D 00 0.28009499D 00 -0.21682459D 00 0.11282166D 00 -0.13935340D 00 0.11291427D 00	-0.39460609D 00 0.28058125D 00 -0.17378545D 00 -0.14046022D 00 0.11409601D 00 -0.92882262D-01	-0.39480997D 00 0.28105035D 00 0.21830270D 00 0.17472102D 00 -0.14153846D 00 0.11525134D 00
BEND POINT	0.11327772D 01 0.14682329D 01 0.20651333D 01 0.30297169D 01 0.68448935D 01	0.11292480D 01 0.14551900D 01 0.20333844D 01 0.29640108D 01 0.44083601D 01 0.66167653D 01 0.99719862D 01	0.11258583D 01 0.14426935D 01 0.20030726D 01 0.42897614D 01 0.64020246D 01 0.95948551D 01	0.11226008D 01 0.14307129D 01 0.19741108D 01 0.28421124D 01 0.41774361D 01 0.61996442D 01
DERIVATIVE	-0.11818159D 02 0.30304821D 01 -0.12800699D 01 0.62356463D 00 -0.32035448D 00 0.16818773D 00 -0.89110668D-01	-0.12144001D 02 0.31226346D 01 -0.13250885D 01 0.64938052D 00 -0.33594256D 00 0.17770518D 00	-0.12474165D 02 0.32160527D 01 -0.13708073D 01 0.67569214D 00 -0.35191282D 00 0.18751827D 00	-0.12808652D 02 0.33107349D 01 -0.14172236D 01 0.70249668D 00 -0.36826364D 00 0.19762709D 00
P# 7.50 ZERO	0.10519253D 01 0.12837307D 01 0.17439973D 01 0.25147542D 01 0.37332360D 01 0.56163833D 01	0.10505545D 01 0.12759793D 01 0.12759793D 01 0.24674746D 01 0.36401448D 01 0.54437231D 01	0.10492375D 01 0.12685451D 01 0.17018743D 01 0.35518836D 01 0.52808285D 01 0.52808285D 01	0.10479715D 01 0.12614109D 01 0.16821726D 01 0.23795184D 01 0.34681195D 01 0.51269732D 01

FUNCTION -0.39500634D 00 0.28150307D 00 -0.21900611D 00 0.17562938D 00 -0.14258894D 00 -0.14258894D 00	-0.39519556D 00 0.28194016D 00 -0.21968705D 00 0.17651149D 00 -0.14361248D 00 0.11748512D 00 -0.96345727D-01	-0.39537798D 00 0.28236231D 00 -0.22034643D 00 0.17736827D 00 -0.14460989D 00 0.11856477D 00 -0.97456429D-01	-0.39555392D 00 0.28277018D 00 -0.22098511D 00 0.17820061D 00 -0.14558194D 00 0.11962036D 00 -0.98545842D-01
0.11194688D 01 0.14192199D 01 0.19464185D 01 0.27855106D 01 0.40709446D 01 0.60086927D 01 0.89092527D 01	0.11164559D 01 0.14081879D 01 0.19199214D 01 0.39698854D 01 0.58283241D 01	0.11135560D 01 0.13975924D 01 0.18945503D 01 0.38738911D 01 0.56577686D 01 0.83034577D 01	0.11107636D 01 0.13874103D 01 0.18702412D 01 0.26309696D 01 0.37826246D 01 0.54963249D 01 0.80267020D 01
DERIVATIVE -0.13147460D 02 0.34066798D 01 -0.14643347D 01 0.72979135D 00 -0.38499340D 00 0.20803163D 00 -0.11359866D 00	-0.13490591D 02 0.35038861D 01 -0.15121381D 01 0.75757347D 00 -0.40210041D 00 0.21873176D 00	-0.13838043D 02 0.36023525D 01 -0.15606313D 01 0.78584040D 00 -0.41958296D 00 0.22972723D 00	-0.14189817D 02 0.37020778D 01 -0.16098121D 01 0.81458956D 00 -0.43743933D 00 0.24101770D 00 -0.13433696D 00
P# 7.90 ZERU 0.10467537D 01 0.12545605D 01 0.23385670D 01 0.33885477D 01 0.49814968D 01 0.73726631D 01	0.10455820D 01 0.12479791D 01 0.12479791D 01 0.22994681D 01 0.33128890D 01 0.48437977D 01 0.71313297D 01	P# 8.10 0.10444538D 01 0.12416525D 01 0.22621094D 01 0.32408870D 01 0.47133275D 01	0.10433671D 01 0.12355676D 01 0.1212760D 01 0.22263873D 01 0.31723063D 01 0.45895849D 01 0.66888113D 01

FUNCTION -0.39572367D 00 0.28316440D 00 -0.22160391D 00 0.17900937D 00 -0.14652936D 00 0.12065246D 00 -0.99614333D-01	-0.39588753D 00 0.28354555D 00 -0.22220362D 00 0.17979537D 00 -0.14745290D 00 -0.12166163D 00 -0.10066227D 00	-0.39604576D 00 0.28391421D 00 -0.22278499D 00 0.18055940D 00 -0.14835326D 00 0.12264843D 00 -0.10169002D 00	-0.28427090D 00 -0.22334872D 00 -0.18130221D 00 -0.12361339D 00 -0.10269796D 00
BEND POINT 0.11080734D 01 0.13776201D 01 0.25840322D 01 0.36957766D 01 0.53433527D 01 0.77656336D 01	0.11054805D 01 0.13682017D 01 0.18245750D 01 0.25391559D 01 0.36130627D 01 0.51982667D 01 0.75190952D 01	0.11029801D 01 0.13591363D 01 0.24962195D 01 0.35342212D 01 0.50605312D 01 0.72860331D 01	0.11005680D 01 0.13504063D 01 0.17824947D 01 0.34551108D 01 0.49296549D 01 0.70654861D 01
DERIVATIVE -0.14545912D 02 0.38030609D 01 -0.16596781D 01 0.84381843D 00 -0.45566775D 00 0.25260276D 00 -0.14170730D 00 0.79933901D-01	-0.14906329D 02 0.39053007D 01 -0.17102274D 01 0.87352457D 00 -0.47426649D 00 -0.26448189D 00 -0.14930773D 00	-0.15271066D 02 0.40087962D 01 0.90370561D 00 -0.49323377D 00 0.27665452D 00 -0.15713888D 00	-0.15640124D 02 0.41135464D 01 -0.18133672D 01 0.93435921D 00 -0.28911*00D 00 -0.16520127D 00 0.94973469D-01
D.10423199D 01 0.12297122D 01 0.15953157D 01 0.21922057D 01 0.31069301D 01 0.44721114D 01 0.64856805D 01 0.94395744D 01	P# 8.40 0.10413102D 01 0.12240747D 01 0.21594757D 01 0.30445586D 01 0.43604866D 01 0.62934818D 01 0.62934818D 01	0.10403364D 01 0.12186444D 01 0.15652522D 01 0.21281147D 01 0.42543250D 01 0.61114442D 01	0.10393967D 01 0.12510847D 01 0.20980458D 01 0.41532722D 01 0.59388629D 01 0.85258185D 01

BEND POINT    0.354895	
DERIVATIVE  DERIVA	0.10359487D 0.11942666D 0.14995185D 0.19893363D 0.27240180D 0.37939964D 0.53309747D
DERIVATIVE  0.109824000 01 0.109824000 01 0.109824000 01 0.1034955050 01 0.1034955050 01 0.1034955050 01 0.1034955050 01 0.1034955050 01 0.1034955050 01 0.103495050 01 0.103495050 01 0.103495050 01 0.103495050 01 0.103495050 01 0.103495050 01 0.103495050 01 0.103495050 01 0.103495050 01 0.103495050 01 0.103495050 01 0.103495050 01 0.103495050 01 0.103495050 01 0.10349500 01 0.10349500 01 0.10349500 01 0.10349500 01 0.10349500 01 0.10349500 01 0.10349500 01 0.10349500 01 0.10349500 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910 01 0.1034910	5500 <b>5</b> 50
DEND POINT	-0.1/159562D 02 0.45450776D 01 -0.20277602D 01 0.10616552D 01 -0.59353665D 00 0.34189546D 00 -0.19977043D 00 0.11754036D 00
FUNCTION 0.39634633D 00 0.22389551D 00 0.18202454D 00 0.12455705D 00 0.12455705D 00 0.28442599D 00 0.28442599D 00 0.28442599D 00 0.18272708D 00 0.18272708D 00 0.18272708D 00 0.18241050D 00 0.18341050D 00 0.1844250D 00 0.1844250D 00 0.1844250D 00 0.1844250D 00 0.184420D 00	0.13185240D 01 0.17076509D 01 0.23069764D 01 0.31902671D 01 0.44662459D 01 0.62919715D 01
7 000000H 000000H 00	0.285587660 0.28587660 0.184075450 0.152530570 0.127265340 0.893336110-01

FUNCTION -0.39689019D 00 0.28589157D 00 -0.22592570D 00 -0.15330550D 00 -0.12812886D 00 -0.10745328D 00 0.90266141D-01	-0.39701539D 00 0.28618618D 00 -0.22639690D 00 0.18535241D 00 -0.15406162D 00 0.12897358D 00 -0.10835000D 00 0.91182985D-01	-0.39713664D 00 0.28647186D 00 -0.22685464D 00 0.18596558D 00 -0.15479949D 00 0.12979994D 00 0.92084375D-01	-0.39725411D 00 0.28674895D 00 -0.22729940D 00 0.18656263D 00 -0.15551962D 00 0.13060841D 00 -0.11009194D 00 0.92970544D-01
DEND POINT 0.10896945D 01 0.13112421D 01 0.22735821D 01 0.31301860D 01 0.43635776D 01 0.61222227D 01 0.86178130D 01	0.10877330D 01 0.13042101D 01 0.16742789D 01 0.22414897D 01 0.30726247D 01 0.42655443D 01 0.59607078D 01 0.83577112D 01	0.10858354D 01 0.12974167D 01 0.16584907D 01 0.22106311D 01 0.30174418D 01 0.41718681D 01 0.58069018D 01 0.581109020D 01	0.10839989D 01 0.12908511D 01 0.16432627D 01 0.29645061D 01 0.40822920D 01 0.56603206D 01 0.78765008D 01
DERIVATIVE -0.17550222D 02 0.46560880D 01 -0.20830388D 01 0.10946391D 01 -0.61467830D 00 0.35581354D 00 -0.20899365D 00	-0.17945202D 02 0.47683489D 01 -0.11280829D 01 0.37001951D 00 -0.21844939D 00 0.12992073D 00	-0.18344502D 02 0.48818590D 01 -0.21956014D 01 0.11619848D 01 -0.65802895D 00 0.38451240D 00 -0.22813760D 00 -0.22813760D 00 -0.81848934D-01	-0.18748122D 02 0.449966177D 01 -0.22528827D 01 0.11963459D 00 0.39929121D 00 -0.23805816D 00 -0.4304133D 00
P# 9.10 ZERO 0.10351576D 01 0.11898868D 01 0.14877792D 01 0.19647504D 01 0.26782208D 01 0.37140740D 01 0.51969981D 01 0.73054961D 01	0.10343924D 01 0.11856548D 01 0.14764567D 01 0.26342836D 01 0.36376442D 01 0.50693166D 01 0.70978725D 01 0.99621628D 01	9.30 0.10336519D 01 0.11815639D 01 0.19183219D 01 0.25921052D 01 0.35645036D 01 0.49475395D 01 0.69005416D 01 0.96484994D 01	0.10329352D 01 0.11776080D 01 0.14549836D 01 0.25953878D 01 0.25515912D 01 0.34944636D 01 0.48313060D 01 0.67128356D 01

FUNCTION -0.39736795D 00 0.28701781D 00 -0.22773165D 00 -0.18714409D 00 -0.1562254D 00 -0.13139943D 00 -0.11093786D 00	-0.39747832D 00 0.28727873D 00 -0.22815185D 00 0.18771046D 00 -0.15690874D 00 0.13217344D 00 -0.11176752D 00 0.94698154D-01	-0.39758534D 00 0.28753204D 00 -0.22856043D 00 0.18826224D 00 -0.13757871D 00 0.13293086D 00 -0.11258126D 00 0.95540066D-01
BEND POINT 0.10822209D 01 0.12845031D 01 0.21523648D 01 0.29136950D 01 0.39965773D 01 0.55205171D 01 0.76536942D 01	0.10804990D 01 0.12783531D 01 0.2124841BD 01 0.2864845D 01 0.39145029D 01 0.53870782D 01 0.74417338D 01	0.10788307D 01 0.12724219D 01 0.16006809D 01 0.28179980D 01 0.5259628D 01 0.5259628D 01 0.99863347D 01
DERIVATIVE -0.19156062D 02 0.51126243D 01 -0.23108287D 01 -0.70279158D 00 0.41435494D 00 -0.24821090D 00 0.14988086D 00	-0.19568321D 02 0.52298783D 01 -0.23694383D 01 0.12664210D 01 -0.72569829D 00 0.42970256D 00 -0.2585956D 00 -0.2585956D 00	-0.19984900D 02 0.53483791D 01 -0.24287101D 01 0.13021374D 01 -0.74895312D 00 0.44533301D 00 -0.26921198D 00 0.16412108D 00
2411D 01 7810D 01 7970D 01 2511D 01 6536D 01 3490D 01 1394D 01 0962D 01	5688D 01 0775D 01 9547D 01 2102D 01 9965D 01 1585D 01 2385D 01	9174D 01 4920D 01 4411D 01 2163D 01 1839D 01 2545D 01 6487D 01 5510D 01
0.050 0.1173 0.1173 0.1173 0.1875 0.3427 0.5534 0.9069	0.00 1170 0.1170 0.1170 0.1854 0.00 1434 0.00 6363	9.70 0.11030 0.11630 0.11656 0.2439 0.3301 0.65201 0.65201

FUNCTION -0.39768915D 00 0.28777802D 00 -0.22895780D 00 0.18879990D 00 -0.15823292D 00 -0.13367209D 00 -0.11337942D 00 -0.81997371D-01	-0.39778988D 00 0.28801694D 00 -0.22934436D 00 0.18932389D 00 -0.15887181D 00 -0.13439755D 00 -0.11416231D 00 0.97181277D-01	-0.39788765D 00 0.28824907D 00 -0.22972048D 00 0.18983466D 00 -0.15949583D 00 -0.13510763D 00 -0.11493026D 00 0.97981044D-01
BEND POINT 0.10772140D 01 0.12666712D 01 0.27729060D 01 0.27729060D 01 0.37604659D 01 0.51377931D 01 0.70476456D 01 0.96879771D 01	0.10756466D 01 0.12611026D 01 0.2746456D 01 0.27295255D 01 0.36881337D 01 0.50212642D 01 0.68642933D 01	0.10741267D 01 0.12557085D 01 0.20243000D 01 0.26877694D 01 0.36187004D 01 0.49097301D 01 0.56893293D 01
DERIVATIVE -0.20405799D 02 0.54681262D 01 -0.24886430D 01 0.13383032D 01 -0.77255452D 00 0.46124523D 00 -0.28005971D 00 0.17152263D 00	-0.20831016D 02 0.55891191D 01 -0.25492359D 01 0.13749169D 01 -0.79650092D 00 0.47743815D 00 -0.29113843D 00 -0.17911227D 00	-0.21260554D 02 0.57113572D 01 -0.26104876D 01 0.14119769D 01 -0.82079081D 00 0.49391068D 00 -0.30244774D 00 0.18689029D 00 -0.11603558D 00
P# 9.80 ZERO 0.10302859D 01 0.11630197D 01 0.14162416D 01 0.24045029D 01 0.32419814D 01 0.44154867D 01 0.60466511D 01	0.10296737D 01 0.11596558D 01 0.17979342D 01 0.1797937D 01 0.31850454D 01 0.43224253D 01 0.58987383D 01 0.80740554D 01	0.10290798D 01 0.11563957D 01 0.13987298D 01 0.23389112D 01 0.31303231D 01 0.42332348D 01 0.57573975D 01

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## MANUSCRIPT REPORT SERIES



No. 23

Effect of a travelling atmospheric pressure disturbance on a narrow lake with a depth-discontinuity

T.S. Murty



# Manuscript Report Series No.23

# DISTURBANCE ON A NARROW LAKE WITH A DEPTH-DISCONTINUITY

T.S. Murty



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#### 0. Abstract

The effect of a travelling atmospheric pressure disturbance on the water level of a narrow lake with a depth-discontinuity has been studied using the method of characteristics. Both semi-infinite stress-bands and finite stress-bands were considered. The set-up defined as the difference in the water level between the right and left boundaries does not become periodic any time after the disturbance crosses the lake. For semi-infinite stress-bands the water level at the left side is predominantly negative while it is positive on the right side and it becomes both positive and negative for finite stress-band widths.

#### 1. Introduction

The calculations reported here have been originally made in connection with evaluating the role of the Saguenay River Barrier on the water level changes in the St. Lawrence River when atmospheric disturbances travel over the river. The paper was presented at the 51st Annual Meeting of the American Geophysical Union held at Washington, D.C. in April 1970 and a condensed version is published (Murty, 1971). Neither in the presentation of the paper nor in the paper published are the details of the calculations dealt with. These are discussed in this manuscript report.

The effect of a travelling pressure disturbance on the water level in a non-rotating one-dimensional lake of uniform depth has been studied by Harris (1957). His solutions are valid only for very short time intervals after the disturbance crosses

the lake. Rao (1967) treated the same problem and also the case of a rectangular bay of uniform depth (1969) making use of the method of characteristics. For both these cases he showed that the maximum value of the set-up (set-up is defined as the difference of water level at the right and left boundaries of the lake in the case of an atmospheric disturbance travelling from left to right) is periodic some time after the passage of the disturbance. In fact, for the uniform depth cases this periodicity could be proved a priori from the characteristics' equations even before obtaining solutions for the height fields  $h_{\rm T}$  and  $h_{\rm P}$  at the left and right boundaries.

In the present study I consider the case of a lake with a depth-discontinuity. An example of such a discontinuity in nature is the so-called Saguenay River Barrier (the region where the Saguenay River joins the St. Lawrence River) where the depth falls approximately by a factor of four in a distance of a few kilometers. However, our interest here is not a particular case (in which case direct numerical integration could have been used) but to study this problem generally keeping other conditions as simple as possible. For this reason the problem is solved analytically through the use of the method of characteristics.

An atmospheric disturbance travelling over a body of water imparts energy to the water through pressure-gradient and wind-stress forces. For a small-scale system such as that considered here, the pressure-gradient force can be neglected and thus only the wind-stress is considered. The rate at which

the water absorbs energy from the atmospheric disturbance depends upon how close the propagation speed of the disturbance is to the propagation speed of long gravity waves in the lake because resonant coupling is possible when they are equal or approximately so. The relative sizes of the atmospheric disturbance and the lake are also important because the finite size of the lake has to be taken into account if the size of the atmospheric disturbance is comparable to or greater than the size of the lake. This means wave reflections at the left and right boundaries of the lake as well as at the depth-discontinuity must be taken into account.

Certain simplifying assumptions have been made to keep the mathematical problem tractable and adapt the method of characteristics to obtain the solution. For convenience, the depth-discontinuity is assumed to be at mid-length of the lake. The two portions of the lake on either side of this discontinuity are assumed to have uniform depths D1 and D2. Hydrostatic approximation has been made for the pressure field and the nonlinear advective terms are suppressed. Platzman (1958) gave a justification for these approximations based on order of magnitude considerations. Though in principle the method of characteristics could be used with a two-dimensional lake (in which case integration must be performed along characteristic surface rather than along characteristic lines), in practice the calculation is complicated. For this reason a narrow lake is considered here in which the earth's rotation effects could be ignored and no transverse motions are permitted. Also, since

our interest is in the transient aspects of the motion and not in the final steady state, bottom frictional effects also are ignored.

2. Formulation of the Problem and the Method of Solution

Consider a lake of length L containing water of uniform density  $\rho$ . The origin of a right-handed cartesian co-ordinate system with the x-axis along the length of the lake and the Z-axis pointing upward is placed on the left boundary at the mean water level. Then the depth-discontinuity is at x=L/2 and the right boundary is at x=L. The vertically integrated equations of motion and continuity for regions to the left of the depth-discontinuity (identified through the subscript 1) and to the right of the discontinuity (identified through the subscript 2) are

$$\frac{\partial M_{1,2}}{\partial t} = -gD_{1,2} \frac{\partial h_{1,2}}{\partial x} + R$$
 (1)

$$\frac{\partial h_{1,2}}{\partial t} = \frac{\partial M_{1,2}}{\partial x} \tag{2}$$

where

$$M_{1,2} \equiv {h_{1,2} \choose -D_{1,2}} U_{1,2} dz$$
 (3)

Here M is the volume transport through a vertical section, U is the horizontal component of velocity along the length of the lake, h is the deviation of the water level from the equilibrium position and R is a forcing function such that

$$R \equiv \frac{\tau}{\rho} \tag{4}$$

where  $\tau$  is the wind-stress which is assumed to be known and R is the quantity to be prescribed in the actual calculation.

The boundary conditions of the problem are prescribed through assuming perfect reflections at the left and right boundaries of the lake. Thus

$$M_1 = 0 \text{ at } x = 0$$

$$M_2 = 0 \text{ at } x = L$$
(5)

At the depth-discontinuity, continuity of M and h is required.

$$M_1 = M_2$$

$$\text{at } x = \frac{L}{2}$$

$$h_1 = h_2$$
(6)

Initially the lake is assumed to be at rest.

$$M_1 = M_2 = 0$$
  
at t = 0 (7)  
 $h_1 = h_2 = 0$ 

Let  $C_1$  and  $C_2$  be the speed of long gravity waves in regions 1 and 2 such that

$$C_{1,2}^2 \equiv g D_{1,2}$$
 (8)

Since this is a non-dispersive system (with the earth's rotation ignored) all the gravity waves propagate with the same speed  $C_1$  in region 1 and with the same speed  $C_2$  in region 2.

Addition and subtraction of equations (1) and (2) and then use of (8) gives four equations for the four unknowns  $M_1$ ,  $M_2$ ,  $h_1$  and  $h_2$  which can be expressed in the following compact

form:

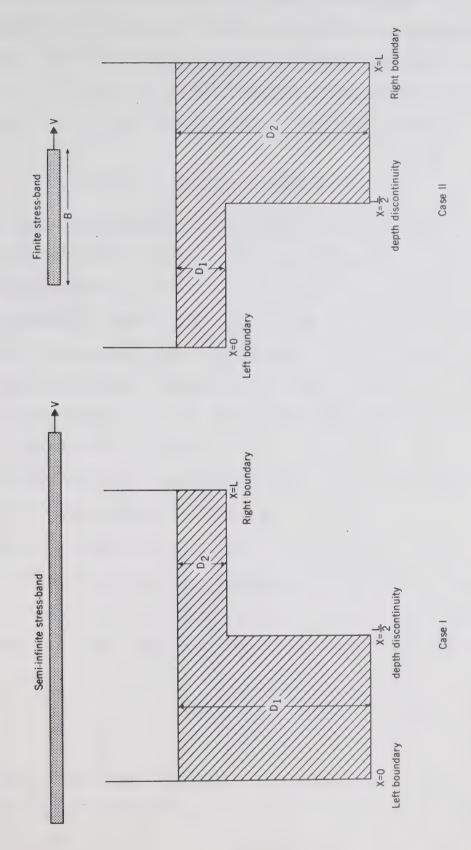
$$\frac{d}{dt} (M_{1,2} \pm C_{1,2} h_{1,2}) = R$$

for

$$\frac{dx}{dt} = \pm C_{1,2}$$

This equation states that the quantity  $M_{1,2} \pm C_{1,2} \cdot h_{1,2}$  is constant along the characteristics  $dx/dt = \pm C_{1,2} \cdot h_{1,2}$ . Since both the regions of the lake have uniform depths individually, the characteristics in both regions are straight lines with slopes  $C_1$  and  $C_2 \cdot h_1$ . Had the nonlinear term  $U \cdot \frac{\partial U}{\partial x}$  been included in the equation of motion, the method of characteristics would still be applicable but the characteristics would no longer be straight lines. In fact, in such a case they have to be determined as a part of the solution with each integration. Through a scale analysis it can be shown that the omission of the nonlinear term is a minor drawback of the model compared to the omission of bottom friction.

The calculations were performed both for semi-infinite and finite stress-bands travelling over deep to shallow (designated as case I) and shallow to deep water (designated as case II). The left side of Figure 1 shows an atmospheric disturbance of the semi-infinite stress-band type moving from left to right for case II. In these calculations we will assume that the atmospheric disturbance travels only from left to right because the results for case I (case II) in which the disturbance travels from right to left can be easily inferred from the



and right side shows a finite stress band moving over shallower to deeper A schematic representation of the geometry of the problem. Left side shows a semi-infinite stress band moving over deeper to shallower water Fig. 1

results corresponding to the travel of the disturbance from left to right for case II (case I). The conditions imposed on the surface stress for the semi-infinite stress-band case which are valid for both cases I and II are

$$R = \begin{cases} 0 & \text{for } t \leq \frac{x}{V} \\ 1 & \text{for } t \geq \frac{x}{V} \end{cases}$$
 (10)

Next we will express the surface-stress conditions for a finite stress-band of width B. Since this is a linear problem, the solutions for the water levels at the left and right boundaries and for the set-up can be obtained by superposing two semi-infinite stress-bands, one positive and the other negative, both of the same intensity and both moving with the same speed V in the same direction but with their jumps separated by a distance B. The solutions for the negative stress-band can be obtained by replacing the time-variable t by t-to in the solutions for the positive stress-band. Here to=B/V is the time interval that elapses between the passage of the leading and the trailing edges of the stress-band over a given point of the lake. Hence the conditions imposed on the wind-stress for the finite stress-band case are

$$R = \begin{cases} 0 & \text{for } t \leq \frac{x}{V} \text{ and } t \geq \frac{x}{V} + t_0 \\ 1 & \text{for } \frac{x}{V} \leq t \leq \frac{x}{V} + t_0 \end{cases}$$
 (11)

In all these calculations the origin of time t=o is taken at the moment the stress-band is over the left boundary of the lake. The leading edge of a semi-infinite stress-band crosses

the right boundary of the lake at  $t_0=L/V$  after which the whole lake is under the influence of the wind-stress. However, for the finite stress-band, as can be seen from equation (11), a given point in the lake distant x from the left boundary is influenced by the wind-stress, only during the interval x/V to  $\frac{x}{V}$  +  $t_0$ .

Although we could have obtained solutions for any depth ratios of the two regions, we did the computations only for two cases; namely  $C_1/C_2=2$  and 1/2. The first case is one in which the region to the left of the discontinuity is deeper than the region to the right by a factor of four. In the second case the situation is reversed. Detailed algebraic solutions were obtained for the water level  $h_L$  at the left boundary and the water level  $h_R$  at the right boundary in terms of the wind-stress R, length L of the lake, the speed V of the atmospheric disturbance and either  $C_1$  or  $C_2$ , which ever is greater depending upon the case.

#### 3. The Characteristic Diagrams

The characteristic diagram is a representation of the characteristics (which identify the long gravity waves) in the x-t plane. Several different characteristic diagrams (subcases) are possible for a given case depending upon the ratio of V to  $C_1$  or  $C_2$ . Figure 2 shows the characteristic diagram for  $V/C_1=3$  and  $C_1/C_2=2$ . The dashed-dotted line shows the position of the leading edge of the stress-band which is usually referred to as the stress-jump line. The sloping lines in the diagram are the characteristics. Although this diagram is drawn for

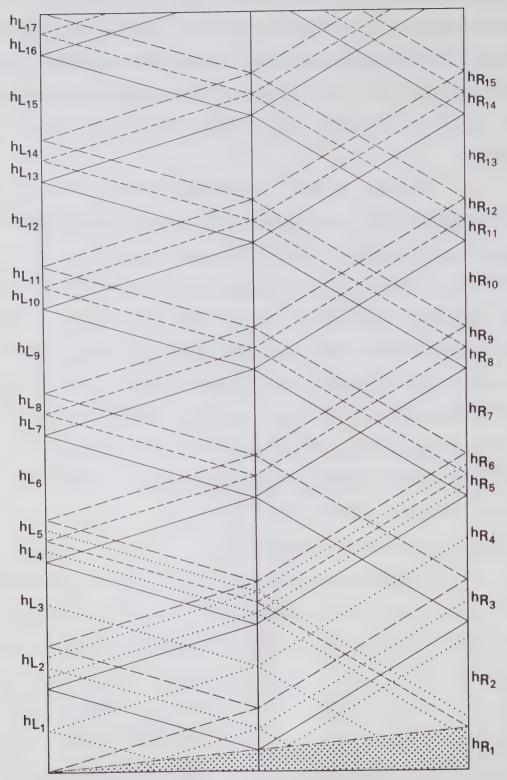


Fig. 2 Characteristic diagram for sub-case 1  $(\infty>V>2)$  of case I. The diagram is  $\overline{C}_1$  actually drawn for V=3.

 $V/C_1=3$ , it can be seen that the nature of the diagram remains the same for  $\infty>\frac{V}{C_1}>2$  except that as  $V/C_1$  increases, the slope of the stress-jump line decreases until at  $V/C_1=\infty$ , the stress-jump line merges with the abscissa. Since  $V>C_1$ , 2 for this sub-case, the stress-band propagates faster than the free gravity waves in both regions.

When the stress-band arrives over the left boundary of the lake it generates a free gravity wave that propagates in region 1 with uniform speed C1. Since in this sub-case V>C1, the stress-band arrives at the depth-discontinuity earlier than the free gravity wave that has originated at the left boundary. The stress-band generates two gravity waves at the discontinuity, one travelling with speed C1 in region 1 and the other travelling with speed C2 in region 2. These gravity wave systems thus created arrive at the respective boundaries in course of time and get reflected. When the gravity wave that has originated at the left boundary arrives at the discontinuity, it creates another two gravity waves. When the stress-band arrives at the right boundary it creates another system of gravity waves.

Thus, essentially we can distinguish among three families of gravity waves based on the location of the source of their creation: the L-family (originated at the left boundary), D-family (originated at the discontinuity) and the R-family (originated at the right boundary). Using these three families of characteristics, the characteristic diagram could be built up in advance before we attempt to obtain solutions for  $h_{\rm L}$  and  $h_{\rm R}$ .

The different possibilities of intermingling among these three families of characteristics and also the stress-band represent the different sub-cases. In principle, for each case an infinite number of sub-cases are possible. However, from a practical point of view we have limited our discussion to  $3 \geq \frac{V}{C_{1,2}} \geq \frac{1}{3}.$  It so happens that  $\infty > \frac{V}{C_{1,2}} > 2$  is only one sub-case. The possible sub-cases for case I in this range are listed in the Table.

Out of these twelve sub-cases, two sub-cases, namely  $\frac{V}{C_1} = 1$  and  $\frac{V}{C_2} = 1$  are not considered here because the method of characteristics fails when  $V = C_1$  or  $C_2$  and the solutions for  $h_L$  and  $h_R$  are discontinuous. Figures 3 and 4 show the characteristic diagrams for sub-cases 2 and 5 of case I and Figure 5 shows the diagram for sub-case 12 of case II. In sub-case 2 of case I there are only two families of characteristics because the R-family merges with the L-family at  $t = \frac{3}{2} \frac{L}{C_1}$  and  $x = \frac{L}{2}$  as can be seen from Figure 3. Since the diagrams are different for each sub-case, the solutions for  $h_L$  and  $h_R$  are different. Note that if  $D_1 = D_2 = D$  then  $C_1 = C_2 = C$  and only three sub-cases are possible, namely  $\frac{V}{C} \stackrel{>}{<} 1$ . If V/C = 1, the method of characteristics fails; however, the maximum value of the set-up is continuous (Rao, 1967).

## 4. Detailed Solutions

The detailed solutions for sub-case 1 of case I will be described here along with the technique of integration along the characteristics to obtain expressions for  $h_{\rm L}$  and  $h_{\rm R}$ . First the solutions on the left boundary of the lake will be obtained.

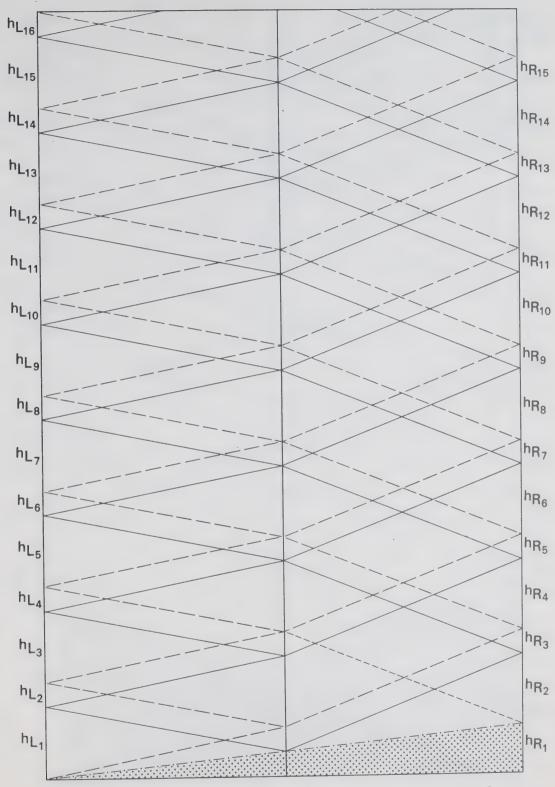


Fig. 3 Characteristic diagram for sub-case 2  $(\underline{V} = 2)$  for case I.

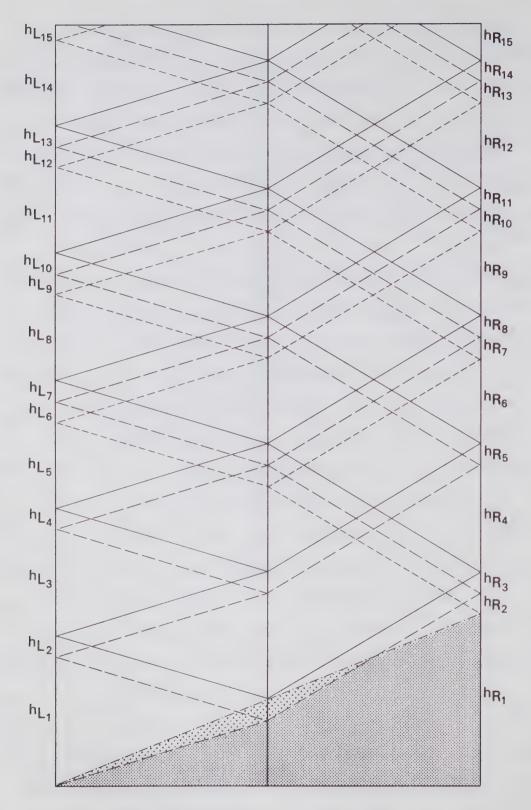


Fig. 4 Characteristic diagram for sub-case 5 (1>V/ $\overline{C}_1$ ) of case I.

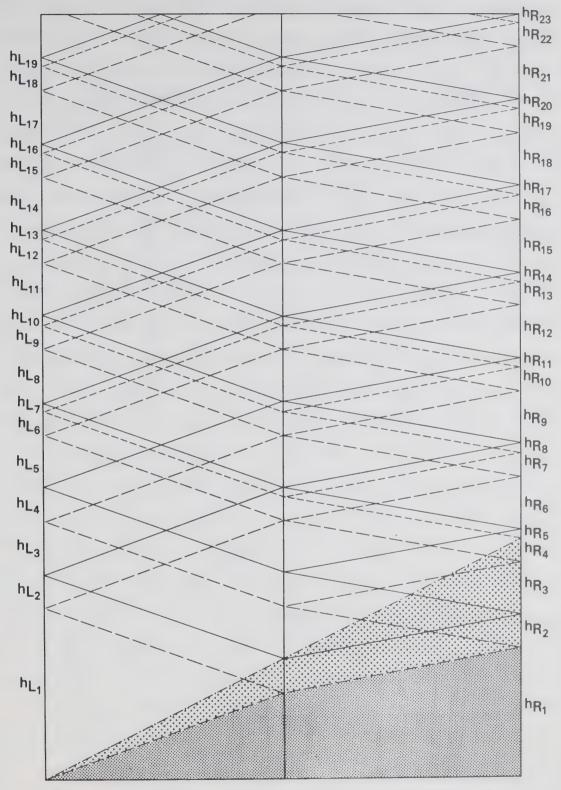


Fig. 5 Characteristic diagram for sub-case 12  $(\frac{V}{C_1} = \frac{1}{3})$  of case II.

Table Sub-case for case I  $(C_1 = 2C_2)$ 

Sub-case	Number	Rang	ge of	$\frac{v}{c_1}$
1		∞ >	$\frac{v}{c_1}$ >	2
2			2	
3		2 >	$\frac{v}{c_1}$ >	1
4			1	
5		1 >	$\frac{V}{C_1}$ >	$\frac{2}{3}$
6			<del>2</del> <del>3</del>	
7		$\frac{2}{3}$ >	$\frac{v}{c_1}$ >	$\frac{1}{2}$
8			$\frac{1}{2}$	
9		$\frac{1}{2}$ >	$\frac{v}{c_1}$ >	0.4
10			0.4	
11		0.4	$> \frac{V}{C_1}$	$->\frac{1}{3}$
12			$\frac{1}{3}$	

Since the lake is assumed to be at rest till the stress-band passes over it,  $h_L = 0$  at t = 0. As can be seen from Figure 2, an expression for hL that is valid for  $0 \le t \le \frac{L}{2V} + \frac{L}{2C_1}$  can be derived through integration of equation (9) along the characteristic AB which intersects the stress-jump line at the point A in region I. The solution for h, which is valid in the above time interval is designated as  $h_{L_1}$ . The solution for  $h_{L_1}$ valid for the time interval  $\frac{L}{2V}$  +  $\frac{L}{2C_1}$   $\leq$  t  $\leq$   $\frac{L}{C_1}$  and designated by  $h_{L_2}$  can be obtained through integration along the characteristics FD, DA, and DL. Of these, FD and DL are negative characteristics (since dx/dt is negative) and DA is a positive characteristic. The solutions  $h_{\mathrm{L}_1}$  and  $h_{\mathrm{L}_2}$  should match at the time  $t = \frac{L}{2V} + \frac{L}{2C_1}$  which is the upper time limit for  $h_{L_1}$  and the lower time limit for hL2. Otherwise the solutions are discontinuous and have no relevance. Solutions for  $h_{L_3}$ ,  $h_{L_4}$  etc. could be built up through integration along the proper characteristics.

For this sub-case 1 of case I (as for some other but not all sub-cases) the water near the right boundary of the lake is at rest till the stress-band passes over it at t =  $\frac{L}{V}$ . Thus  $h_{R_1}$  = 0 for 0  $\leq$  t  $\leq$   $\frac{L}{V}$ . The solution  $h_{R_2}$  for  $\frac{L}{V} \leq$  t  $\leq$   $\frac{L}{2V}$  +  $\frac{L}{2C_2}$  can be obtained through integration along the characteristics PD, DL and DA. The solutions for  $h_{R_3}$  and  $h_{R_4}$  can be constructed in a similar manner by integrating along the proper characteristics.

To illustrate the actual integration process along the characteristics, consider the integration along BA to obtain the solution for  $h_{\rm L_1}$ . Equation (9) can be written as

 $M_{1B} - C_1h_{1B} - M_{1A} + C_1h_{1A} = R(t_B-t_A)$ . From the boundary and initial conditions

$$M_{1R} = M_{1A} = h_{1A} = 0$$

Since  $h_{lB} = h_{L_1}$  the above equation becomes

$$- C_1 h_{L_1} = R(t_B - t_A)$$
 (12)

To determine  $t_B^-t_A^-$  one can use the information that A is the point of intersection of the stress-jump line x=V.t with the negative characteristic  $dx/dt = -C_1$ . This then gives

$$v.t_A = C_1(t_B-t_A)$$

Noting that  $t_B$  can be replaced by t where t denotes any time in the interval of 0  $\le$  t  $\le \frac{L}{2V}$  +  $\frac{L}{2C_1}$  , equation (12) can be written as

$$h_{L_1} = \frac{-R.V.t}{C_1(C_1+V)}$$
 (13)

This solution satisfies the condition  $h_{L_1} = 0$  at t = 0.

The solution for  $h_{L_2}$  requires integrations along FD, DA and DL. Integration along FA gives

$$M_{1F} - C_{1.h_{1F}} - M_{1D} + C_{1.h_{1D}} = R(t_{F}-t_{D})$$

but

$$M_{1F} = 0$$
;  $h_{1F} = h_{L_2}$  and  $t_F - t_D = \frac{L}{2C_1}$ 

Hence

$$h_{L_2} = \frac{-M_{1_D}}{C_1} + h_{1_D} - \frac{RL}{2C_1^2}$$
 (14)

Integration along DA gives

$$M_{1D} + C_{1}.h_{1D} - M_{1A} - C_{1}.h_{1A} = R(t_{D}-t_{A})$$

Noting that  $M_{1A} = h_{1A} = 0$  and determining  $t_D - t_A$  making use of the property that A is the point of intersection of the stress-jump line with the positive characteristic AD we get

$$M_{1_D} + C_1.h_{1_D} = \frac{R(\frac{L}{2} - V.t_D)}{(C_1 - V)}$$
 (15)

Integration along DL gives

$$M_{2D} - C_{2}.h_{2D} - M_{2L} + C_{2}.h_{2L} = R(t_{D}-t_{L})$$

But

$$h_{2D} = h_{1D}$$
;  $h_{2L} = 0$ ;  $M_{2L} = 0$  and  $M_{1D} = M_{2D}$ 

Also determining  $t_D$ - $t_L$  in a manner as before, the above expression becomes

$$M_{1_D} - \frac{C_1}{2} h_{1_D} = \frac{-R(L-2V.t_D)}{(C_1+2V)}$$
 (16)

Solving (15) and (16) for  $\mathrm{M}_{\mathrm{1D}}$  and  $\mathrm{h}_{\mathrm{1D}}$  we get the solutions for the volume transport and the water level at the discontinuity

$$M_{1D} = \frac{R(L-2V.t_D) (2V-C_1)}{2(2V+C_1) (C_1-V)}$$
(17)

and

$$h_{1D} = \frac{R(L-2V.t_D)}{(C_1-V)(C_1+2V)}$$
 (18)

Substitution of (17) and (18) into (14) gives the solution for  $h_{\rm L2}$ .

$$h_{L2} = \frac{R\left\{L\left(1 + \frac{V}{C_1}\right) - 2Vt\right\}(3C_1 - 2V)}{2C_1\left(C_1 - V\right)\left(C_1 + 2V\right)} - \frac{RL}{2C_1^2}$$
(19)

which is valid for

$$\frac{L}{2V} + \frac{L}{2C_1} \le t \le \frac{L}{C_1}$$

The solutions for  $h_{L_2}$  and  $h_{L_1}$  match correctly at  $t = \frac{L}{2V} + \frac{L}{C_1}$ . In a similar manner the solutions for  $h_{L_3}$ ,  $h_{L_4}$  etc. and  $h_{R_1}$ ,  $h_{R_2}$  etc. can be constructed. These solutions are summarized below for this sub-case. For i=3 and 4 only

$$h_{\text{Li}} = \frac{1}{3} h_{\text{Li-2}} - \frac{8}{3} \frac{R}{C_1} \left\{ \frac{\text{Vt} - \frac{L}{2} (1 + \frac{\text{V}}{C_1})}{(2\text{V} + C_1)} \right\} - \frac{RL}{3C_1^2}$$

$$(t - \frac{L}{C_1})$$

$$i = 3: \frac{L}{C_1} \le t \le \frac{L}{2V} + \frac{3}{2} \frac{L}{C_1}$$

$$i = 4$$
:  $\frac{L}{2V} + \frac{3}{2} \frac{L}{C_1} \le t \le \frac{L}{V} + \frac{3}{2} \frac{L}{C_1}$ 

These give the upper and lower time limits during which the solutions for  $h_{L_3}$  and  $h_{L_4}$  are valid. Note that the upper time limit for  $h_{L_3}$  coincides with the lower time limit for  $h_{L_4}$ , as it should be. The solution for  $h_{L_5}$  is given by

for 
$$\frac{L}{V}$$
 +  $\frac{3}{2}$   $\frac{L}{C_1}$   $\leq$  t  $\leq$   $\frac{2L}{C_1}$ 

For i = 6 and 7 only, the solutions are

$$i = 6: \frac{2L}{C_1} \le t \le \frac{L}{2V} + \frac{5}{2} \frac{L}{C_1}$$

$$i = 7: \frac{L}{2V} + \frac{5}{2} \frac{L}{C_1} \le t \le \frac{L}{V} + \frac{5}{2} \frac{L}{C_1}$$

Now we are in a position to write the general solutions on the left side. For  $i \geq 8$  (i takes only integer values)

Depending upon the time limits these general solutions fall into three types. For i=8, 11, 14, 17, etc. (general solution of first kind) the lower time limit is

$$T_{8+3J}' = \frac{L}{V} + \frac{L}{2C_1}$$
 (5 + 2J); the upper limit =  $T_{8+3J}^u = \frac{L}{C_1}$  (3 + J)

where

$$J = 0,1,2,3, etc.$$

For i=9, 12, 15, 18 etc. (general solution of second kind)

the lower time limit = 
$$T_{9+3J}' = \frac{L}{C_1}$$
 (3 + J)

the upper time limit = 
$$T_{9+3J}^{u} = \frac{L}{2V} + \frac{L}{2C_1} (7 + 2J)$$

For i=10, 13, 16, 19, etc. (general solution of third kind)

the lower time limit = 
$$T_{10+3J}^{\tau} = \frac{L}{2V} + \frac{L}{2C_1}$$
 (7 + 2J)

the upper time limit = 
$$T_{10+3J}^u = \frac{L}{V} + \frac{L}{2C_1} (7 + 2J)$$

This completes the solutions on the left side of the lake.

Next we will summarize the solutions for the water level on the right side of the lake.

$$h_{R1} = 0$$
 for  $0 \le t \le \frac{L}{V}$ 

$$h_{R2} = \frac{4R(L-Vt)}{C_1(C_1-2V)} \text{ for } \frac{L}{V} \le t \le \frac{L}{2V} + \frac{L}{C_1}$$

$$h_{R3} = \frac{2RV}{C_1} \left\{ \frac{2Vt - L(1 + \frac{2V}{C_1})}{(C_1 + 2V)(V - C_1)} \right\} + \frac{2RL}{C_1^2}$$

for 
$$\frac{L}{2V}$$
 +  $\frac{L}{C_1} \le t \le \frac{3}{2} \frac{L}{C_1}$ 

The solutions for  $\mathbf{h}_{R_4}$  and  $\mathbf{h}_{R_5}$  are similar. Thus for i=4 and 5 only

$$h_{R_{\dot{1}}} = \frac{4}{3} h_{L_{\dot{1}}-3} + \frac{2}{3} \frac{R}{C_{\dot{1}}} \left\{ \frac{2Vt - L(1 + \frac{2V}{C_{\dot{1}}})}{(2V + C_{\dot{1}})} \right\} + \frac{8}{3} \frac{RL}{C_{\dot{1}}^2}$$

The time limits are given by

$$i = 4: \frac{3}{2} \frac{L}{C_1} \le t \le \frac{L}{2V} + \frac{2L}{C_1}$$

$$i = 5$$
:  $\frac{L}{2V} + \frac{2L}{C_1} \le t \le \frac{L}{V} + \frac{2L}{C_1}$ 

The solution for  $h_{R_6}$  is given by

$$h_{R_6} = \frac{4}{3} h_{L_2} - \frac{1}{3} h_{R_2} + \frac{10}{3} \frac{RL}{C_1^2}$$

for 
$$\frac{L}{V}$$
 +  $\frac{2L}{C_1} \le t \le \frac{5}{2} \frac{L}{C_1}$ 

The solutions for  $h_{R_7}$  and  $h_{R_8}$  are similar and are given by the following expression. For i=7 and 8 only

$$h_{R_{i}} = \frac{4}{3} h_{L_{i-4}} - \frac{1}{3} h_{R_{i-5}} + \frac{10}{3} \frac{RL}{C_{1}^{2}} + (t - \frac{3}{2} \frac{L}{C_{1}})$$

The time limits are given by

$$i = 7: \frac{5}{2} \frac{L}{C_1} \le t \le \frac{L}{2V} + \frac{3L}{C_1}$$

$$i = 8: \quad \frac{L}{2V} + \frac{3L}{C_1} \le t \le \frac{L}{V} + \frac{3L}{C_1}$$

Now the general solution on the right side can be written. For  $i\,\geq\,9$ 

Depending upon the time limits, these general solutions fall into three categories. For i=9, 12, 15, 18, etc. (general solution of the first kind) we have the following limits;

the lower time limit = 
$$T_{9+3J}' = \frac{L}{V} + \frac{L}{C_1}$$
 (3 + J)

the upper time limit = 
$$T_{9+3J}^{u} = \frac{L}{2C_1} (7 + 2J)$$

For i=10, 13, 16, 19, etc. (general solution of the second kind) the lower time limit =  $T_{10+3J}' = \frac{L}{2C_1}$  (7 + 2J)

the upper time limit =  $T_{10+3J}^{u} = \frac{L}{2V} + \frac{L}{C_1} (4 + J)$ 

For i=11, 14, 17, 20, etc. (general solution of the third kind)

the lower time limit = 
$$T_{11+3J}' = \frac{L}{2V} + \frac{L}{C_1} (4 + J)$$

the upper time limit = 
$$T_{11+3J}^{u} = \frac{L}{V} + \frac{L}{C_1} (4 + J)$$

This completes the solutions on the right side.

The setup is calculated from the relation

$$\Delta h_t = h_{R_t} - h_{L_t}$$

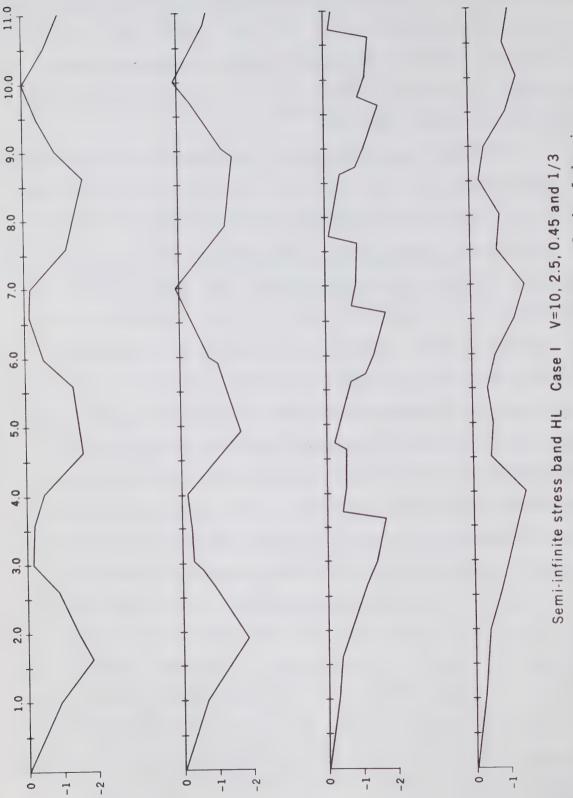
This completes the solutions for sub-case 1 of case I. In a similar manner the solutions for all sub-cases of cases I and II for the semi-infinite stress-band could be determined. The method of obtaining the solutions for the finite stress-band starting from those for the semi-infinite stress-band has been described in Section 2. The solutions for the finite stress-band were determined for band widths of 2L, 1.5L, L, 0.75L, 0.5L and 0.25L where L is the length of the lake.

#### 5. Results

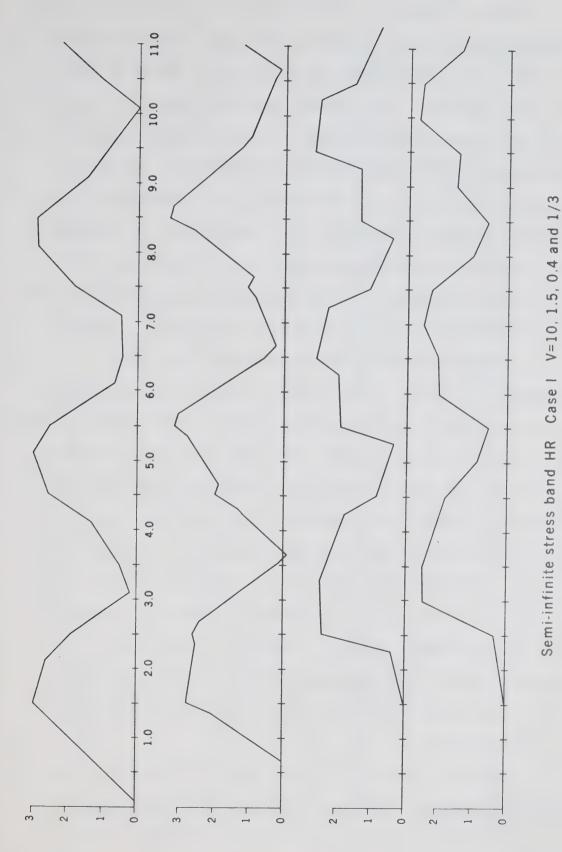
Most of the results of the computations are presented in Figures 6 to 16 all with the same format, namely the amplitude of the water level as ordinate and the time as abscissa. The numbers marked along both the abscissa and the ordinate are nondimensional quantities. To get the corresponding dimensional

numbers, for cases I and II, the ordinate values have to be multiplied by  $RL/C_1^2$  and  $RL/C_2^2$  respectively. The abscissa values have to be multiplied by  $L/2C_1$  for case I and  $L/2C_2$  for case II. Note that the values of V quoted in the diagrams are non-dimensional, being expressed as a ratio to  $C_1$  for case I and as a ratio to  $C_2$  for case II.

Figure 6 shows the water level variation with time at the left boundary of the lake for a semi-infinite stress-band moving over deep to shallow water for four different speeds. The interesting result is that the water level is predominantly negative. On this diagram the water level becomes positive only for V = 2.5 around t = 7 and t = 10. Since the figures are self-explanatory, lengthy descriptions will be omitted. Since bottom frictional effects are not included in this computation our results become less reliable with increasing In every case the computations have been made for V ranging from 10 to 1/3 but only those cases which showed interesting results are included in the diagram. Figure 7 shows the variation of water level at the right boundary of the lake for a semi-infinite stress-band moving over deep to shallow water for four different speeds. It is immediately clear from this diagram that the water level is predominantly positive. In fact, in this diagram, it becomes negative only for V = 1.5 around t = 3.6. It is clear from Figures 6 and 7 that the water level, even in the absence of friction, does not become periodic, although from some values of V it assumes a quasi-periodic nature.



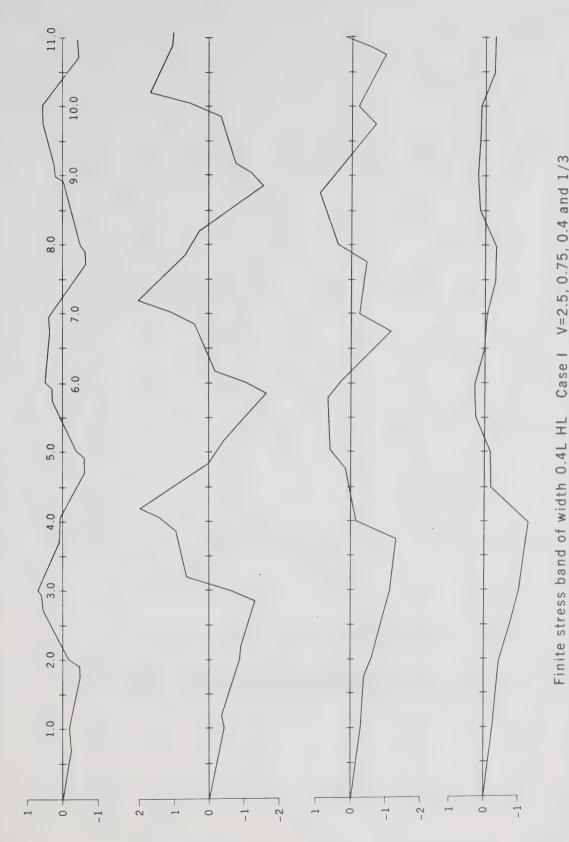
Variation of water level at the left boundary of the lake in case I for a semi-infinite stress-band. Fig. 6



Variation of water level at the right boundary in case I for a semi-infinite stress-band. Fig.

Figure 8 shows the time variation of  $\mathbf{h}_{\mathsf{T}_{\mathsf{L}}}$  for a finite stress-band of width 0.4L in case I for four different speeds. First of all it is clear that the water level could be both positive and negative. The second important point is that the nature of the water level variation is drastically different for different speeds of movement of the stress-band. In this diagram the maximum range of the water level variation occurs for V = 0.75 probably because of the possibility of resonant coupling between the stress-band and the gravity wave in region I. The cases of V = 1 and 1/2 cannot be treated by the method of characteristics for both cases I and II. Figure 9 shows the situation at the right boundary of the lake corresponding to Figure 8 for speeds V = 1.5, 2/3, 0.6 and 0.4. Again, as on the left side the water level is both positive and negative but in contrast to the left side the range of water level variation on the right side is uniformly high for the four different speeds. By comparison with the previous figures it can be seen that the qualitative nature of the variation in case II is quite different from case I. Figure 10 shows the water level  $h_{
m L}$  in case II for a finite stress-band of width 0.1L for four different speeds. Figure 11 shows the corresponding water level variation at the right boundary of the lake. The range of variation is considerably lower than at the left boundary.

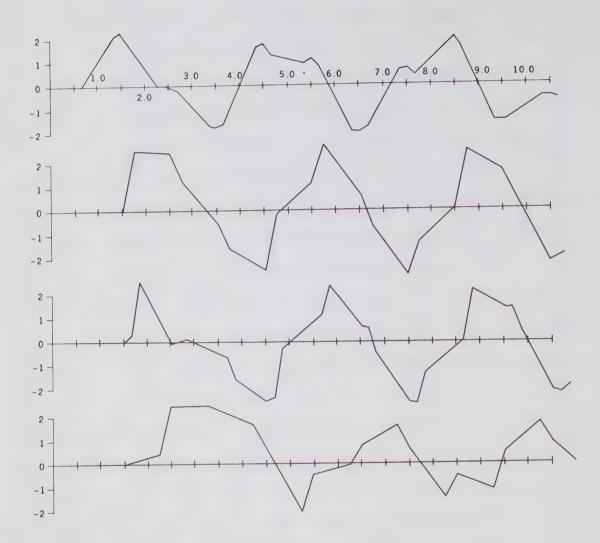
Figure 12 shows  $h_{\rm L}$  in case I for V = 1/3 for four different band widths. It is interesting to note that the amplitude of variation decreases with the decrease of width of



Variation of water level at the left boundary in case I for a finite stress-band of width 0.4L where L is the length of the lake. Case 1 V=2.5, 0.75, 0.4 and 1/3

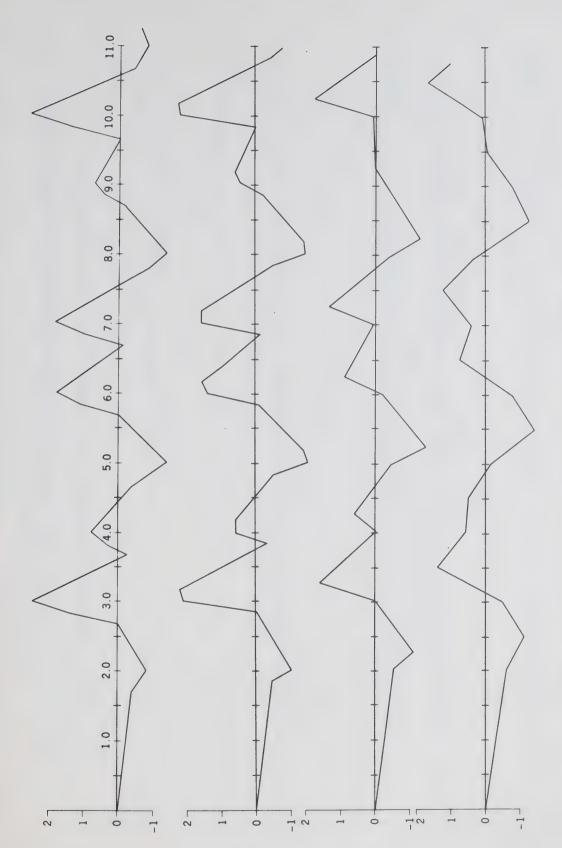
 $\infty$ 

Fig.

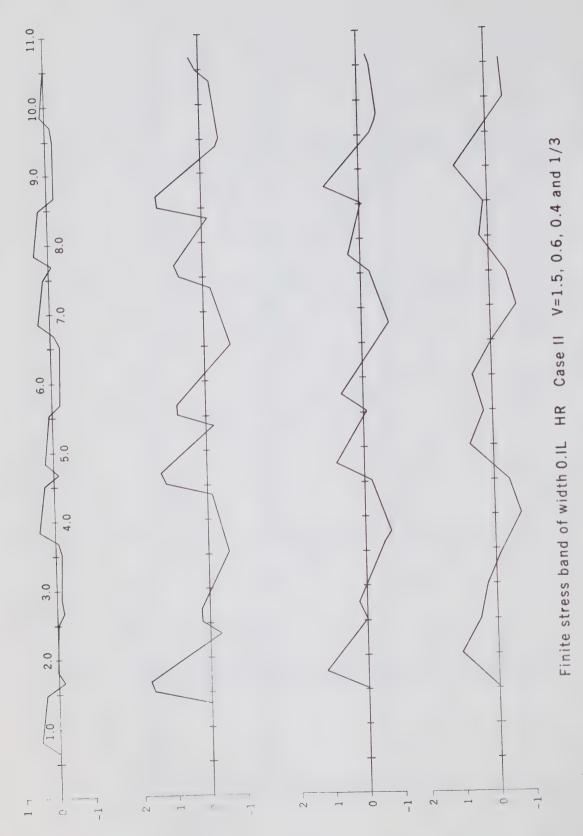


Finite stress-band of width 0.4L HR Case I V=1.5, 2/3, 0.6 and 0.4

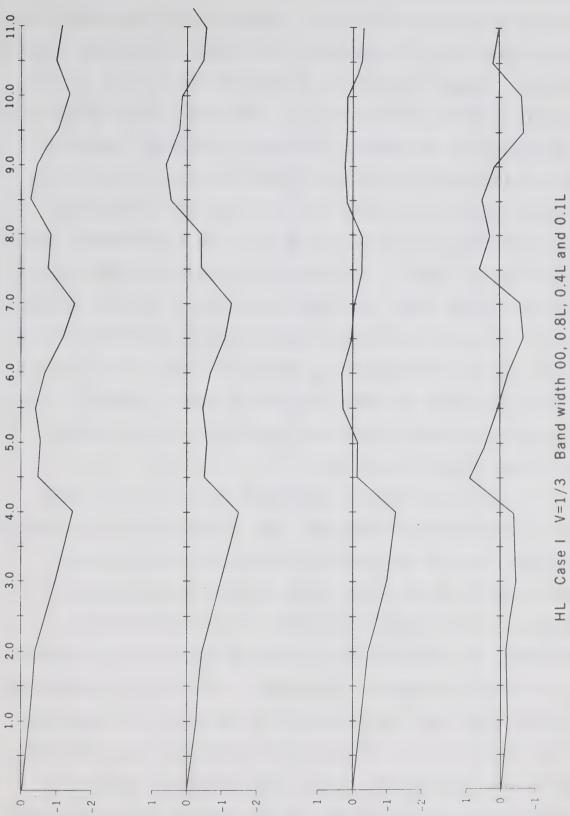
Fig. 9 Water level variation at the right boundary in case I for a finite stress-band of width 0.4L.



Q for Water level variation at the left boundary in case II finite stress-band of width 0.1L. Case II V=0.75, 0.6, 0.4 and 1/3 Finite stress band of width 0.1L HL Fig. 10



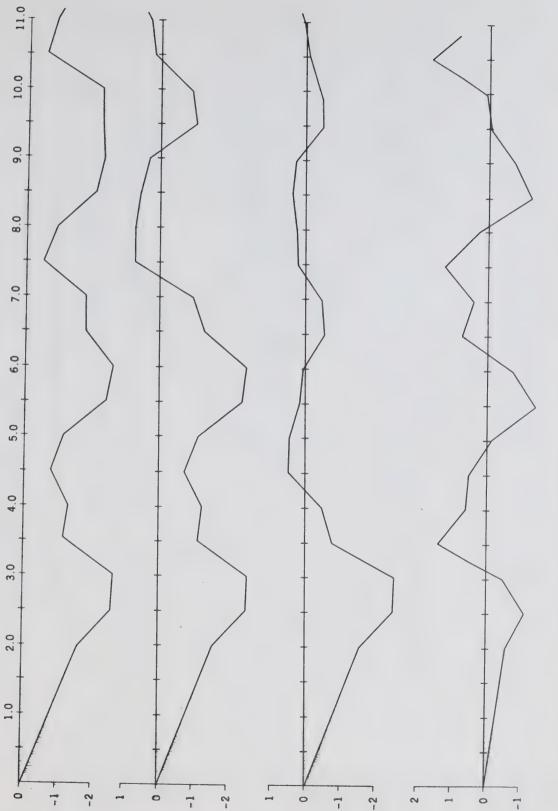
ಹ Water level variation at the right boundary in case II for finite stress-band of width 0.1L. 11 Fig.



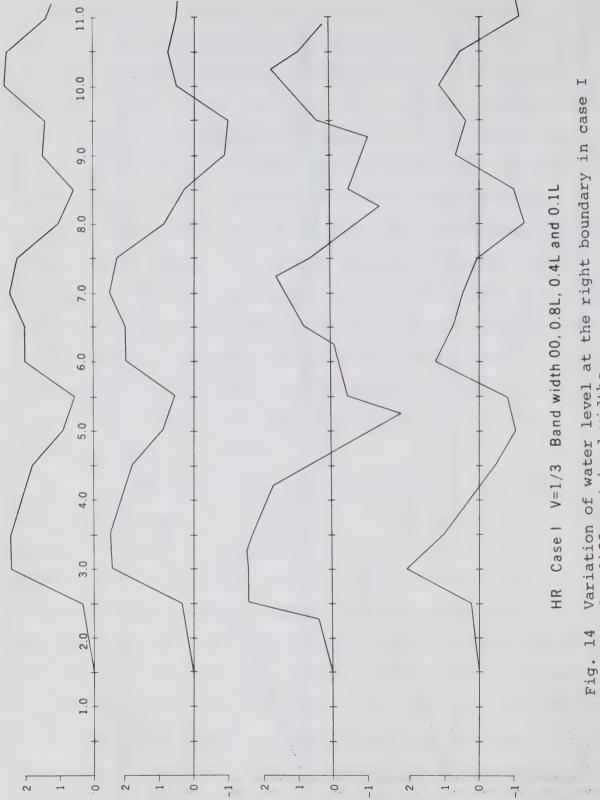
Variation of water level at the left boundary in case I for different band widths. Fig. 12

the band until 0.4L and it increases again slightly for 0.1L. Thus the width of the band with respect to the dimension of the lake is important in determining the amount of resonant coupling possible. Figure 13 shows  $h_r$  in case II for V = 1/3 for the same band widths as in Figure 12. The range of variation on the whole appears to be greater in Figure 13 than in Figure 12. Thus the amplitude of the water level variation at the left boundary appears to be more for a stress-band moving over shallow to deep water compared to that for a band moving over deep to shallow water. By analogy we expect the amplitude of variation at the right boundary to be greater in case I than in case II. This is exactly the case as can be seen from Figures 14 and 15. Figure 16 compares h<sub>I.</sub> and h<sub>R</sub> in cases I and II for a finite stress-band of width 0.4L moving with a speed V = 1/3. It can be seen that  $h_R$  for case I and  $h_I$  for case II have higher amplitude of variation.

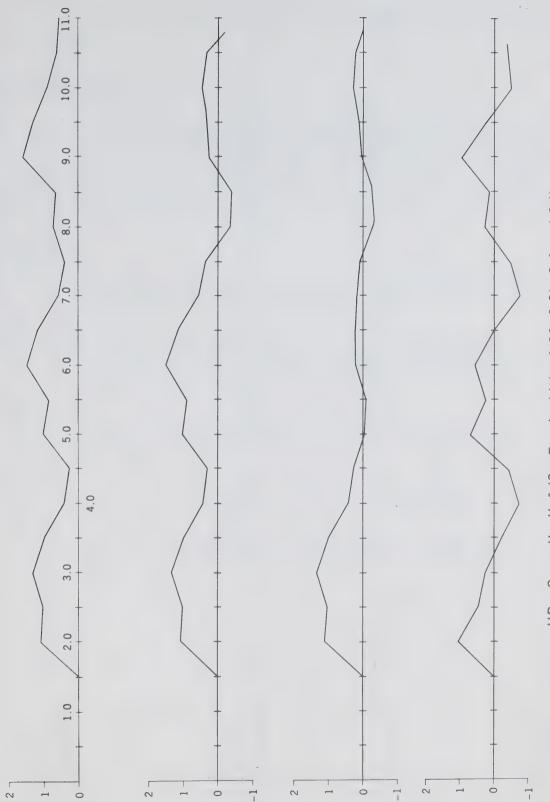
In conclusion the important results of this study could be summarized as follows. For a semi-infinite stress-band the water level at the left side could be predominantly negative while at the right side it could be predominantly positive. As the speed of movement of the stress-band approaches the speed of long gravity waves either in regions I or II, resonant coupling is possible. For finite stress-bands of width about half the dimension of the lake, the amplitude of the water level variation could be smaller than that for either wider or narrower bands. The additional complexity in reflections introduced due to the depth discontinuity does



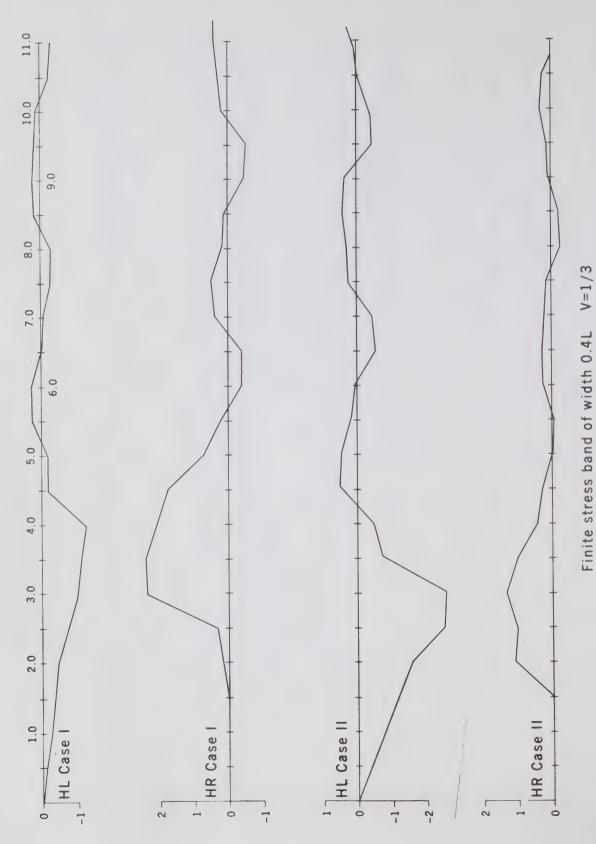
Variation of water level at the left boundary in case II for Band width 00, 0.8L, 0.4L and 0.1L different band widths. HL Case II V=1/3 Fig. 13



Variation of water level at the right boundary in case I for different band widths. Fig. 14



Variation of water level at the right boundary in case II for different band widths. HR Case II V=1/3 Band width of 00, 0.8L, 0.4 and 0.1L Fig. 15



Variation of water level at left and right boundaries in cases I and II for a finite stress-band of width 0.4L. Fig. 16

not allow exact periodicity of the setup variation even in the absence of friction although quasi-periodicity is achieved in some cases.

#### **ACKNOWLEDGEMENTS**

I thank Dr. D.B. Rao for introducing me to this topic and for enlightening discussions throughout the progress of this study. Mr. Robert Engley did the programming for the computer and Mr. John Cookson provided the illustrations. I also thank Miss T. Millington for helping me in the final preparation of this manuscript and Mrs. P. McAinsh for typing the final copy.

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# JAMES BAY

MANUSCRIPT REPORT SERIES

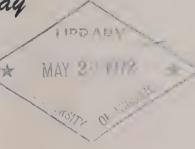
No.24

On the oceanography of James Bay F.G. Barber

The tides of James Bay
G. Godin

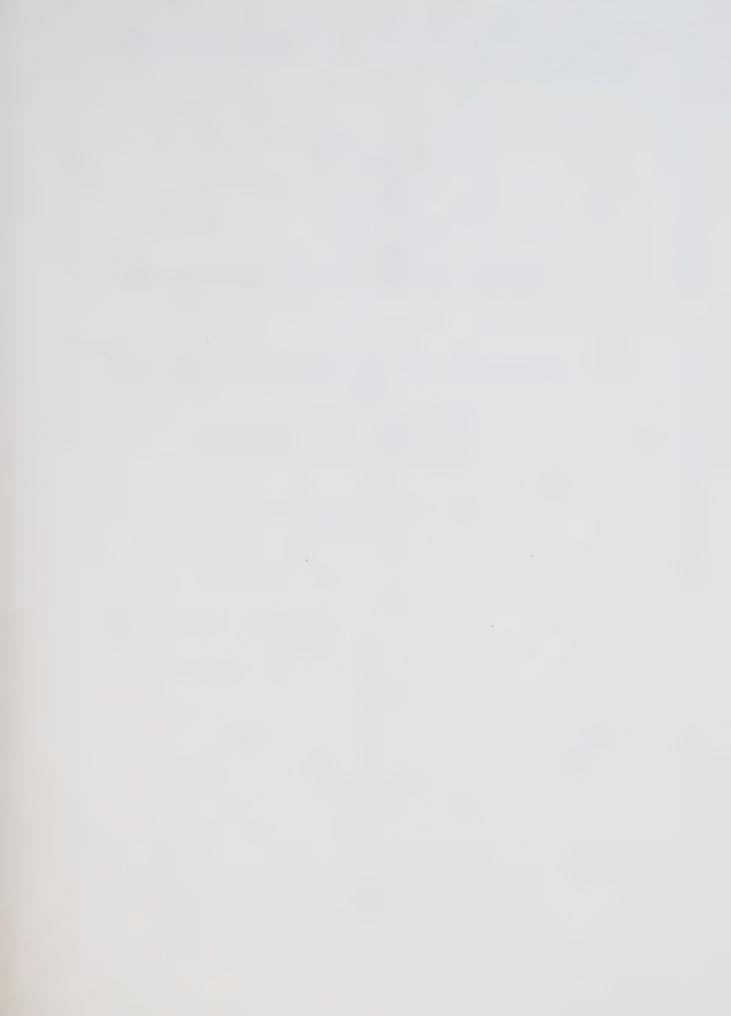
Circulation in James Bay

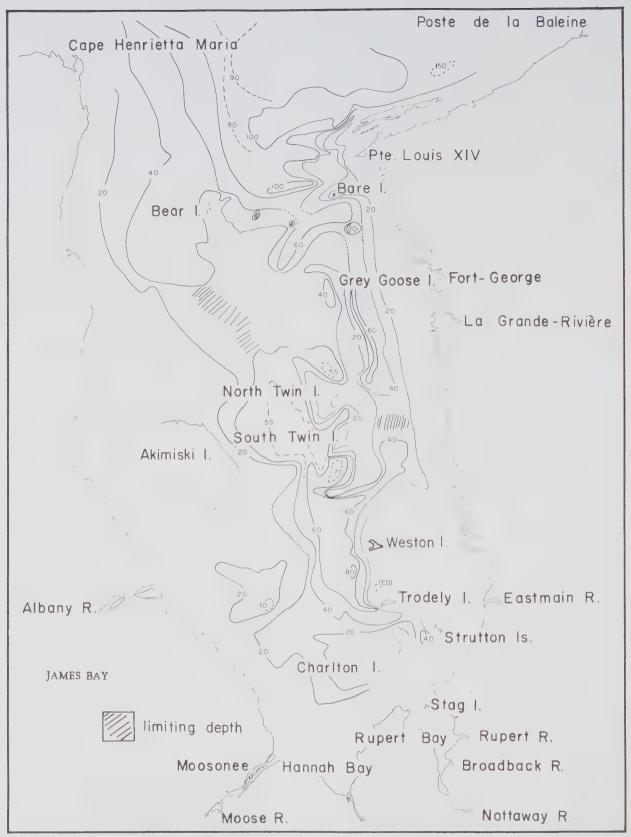
T.S. Murty



Marine Sciences Branch
Department of the Environment, Ottawa







Frontispiece Bathymetry (metres) of James Bay and some place names. An interpretation of bathymetric data for James Bay from Canadian Hydrographic Service Chart 5800, edition 1971, and from topographic maps of the Surveys and Mapping Branch.

## Manuscript Report Series No.24

ON THE OCEANOGRAPHY OF JAMES BAY
F.G. Barber

THE TIDES OF JAMES BAY
G. Godin

CIRCULATION IN JAMES BAY
T.S. Murty



#### Foreword

The reports contained in this number of our manuscript report series were prepared in response to a request to Federal authorities from the government of the Province of Quebec for an initial consideration of the possible influence of the James Bay Project on the environment; the three reports are directed to the marine environment and comprise but a portion of the total activity. Specifically, the study was to be directed to a question concerning the area to be developed first, i.e. whether construction should begin on the southern rivers or on the northern rivers. It was realized from the outset that we could not, for a number of reasons, contribute significantly to the resolution of the question; nevertheless, it was deemed useful that we attempt to understand the oceanography of James Bay as well as available data would permit.

A.E. Collin,

Director,

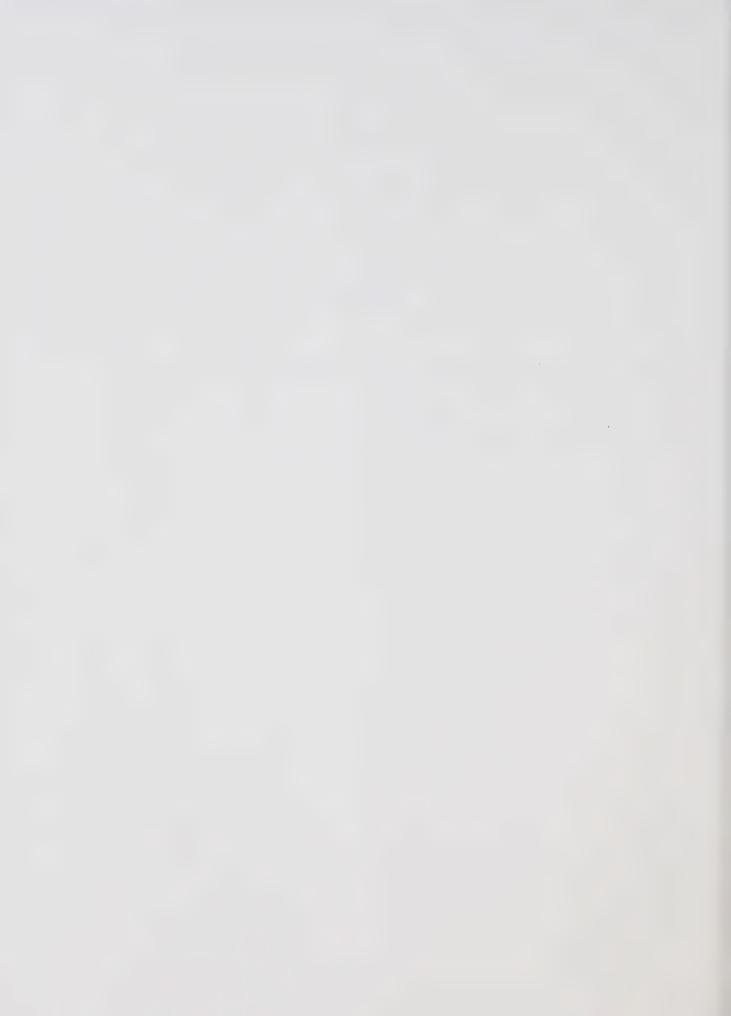
Marine Sciences Branch.

December 2, 1971 Ottawa



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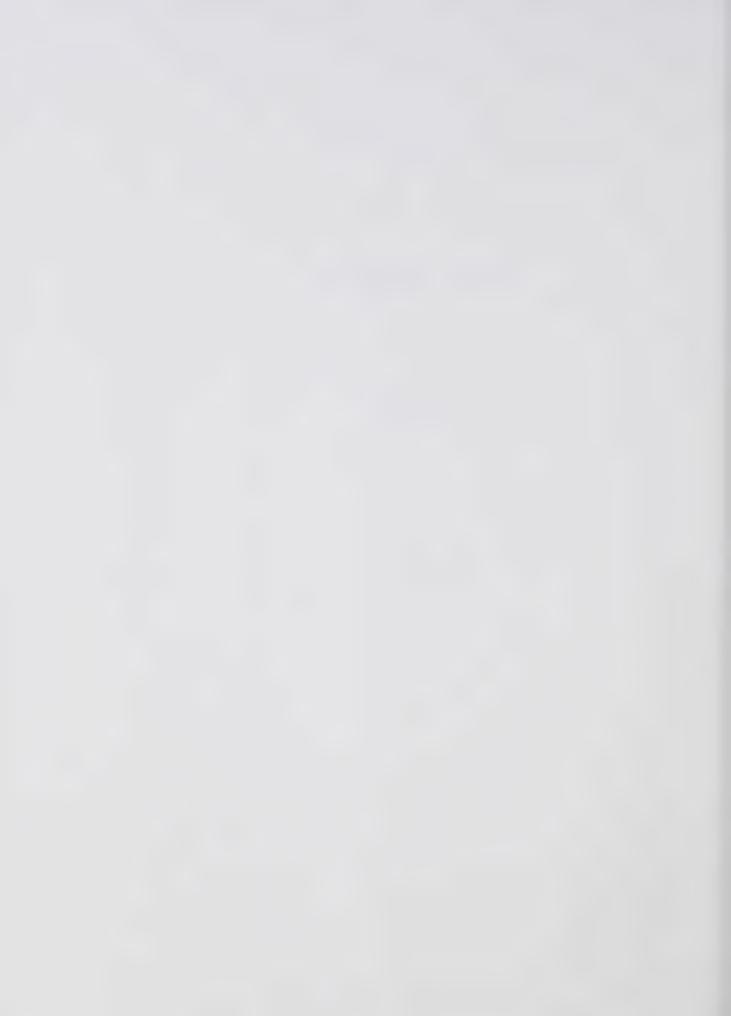


On the Oceanography

of

James Bay

F.G. Barber



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#### 1. Introduction

The examination of certain aspects of the oceanography of James Bay (frontispiece) contained in the following was carried out in a consideration of the possible influence on the water of James Bay of the development of the hydroelectric potential there. The work constitutes more of a review or reappraisal of knowledge of the region than it does an impact study, but was believed necessary in our initial approach to the problem. The review aspect emphasizes the importance to James Bay of processes operating in adjacent waters (Figure 1) and the fact that all of these waters generally reflect to similar degree the influence of similar internal and external factors; in consideration of this, the word "system" has been used to include all the region (Barber, 1967, p20).

It was not surprising that an assessment of the influence of the development would prove almost intractable for, while data on the region are few, our understanding of both the qualitative and quantitative changes which may occur through oceanographic processes is limited.

The consideration applies to the system as a whole as well as to James Bay, for not only are certain major features not understood but also much of the "understanding" is not well-founded; we have a great deal to learn about our northern waters generally. Part of the understanding could result through a study of the changes actually wrought on the system by the project. It is recommended therefore



Figure 1. Some place names within the system.

that consideration be given this so as to ensure that an adequate description of the system as it existed prior to development is available. Of course this will not ensure that post-development changes will be related in a causal way to the development, for if a measurable change does occur over a decade or two, or longer, other influences may appear by that time to be of equal or greater significance (Anon., 1971a, p45).

requirement is rather formidable, not only in oceanography (Hamill, 1969, p37) but also in hydrography and perhaps in geography. Robinson (1968) reviewed knowledge of the geography of the surrounding land areas where two subregions could be delineated (p202). One, "the South Coast Lowland", extends southeast of Churchill to west and south of James Bay where along the coast (p217),

...there is a flat strip five to ten miles wide, with the widest parts generally being to the north. The coastal zone is treeless, but grass or marshes are common. Storm beaches, a few feet high, are the only topographic features. Tidal flats may be exposed for one to six miles; even at high tide shallow water extends far offshore. Much of the monotonous coast has been unapproachable from the sea even by small ships; this being one of the main reasons for the lack of settlement. Deeper water may be found at the drowned river mouths, but shifting sand bars and minor deltaic deposits are navigation hazards.

The other, the "East Coast Upland" (p222) extends from east of James Bay to the north coast of Hudson Bay a north-south distance of 750 miles. James Bay constitutes almost a third of this, i.e. 230 miles (389 km), and is up to about 100 miles

(185 km) wide with an average depth of about 32 m. It seems likely that given adequate hydrographic data, distinct sub-regions would be recognizable within the bay. The project will likely influence the priority for such data so that the required knowledge of the bathymetry so important to an understanding of the oceanography may become available reasonably soon.

As yet only limited information is available concerning the project beyond that reported in newspapers, eg. Anonymous, 1971b, which in the main have been about the development of hydroelectric potential on the three southern rivers, Broadback, Nottaway and Rupert. A parallel development is proposed for certain of the northern rivers with diversion of water into La Grande-Rivière, including part of Grande-Rivière de la Baleine (Anon., 1971c). Considerable interest in possible long-term climatic effects of the project arose subsequent to newspaper, radio and television commentary based on expressions of two well-known Canadian scientists (R.W. Stewart and L.M. Dickie), which in turn led some to the conclusion that marked change was to be expected, eg. Time Magazine for September 27 (Anon., 1971d). In a subsequent letter to the magazine, Stewart and Dickie (1971) clarified their views somewhat and seemed to suggest that although it was not possible to predict the influence of the project, it was possible to predict that it will have one. Apparently the influence as they view it will be only on the local climate, or microclimate, and could occur in a number of ways, some rather subtle, requiring

detailed and extensive data in order to recognize the change and process (Landsberg, 1970, pl269). While it is by no means certain whether man's activities have yet influenced global climate (Frisken, 1971), it is the experience that this has not yet occurred (Landsberg, 1970) and may not be likely (Anon. 1970a, p97).

In the following it is assumed that a major feature of the oceanography of the region, the annual ice cover, is due to the global atmospheric circulation and hence not likely influenced by man. The ice cover limits the exchange of heat with the atmosphere so that both the gain of heat during the open season and the loss of heat during the period of ice cover are much smaller than were the area without an ice cover. It is believed that for the region, i.e. Hudson Bay and James Bay, the gain is about equal to the loss although smaller, so that a small deficit occurs which is balanced by an advection from the ocean. In this regard the main parts of the system may be even more uncoupled from the world ocean than is the Arctic Ocean, where a perennial ice cover is general and the balance is maintained by advection from the world ocean and an excess of export over import of ice. It is believed that the export and import of ice from and to our system are small and about equal, so that the amount of ice which melts is equivalent to that which forms there, although considerable redistribution of the ice occurs within the system.

It is possible that the project could influence this redistribution of ice (Murty, 1971) and some attention is given this aspect later. The treatment generally is quite inadequate and does not provide significant insight to those factors of the circulation which might influence ice distribution. This is of particular significance when we realize that not only is the ice cover unstable but also in some years a portion very nearly survives through a summer, eg. in 1969 (see section 2.3). If the project were to lower the average air temperature in summer through an impact on local climate, the ice cover would persist for a longer time causing a further decline in average temperature. It may be that through observation and study it could be determined whether the project will constitute a "major environmental disturbance" (Kershaw, 1971) or whether our considerations which suggest little impact are, in fact, "miscalculations" (Heyneman, 1971).

#### 1.1 Climate

In contrast to the broad expanses of Arctic tundra that surround Hudson Bay the sub-Arctic lands bordering James Bay are partially forested and thus protected from strong winds. As a result two of the most distinguishing features of Hudson Bay winter climate — wind chill and blowing snow — are not nearly as evident near James Bay. With this important exception and the fact that James Bay is several degrees of latitude farther south the factors that influence the climate of James Bay are essentially the same as those outlined for Hudson Bay.

Along southern coastlines more than 100 inches of winter snowfall are evenly distributed during the months November through March. Fort George, on the east coast, receives almost as much snow but approximately one half of the annual total falls during November and December. Lesser amounts are measured in the mid-winter months

when almost solid ice cover over James Bay cuts off the supply of moisture to the air moving over it. Throughout the winter the snow is soft and deep in the woods but drifted and hard packed over the exposed Bay ice.

In general the lands bordering James Bay are free of snow from late in May until the middle of October. During this period precipitation, mostly in the form of rain, accounts for more than one-half of the 25 to 30-inch annual total. At Moosonee, in fact, rain or drizzle are reported on approximately 15 days of each month in summer. As would be expected from these figures, cloudy days are frequent. While Moosonee and Fort George receive less than 240 hours of bright sunshine in July, stations at approximately the same latitudes in the prairie provinces record close to 320 hours. Moosonee ranks above Churchill in the number of thunderstorms that occur in summer with an average of three or four per month. The frequency of thunderstorms decreases markedly to the north and east however, where the cold waters of Hudson and James Bays tend to inhibit shower development.

Although reports from the weather stations at Moosonee and Fort George do not fully substantiate it, fog is probably quite prevalent over James Bay in June and July when sea ice is still present.

As is the case with Hudson Bay, the cold waters are more influential than latitude in determining summer air temperatures along the shores of James Bay. July temperatures at Moosonee and Moose Factory average about 60°F while daily maxima are near 70°F. Corresponding temperatures for Fort George are about 5 degrees lower. Along the east coast of James Bay daytime high temperatures over 85°F are unusual. At Moosonee, on the other hand, where southwesterly winds bring warm air from the heart of the continent, 90°F temperatures have been recorded in all months from May to August. Sharp and rather frequent temperature fluctuations may be expected during these months depending on whether the winds are off the land or off the water.

Freezing temperatures have been reported in every month of the year at weather stations on the shores of James Bay. While there are wide local variations in the incidence of frost, depending on the proximity of the Bay or on the presence of lakes or muskeg, the first frosts of autumn usually occur early in September. Average daily temperatures generally remain above 32°F, however, until late October.

January is the coldest month in the James Bay region with average temperatures in the neighbourhood of -5° to -10°F. An extreme minimum temperature of -52°F has been recorded at Moosonee and temperatures of -40°F are not uncommon. Readings as low as -20°F may be expected on one day in four at the height of the winter season. In contrast to the Hudson Bay area where above freezing temperatures are rare in winter, James Bay has the occasional mild spell in January and February. An example of such mild conditions occurred in February 1954, when maximum temperatures at Moosonee exceeded 40°F on five successive days, culminating in a reading of 51°F on the 19th.

In summary, the climate of James Bay is not as severe as that of Hudson Bay. Winters are long and cold however, and summers cool, and on an annual basis the lands surrounding James Bay are colder than most in Canada at similar latitudes.

That the "lands surrounding James Bay are colder than most in Canada at similar latitude" is probably due to the "deep southward penetration of Arctic climate in eastern Canada in winter" caused by alteration of "global wind patterns" by geographic features. The exerpt and the guotes are from Thompson (1968) who also provided climatological data for a number of locations around Hudson Bay and James Bay. Stressed in the article (p267) is the effect of ice and relatively cold water on the local climate whereby winter conditions become quite continental and those of summer quite maritime. Earlier Burbidge (1951) discussed the influence of the open water of Hudson Bay on continental polar air moving over the region. Apparently the influence is greatest over the eastern portion in mid-summer and in early winter prior to freeze-up; although the tempering influence does persist past freeze-up into January (Hagglund and Thompson, 1964). It is clear that a persistent change in the average extent of ice cover would be followed by

a change in climate there. As noted, it is generally believed that ice extent in the northern hemisphere is determined largely by global atmospheric patterns (see also Fletcher, 1969); however, in Hudson Bay the persistence of ice is particularly sensitive to warm air from the south (Mackay, 1952) or, as we shall see, to cold air from the north.

#### 1.2 Ice cover and a possible consequence

The "deep southward penetration of Arctic climate into eastern Canada in winter" is such that ice thickness over much of Hudson Bay approaches that for Arctic regions generally, although for James Bay it is considerably less (for example see Bilello and Bates, 1971).

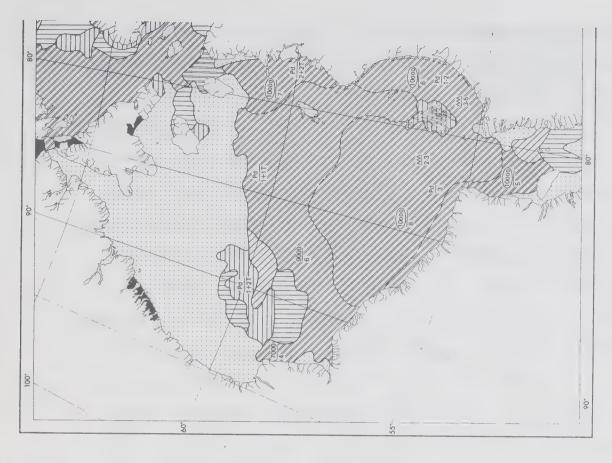
The greater part of James Bay is frozen over in the winter. By the end of January, local inhabitants have been known to go across the ice from the eastern shore to almost every island lying in the middle of the bay. At the head of the bay, in mid-winter, frequent journeys are made across the ice in a direct line between Charlton Island (Lat. 52°00'N., Long. 79°25'W.) and Moose River and between Charlton Island and Rupert River. The movement of ice in spring is greatly affected by the direction of the wind. If southerly winds predominate, the wind and outflow of the rivers acting together clear the southern part of the bay by the middle of June. But if northerly winds predominate, the ice will remain in the bay and seriously obstruct navigation till the sun becomes strong enough to melt it.

Coasting vessels are put into commission during the last days in June, and make their passage from one post to another in lanes of open water between the coast and the ice in the centre of the bay. (Anon., 1965, p441).

Larnder (1968) stressed the variability of the ice cover in Hudson Bay, both in its formation and breakup, which she noted differed "widely from year to year and from one locality to

another" (p319) and, "Although, as already noted, the ice in James Bay generally melts or moves out of the southern two-thirds of the bay by early July, the ice that moves into its northern part from Hudson Bay may be so heavy, close-packed and persistent" (p335) that local transportation may be limited. Danielson (1971, pl02) also considered that the northern "parts of James Bay receive ice from Hudson Bay" and implied that like southwest Hudson Bay much more ice melts there than forms there (p98). From the charts of ice conditions provided annually by the Meteorological Service it is possible to obtain an appreciation of the extent to which breakup can vary. For example, observations on July 9 of each of 1968 and 1969 (Figure 2) indicated quite a wide divergence in extent of ice cover. This may be due to variations in direction of persistent winds early in the summer which can accumulate or disperse ice. This in turn can influence the surface albedo and hence the amount of insolation eventually absorbed. A pattern of circulation associated with mixed freshwater moving seaward at the surface probably also influences the distribution of ice, but not likely to the same extent as can this early wind.

As James Bay is a region of annual, rather than perennial ice cover and as open water generally occurs at a relatively early time in summer, a seasonal variation in temperature and salinity may be anticipated. In the approach to James Bay the annual sequence of events is, in the main features, likely similar to that suggested





9 of 1968 and 1969 (from Anon., 1970e 1969. cover on July (a) 1968. (b) The extent of ice pl3; 1971e pl3). Figure

for a more northerly location in Hudson Bay (Figure 3), where the variation of salinity within a perennial surface layer is largely related to change in the ice cover. Within James Bay, the salinity structure is likely strongly influenced by the amount of freshwater from runoff and in the southern portion it may be that in winter a layer of freshwater exists under the ice (see also section 2.2.2). The existence of such a layer would be a consequence of the marked decrease of mixing at the surface due to a reduction in the transport of energy, from wind and exchange processes, across the sea surface resulting from the existence of an ice cover. Tidal currents would provide some energy for mixing but as these are not likely much above about 1 knot away from shore (Godin, 1971), it is possible the layer persists. Thus, in the absence of observations, the possible late winter distribution of salinity and temperature (Table 1) for the deep water west of Trodely Island is highly speculative, perhaps even fanciful. The small maximum temperature in a halocline is characteristic of arctic seas, as is the halocline itself. It seems almost certain that the halocline exists in winter, so that the water below a surface layer may be effectively uncoupled from local surface processes, i.e. convective processes are limited to depths above the halocline. This situation is believed to exist throughout the year over most of James Bay, so that the deeper water probably is from Hudson Bay with relatively little change.

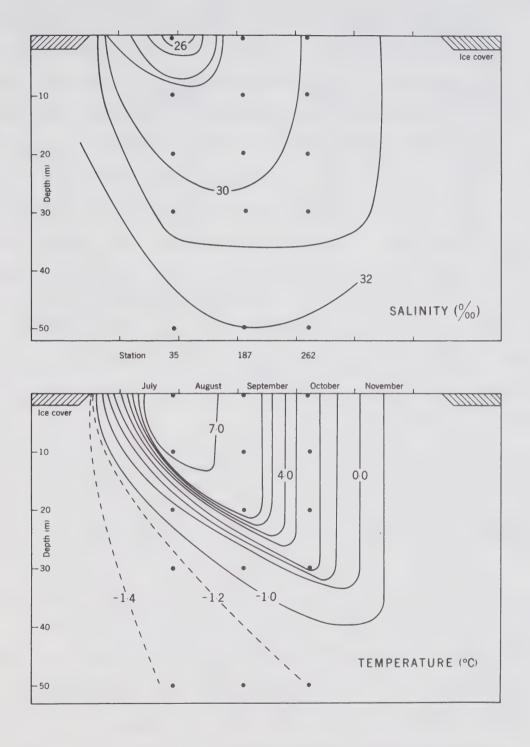


Figure 3. The annual sequence in the near-surface at a location in Hudson Bay (from Barber, 1967, p29). (a) Salinity. (b) Temperature.

Table 1 A possible distribution of salinity and temperature under the ice cover at a point just west of Trodely Island.

Depth (m)	Temp.	Sal. (0/00)
2	0.0	0
4	0.0	0
10	0.4	15
20	-0.8	25
30	-1.3	27
50	-1.3	27
<b>7</b> 5	-1.3	27

#### 1.3 Bathymetry

Most knowledge of the bathymetry (frontispiece) is contained in Canadian Hydrographic Service chart number 5800 and pilotage information is available in another publication of the Service (Anon., 1965) from which the following was extracted (p441):

The low, flat shores of the bay are fringed with wide mud flats. The bay is filled with numerous islands, rocks and shoals, with no harbours for large vessels. In the fairway which, in the outer part, lies close westward of the middle of the bay, there are general depths of from 10 to 40 fathoms (18m3 to 73m2) to within about 20 miles of the head of the bay.

As in the case of Hudson Bay itself, the eastern shore of James Bay is very irregular in outline, with many islets and rocks lying close off it, while the western shore is even and almost free of islands. There is, however, Akimiski Island lying close off the middle of the western shore. In the eastern half of the bay are many islands and shoals. The western shore of the bay affords no harbours, but there are several fairly good harbours on the eastern shore for vessels drawing up to 12 feet (3<sup>m7</sup>) of water.

Limiting depths exist in certain parts of the bay so that the water in some areas may be isolated from that at similar depth in Hudson Bay. For example, it is thought that for a station occupied in the deep water off

Trodely Island a limiting depth of 40 m exists so that a salinity of about 27 % oo could be expected there. The actual situation is quite uncertain and it is possible to interpret available bathymetric data as though there were shallows, less than 20m, extending continuously off the east shore from Cape Jones (Pte. Louis XIV) to Bare Island southward through Grey Goose Island, the Twin Islands, Weston and Charlton Islands to the east coast in the vicinity of Rupert Bay, so that a depth of 60 m would be isolated to the eastward. However, the interpretation here (frontispiece) indicates just one such area eastward of the shallows between Grey Goose Island and the Twin Islands, but with deeper areas south of Bare Island and north of Weston Island. Depths of 80 m occur almost as far south as Trodely Island; a depth which it was noted is not likely continuous to Hudson Bay. The areas of limiting depth of about 40 m are believed to exist between Bear Island and Akimiski Island and to the east of South Twin Island.

With regard to the water which may occur in the northern approach to James Bay, oceanographic data indicate that a limiting depth of between 50 to 75 m occurs between the Belcher Islands and Cape Henrietta Maria (Barber, 1967, pl5), thus the deeper water north of Pte. Louis XIV would be slightly isolated from Hudson Bay (an interpretation of more recent, but still incomplete, data on Canadian Hydrographic Service field sheet 025A suggests a limiting depth of 80 m) so that the range of salinity observed at

depth in the approach, i.e. 31.7 to 32.0 °/oo, is somewhat less than is observed in the open areas of Hudson Bay generally, i.e. up to 33.0 °/oo at 100 m (Figure 4).

1.4 Surface wind

Knowledge of surface wind within the system is based largely on observations at coastal stations, which indicate that winds are strong in all but the summer months when they are:

...generally lighter and are variable in direction with a higher proportion of onshore winds at coastal stations, the effect of local sea-breeze circulations.

(Anon., 1965, p395).

Danielson (1971, p97) remarked on the influence of persistent north winds which move ice southward and evidence provided by Archibald (1969) suggests that wind strength may decrease from north to south, being less in southern James Bay than elsewhere in the system. the author's (limited) experience that storms can occur for short periods at any time during the summer with winds from the southeast to southwest. At Churchill in late September strong northwest winds may persist for 4 to 5 days and in early October in 1961 in western Hudson Strait strong wind (290°, 50 knots) was experienced. Occasionally it is possible to see in the oceanographic distributions, correlations with surface wind. Of particular interest here is the influence of wind during the summer, i.e. when it is weak and variable with a "high proportion of onshore winds". As cloud and fog are frequent over the water such



Figure 4. The distribution of salinity at the surface
(a) and 50 m (b) in Hudson Bay and James Bay
(from Barber, 1967) and the distribution in
Hudson Bay of salinity (c) and temperature (d)
at 100 m (from Barber and Glennie, 1964).

a wind would tend to increase the frequency of cloud and fog at coastal stations; an increase in each has been observed in both Hudson Bay and James Bay (Thompson, 1968). However, the main movement of air at the surface is from west to east so that a relatively warm air mass is moved over the southern-most parts of the system. It is useful to consider the likely surface temperature toward the end of July based on the 1961 experience (Figure 5) by which time considerable warming of the water had occurred over northern Hudson Bay, but with little or none in the southern part due to the ice cover there. Some warming has also occurred on land by this time, for example to the southwest (Figure 6). A warm air moving west to east from the land over the ice covered area would likely become relatively stable with which would be associated a decrease in surface wind. The implication is that this small-scale region of atmospheric stability through a decrease of surface wind, increase of cloud and fog and with the high albedo of ice cover, contributes to the stagnation of the ice in the southwest usually observed (Anon., 1970e). Also implied is a lack of water movement due to other factors; however, this may result as a consequence of the James Bay circulation (Murty, 1971).

#### 1.5 Freshwater from runoff

It seems likely that the pattern of runoff to James Bay will continue to be influenced by man, either through the increasing development of hydroelectric potential as has already occurred to some extent on the

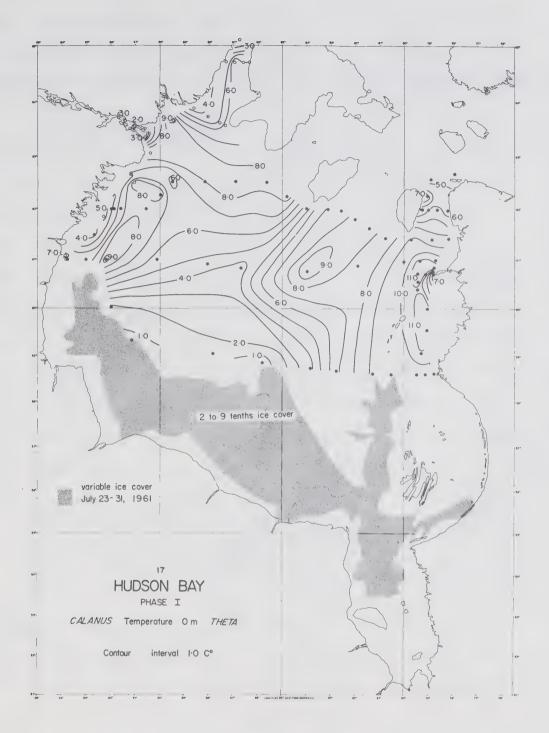


Figure 5. The likely distribution of surface temperature in Hudson Bay toward the end of July based on temperature data (Barber and Glennie, 1964 Figure 17) and ice data (Anon., 1962) in 1961.

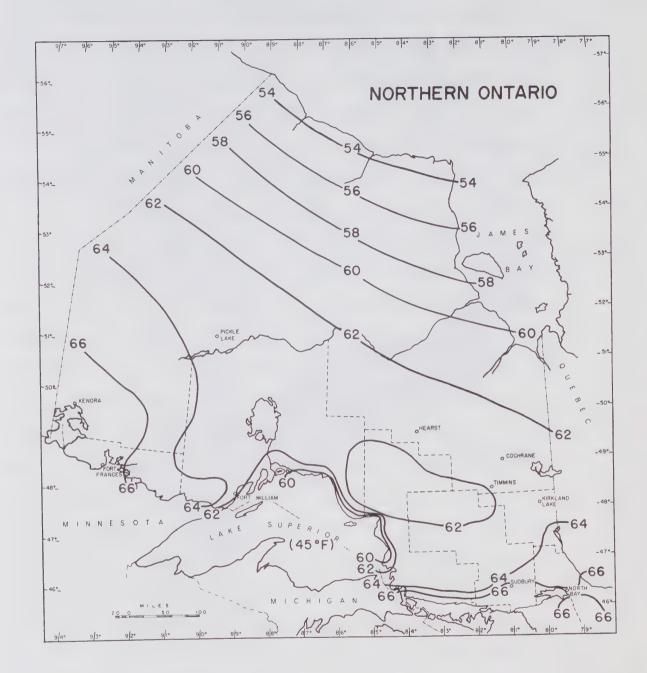


Figure 6. Mean daily temperatures (°F) in northern Ontario for July (Chapman and Thomas, 1968 their figure 8).

Moose River (Robinson, 1968, p219), or perhaps through diversion of water southward (p219). At the present time the annual runoff volume appears to be about  $3 \times 10^{11} \text{m}^3$ , that is if there were no diversions. The value was obtained by extrapolation of the data on runoff from the Québec rivers (Figure 7) over the total drainage basin including James Bay. A check was obtained through the estimation of a value for the annual excess of precipitation over evaporation (P-E) over the total region. Hydroelectric development, in the absence of diversion, would not be expected to influence the annual value appreciably (except that it modified P-E), but would markedly influence the present pattern of runoff (Figure 8), which varies from very high levels in spring to extreme low levels in March and occasionally in summer. Black remarked (1968, p843):

The rivers become quite shallow by July, even during wet weather. At this time the low water exposes numerous bars, bouldery shoals and rocky outcrops which render stretches of river difficult to navigate, ... A short period of navigation is possible before freeze-up when a second high-water period occurs during heavy autumn rains, ...

Presumably the rivers are not used for navigation during the period of low runoff, i.e. in March, at which time the flow can vary between a fifth to more than an eighth of the much more variable spring volumes. No doubt this large spring runoff has a significant influence on conditions in James Bay, including the distribution of ice, salinity and currents, particularly as it occurs when the ice cover is still extensive. However, even though

Table 2 A tabulation of 1968 surface water data from six stations indicating the name, size of basin, mean flow and the page reference to the source (Anon., 1971f).

Other data are available (Anon., 1967a; 1970b; c; d).

Name (St. No.)		Area (Sq. Miles)	Mean Flow (cfs)	Page
La Grande-Rivièr (092704)	îe	37,600	60,000	363
Eastmain (090601)		17,100	31,400	352
Grande-Rivière d la Baleine (093803)	le	16,500	23,700	379
	sub-total	71,200	115,100	
Rupert (081002)		15,800	30,300	345
Broadback (080801)		6,610	11,500	343
Nottaway (080701)		22,200	36,600	333
	sub-total	44,610	78,400	
	TOTAL	115,810	193,500	

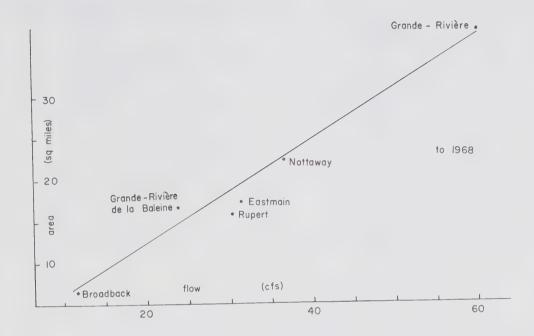


Figure 7. A presentation of the data of Table 2. The extrapolation referred to in the text is based on the "best fit" of the straight line.

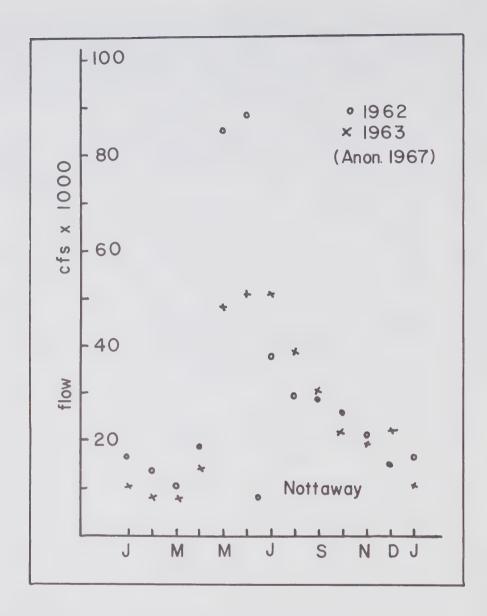


Figure 8. A presentation of the monthly values of runoff for the Nottaway River for 1962 and 1963 (Anon., 1967a).

data on the flow from the Québec rivers are quite adequate, lack of oceanographic data limit their application. For the purpose here the significant annual values appear to be:

- 1) amount from rivers and
   excess of P-E over the
   bay
   but not including
   Grande-Riviére de la
   Baleine
- 2) amount from rivers
  to be developed
  i.e. the sum of the
  southern rivers
  and of the
  northern rivers

 $3x10^{11}m^3$ 

0.02x10<sup>11</sup>m<sup>3</sup>

1.88x10<sup>11</sup>m<sup>3</sup>

0.78x10<sup>11</sup>m<sup>3</sup>

1.1x10<sup>11</sup>m<sup>3</sup>

Thus the development will increase the total annual value slightly and markedly influence the pattern for about half the runoff to the bay. With the potential for development of the southern and west coast rivers the runoff to the bay could be almost completely smoothed or at least reflect a demand for energy, which may be maximum in winter, i.e. at a time when the runoff prior to development, was at a minimum.

### 2. The system

# 2.1 Data in James Bay

It is curious that James Bay has not been examined by oceanographers to nearly the same extent as have areas much further north. Apparently it is considered a difficult area, particularly for vessels suited to ocean survey, so that during the development of the 1961 oceanographic programme

for Hudson Bay the work in James Bay was limited to two sections across the entrance. Earlier, the negative results of an investigation into the fisheries potential of Hudson Bay by Hachey (1931a) considerably deflated interest in the region generally (Hunter, 1968, p373). The 1961 material and the 1959 work of "Calanus" (Table 3) constitute the only significant data yet observed there, although temperature and salinity data were observed in the Moose River with tidal observations (Langford, 1963, p90), but these have not been located. All of the available material were observed during the period of open water, though some of the "Calanus" stations were occupied at a time (Table 4) when some ice cover occurred (see data report for existence of ice nearby at the time the station was occupied). Distributions based on the "Calanus" data were presented by Grainger (1960) and these were also utilized in a description of Hudson Bay (Barber, 1967; 1968a). The recent compilation of information about the region (Beals, 1968) proved useful in a number of ways, for it is "an impressive assemblage of facts not readily found elsewhere" (Jackson, 1970, p841).

## 2.2 Review

James Bay is part of a general system (Figure 1), comprising it, Hudson Bay and Foxe Basin, connected to the world ocean through Fury and Hecla Strait and Hudson Strait. Fury and Hecla Strait is likely much less important than is Hudson Strait to the character of the water within the system, for not only is Hudson Strait relatively wide

Table 3 A tabulation of the data available indicating the year obtained, name of ship, cruise reference number and a reference.

Year Ship		CRN	Reference	
1958	"Calanus"	320	Grainger, 1960	
1959	do	321	do	
1961	"Theta"	337	(Anon., 1964a; b)	

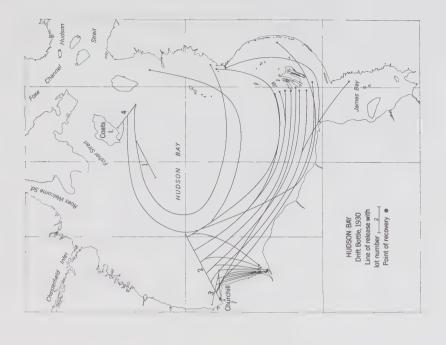
Table 4 A tabulation of the "Calanus" 1959 data in James Bay indicating the station number and day. The latitude of station 59-4 has been taken as 52°30' and that of 59-15 to be 53°56'.

Number	Day	Number	Day
1 2 4 6 7 8 9 11 13 15	June 20 22 22 23 25 26 27 30 July 1	57 58 57* 59 60 61 62 63	Aug. 26 26 29 29 29 30 30 30

<sup>\*</sup>Apparently a re-occupation.

and deep, but also strong mixing due to tides occurs there. Thus the water reflects to large extent an influence of the Atlantic Ocean and an influence of a process occurring within the system, and even though James Bay is 600 to 1,000 miles away from these influences they are still strongly evident there. In the present context however, distance need not be particularly significant, for factors such as the general location and the associated climate and pattern of runoff in such a uniquely shallow region as James Bay can be paramount.

The foregoing assumes certain physical properties of water which are of prime significance. For example, at salinities generally encountered in the ocean a temperature of maximum density does not exist at any temperature warmer than the freezing point (Figure 9), as exists for freshwater at 4°C. It will be shown that much of the water is less than oceanic salinity, specifically less than 24.7 0/00, so that a point of maximum density can occur at temperatures warmer than that of ice formation at the surface. At such low salinities, surface cooling at temperatures close to freezing would lead to stability. Furthermore, at those locations where the surface layer is fresh or nearly so, the amount of salt made available to the surface layer through freezing is negligible. Thus, the evolution of a surface mixed layer through the winter would not relate to the formation of denser water at the ice-water interface (Barber, 1968b). Also of utmost significance, here as elsewhere, is the

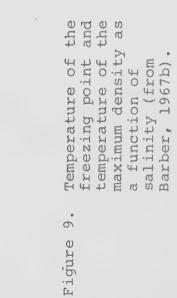


Temperature (°C)

Temperature of the freezing point and temperature of the maximum density as Barber, 1967b). salinity (from a function of 6

A re-presentation of the drift bottle data of Hachey (1935), (from Barber, 1967).

Figure 10.



Salinity (0/00)

20

10

fact that ice formed at the surface remains there, as it is less dense than water, providing a very different boundary than is the air-sea interface.

#### 2.2.1 Currents

It may be significant that only one of the drift bottles released in 1930 (Hachev, 1935) was returned from James Bay (Figure 10). If it is significant, then it is envisaged that the inflow from Hudson Bay occurs at depth rather than the surface and comprises a cold, relatively saline water. On the other hand, Grainger (1960) inferred an anti-clockwise surface movement with water entering from Hudson Bay along the west coast. Data observed in the autumn (October, 1961) suggest the development of a well-defined movement from James Bay north along the east coast of Hudson Bay and to considerable depth (for example see salinity at 100 m, Figure 72 of Barber and Glennie, 1964), which was partly due to an increase in the amount of freshwater there in October over August (Figure 11). Thus the pattern of circulation in these areas may exhibit a time-dependence in association with the distribution of freshwater from runoff.

Presumably a portion of this runoff would have absorbed a significant amount of shortwave radiation and would enter the bay at a temperature above that of the surface water there, so that it would constitute a gain of sensible heat. Presumably too, much of the outflow at the surface would have, in the absence of ice, absorbed heat and would constitute a loss. The heat transported

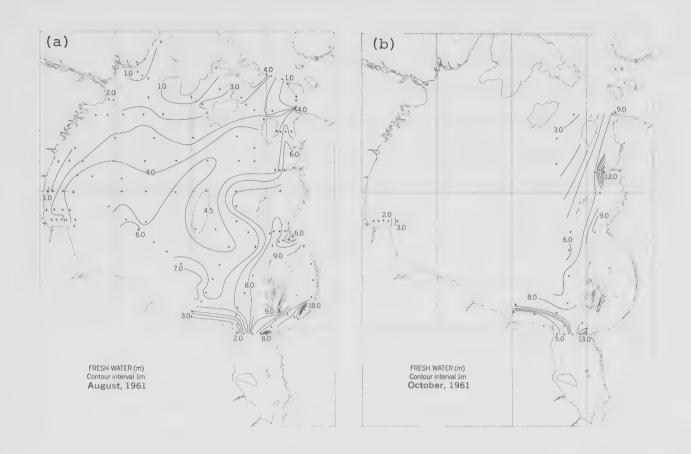


Figure 11. The distribution of the depth of freshwater (m) in August and October, 1961 (from Barber, 1967). (a) August. (b) October.

with the inflowing deep water would be very much less so that in the section across the entrance a loss of heat occurs due to transport. The loss may be balanced by an excess of export over import of ice, but there is no direct evidence that either occurs, and it may be that it is balanced by the input of heat with the runoff. In this circumstance the annual sum of the flux terms would on the average be zero. Danielson (1969, pl62) considered that a variation in the flux occurred over the bay ranging from a deficit of about 100 g cal cm<sup>-2</sup> year<sup>-1</sup> in the north to about 6000 in the south and leading to an average gain of about 2500 g cal cm<sup>-2</sup> year<sup>-1</sup>. A rough heat budget for a position in the northern part of the bay was attempted and although a value which might be considered better than Danielson's was not achieved, a comparison of the actual with the potential heat storage for two periods was made from existing data. Between August 11 and 29 (1959) an increase in heat storage of 6500 g cal cm<sup>-2</sup> occurred, which is about twice the amount to be expected in the absence of advection. This suggests a movement out of the bay of the warmed water. During the period August 29 to October 2 the mixed layer depth increased to 25 m and the heat storage was unchanged. This would have been expected in the absence of advection. The data therefore suggest a variation in a movement of surface water from the bay, again suggesting a time-dependence in the outflow.

Such a time-dependence could be expected to have a distinct influence on the pattern of surface movement

seaward of James Bay, but direct evidence is not available. Some secondary evidence has been mentioned, for example drift bottle data and the distribution of freshwater, but there is one other which should be noted even though quite speculative. This relates to the distribution of ice within Hudson Bay during the summer which exhibits lingering ice, or a pattern of last ice, in the southwest (Anon., 1970e; Danielson, 1971) due in part to movement of ice from the north during the early part of the breakup period. Later the ice appears to stagnate in the southwest with little or no tendency to move and, in particular, no tendency to move with anti-clockwise movement around the coasts presumed to exist.

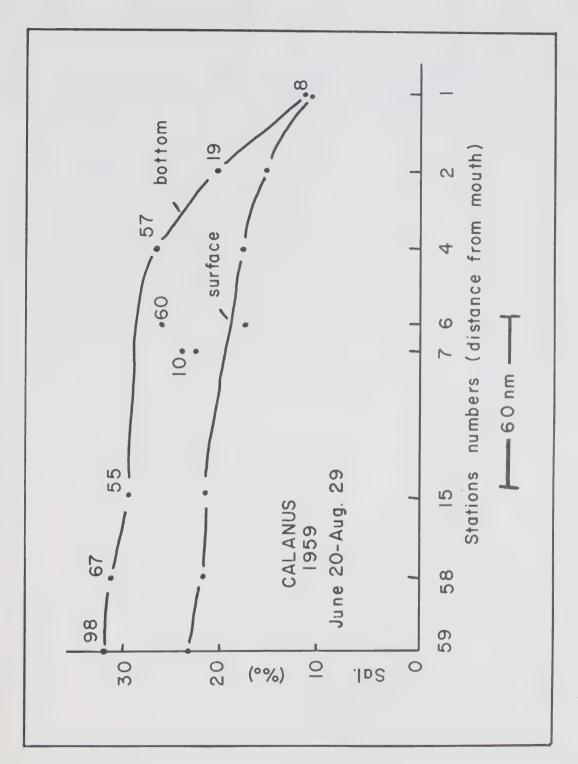
## 2.2.2 Salinity

Various distributions based mainly on the 1959 data are included in the final group of figures (Figures 21 to 24). Generally, the material was difficult to interpret and in some distributions certain of the data are not included.

The distribution of salinity at the surface and at 50 m (Figure 4) within Hudson Bay and James Bay is based on both the 1959 and 1961 data. Surface salinity was generally low and ranged from 32.5 % on in northernmost Hudson Bay, to 27 % on in the approach to James Bay and to less than 10 % on at the head of the bay, most of the decrease occurring in the southern part. Salinity values in excess of 31 % of owere observed at 50 m in the approach

to the bay and a marked north-south gradient extended well into the bay.

The 1959 data are utilized in the illustration of the longitudinal distribution of surface and bottom salinity (Figure 12) which emphasizes both the strong gradient at the surface toward the head and the marked vertical stratification or structure. It seems likely that stratification is a permanent feature of the oceanography of the bay; a feature which may become more pronounced seaward of the major rivers in winter to the extent that freshwater occurs under the ice. Such a circumstance has been observed in Kugmallit Bay (Barber, 1968b) just east of the delta of the Mackenzie River and in Tuktoyaktuk Harbour. The spatial distribution of the freshwater layer seaward of Kugmallit Bay was not observed, but it was visualized that the freshwater moved eastward close to the coast within a surface layer and with little mixing with deeper salt water. As entrainment of salt was minimal the normal estuarine type of circulation did not exist and it was possible to discern an exchange process due to tides. In James Bay, however, the rms value of the tidal current appears somewhat larger than in Kugmallit Bay so that tidal mixing may be considerable. This would tend to remove stratification so that the layer of freshwater suggested here may exist only in the immediate vicinity of the various estuaries. Nevertheless, in order to emphasize the lack of data and the considerations in



A longitudinal distribution of surface and bottom salinity based on "Calanus" data during the period June 2 to August 29, 1959. The depth in metres of the near-bottom sample at each station is indicated. Figure 12.

other sections about the coupling to Hudson Bay, the layer is shown (Table 1) as extending to Trodely Island.

It was suggested earlier (section 1.2) that in the approach to James Bay a surface layer probably existed within which seasonal changes of temperature and salinity could be recognized; the salinity change was due to variations in the annual ice cover. Observations at about one location in a section across the mouth of the bay (Figure 13) support the view that such a change occurs, but the extent to which it reflects the influence of the annual ice cover as opposed to annual runoff is not known, expect that at positions in the section close to the west and east coasts, a variation in salinity occurred which is believed due to runoff (Figure 24).

## 2.2.3 Secchi disk depth

The distribution of Secchi disk depth in James
Bay in 1959 and Hudson Bay in 1961 during the period of
navigation of those years is shown in Figure 14. Least
values occurred in James Bay where depths less than 2 m
were observed in the south and where in general the depth
appears to have been less than 5 m. In the approach to
James Bay values to 10 and 15 m occurred. Over much of
Hudson Bay readings were close to 15 m with a significant
number to 20 m and an occasional (2) reading to 25 m.

The relatively low values in James Bay may be due to a sediment load associated with the large inflow of freshwater from runoff, although other factors could be of equal or greater significance. Langford (1963)

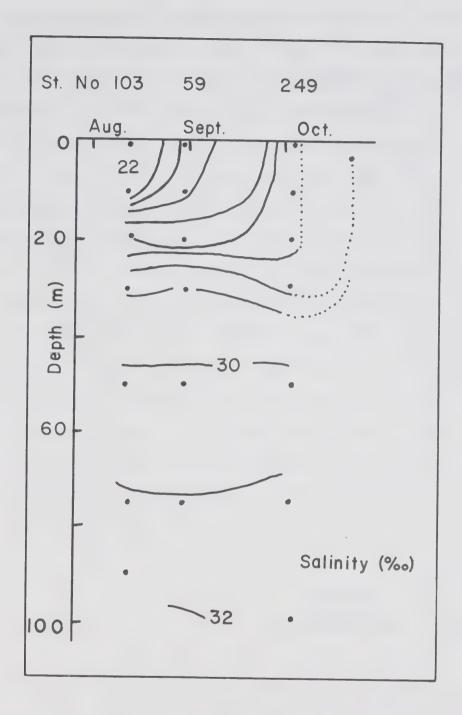


Figure 13. The distribution of salinity at about one location in the mouth of James Bay derived from "Calanus" station 59 in 1959 and "Theta" stations 103 and 249 in 1961.

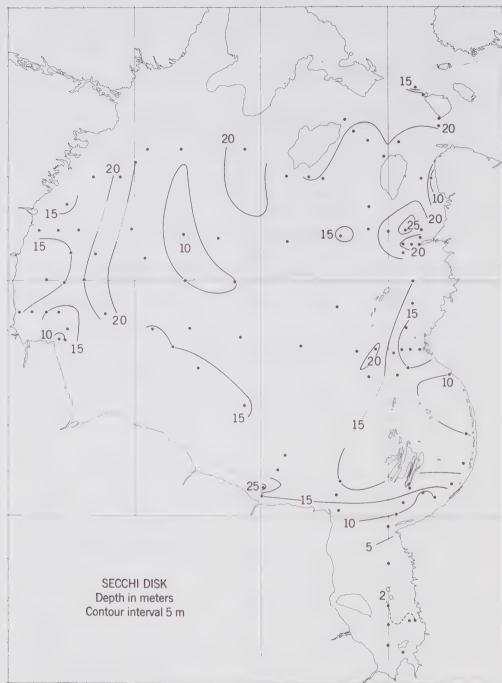


Figure 14. The distribution of Secchi disk observations in James Bay and Hudson Bay based on data observed in 1959 (Grainger, 1960) and in 1961 (Anon., 1964a). The solid circles indicate the location of the observations. A dashed contour indicates one additional to the regular contour interval; a dotted contour indicates a doubtful interpretation.

considered that as the sediment load of the Moose River appeared small in summer months it must be large during the spring freshet. The north-south gradient at the mouth of James Bay does not appear to be reflected in the distribution of either the surface salinity (Figure 4) or the freshwater content (Figure 11) as might be anticipated. A tentative conclusion is that the relatively low Secchi disk readings in James Bay represent the influence of that portion of the input of freshwater derived from runoff\*, as opposed to that derived from the ice cover of the previous winter; the gradient across the mouth would represent the "front" of the seaward moving freshwater from runoff. This suggests that the stored volume of freshwater within James Bay derived from rivers entering there may be distinctly time-dependent. Secchi disk depth would then be time-dependent, and so might sediment deposition.

<sup>\*</sup>It was anticipated that the optical difference of the sea surface caused by the sediment laden runoff into James Bay during the spring and early summer might be visible in APT optical photography. While the available data are not sufficient to allow a study of this, it does seem that the contrast is so slight that a muddy water would not be recognized. On the other hand, Taggart et al., (1965, p190) were able to detect tone variations in the APT photography of the bay in September which they related to the shallow water there.

The latter feature could lead to the existence of annual varves in the sub-bottom vertical sediment structure. Leslie (1965, pl36) studied the sediment core obtained at "Theta" station 104 in the approach to James Bay but did not indicate the existence of varves. However, he did suggest that the bay is an important source of finer material in the bottom sediments of Hudson Bay. He remarked (p18), "North and east of James Bay the bottom sediments consist mainly of medium grey silty clay," and (p20), "James Bay is the source of medium grey sediment along the southeast coast and around the Belcher Islands. Rivers flowing into James Bay drain the region to the south which is underlain by soft Mesozoic shales and siltstones. Much of this fine detritus is carried into James Bay and thence into Hudson Bay". This distribution of sediment finer than 2mm in diameter as deduced by Leslie (1964) is shown in Figure 15. The area of silty clay, which apparently originates in James Bay, extends to the north and east and the westward in the northernmost part of its distribution. The distribution is compatible with present knowledge of the circulation; the main feature of which is believed to be an anticlockwise movement around the bay with relatively strong northerly currents along the east shore. The westward extension may reflect the

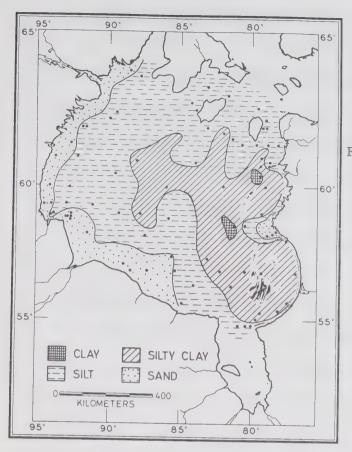


Figure 15. A re-presentation of Figure 8 of Leslie (1964) the caption to which read, "Distribution of bottom sediment finer than 2mm in diameter, or the predominately waterdeposited material".

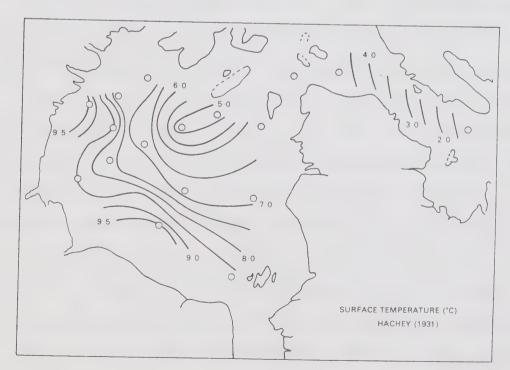


Figure 16. A reproduction of the distribution of surface temperature in Hudson Bay in 1930 (from Hachey, 1931b his Figure 8).

influence of a pattern of water re-circulation\* within Hudson Bay.

2.3 Some recent observations within the system

Wendland and Bryson (1966; 1967) obtained surface temperature data in Hudson Bay using airborne radiation thermometry and concluded, "A signature of last ice of Hudson Bay apparently can be located throughout the remainder of the ice free season". They attributed the persistence of the feature to the "Stratification of the upper layer of the Bay"; a stratification in salinity due to melting ice. It seems that such stratification could have quite the opposite effect and also it seems that other factors could be important including the ice cover itself. Nevertheless, the existence of such a "signature" in the temperature distribution at the surface appears to be entirely possible as other summer data reflect such an influence and it is known that each year a characteristic pattern of ice dissipation, from north to south, can be expected (Markham, 1962, p6; Larnder, 1968, p335; Danielson, 1971). The 1930 data of Hachey (1931b) are an exception, for not only is a signature not apparent (Figure 16) but also the surface temperatures, particularly in the west and

<sup>\*</sup>A (drift) bottle set adrift in 1952 at a position off the northeast coast of Hudson Bay, 10 miles west of Povungnituk, was found in 1968 on the southwest shore, 4 miles northwest of the Kaskattama River (Richard H. Russell, personal communication).

in the south, are relatively high. Of the 1930 data Barber (1967, p7) remarked, "... There does not appear to be any effect due to a recent ice cover or accumulation of ice as in the 1961 season". He also remarked (p7) that 1930 may have been a light ice year, but this is quite speculative as the data (Anon., 1931) on the distribution of ice are few. Of these the second voyage of the tug "Ocean Eagle" during July 10 to 18, 1930 suggests a scattering of the ice (pl4 and 15) eastward of Churchill rather than a concentration. If a subsequent movement did not concentrate\* the ice in 1930 it is possible that a condition similar to the unusually open season of 1962 occurred. In this, the main portion of the bay seems to have been effectively clear of ice by August 1, (Anon., 1963) so that there would have existed sufficient time and open water for the surface waters to have been warmed by insolation in 1930 to the extent indicated by Hachey.

<sup>\*</sup>Ice was observed off Port Harrison on August 8 and 9 in SS "Nascopie" (Anon., 1931, p19) and "loose ice" was observed in SS "Ungava" (p21) on August 2 at a position (57°06', 82°58') west of the northernmost Belcher Islands and north of Cape Henrietta Maria on a voyage to Charlton Island.

Both support the considerations that the ice condition in 1930 was similar to that of 1962 and not similar to either 1967 (Anon., 1969, Figure 16) or 1969 (Anon., 1971e, Figure 16).

It is a tentative conclusion that in some years the "signature of last ice" may not be apparent.

That 1967 would not likely be such a year was predicted from study of satellite imagery and of ice forecasts (Anon., 1967b) which indicated that an open water situation was well-defined in northwest Hudson Bay and close along the east coast by about May 11. Subsequent imagery indicated increasing open water in the northwest and little or no open water along the east coast where it had been earlier observed. The latter change suggested a movement or "pressure" from the west to east. As well, there was little evidence of open water in the eastern half of the bay prior to the end of June; indeed, some ice persisted in the area south of Coats and Mansel Islands to the end of July (Figure 17). An open water condition along the west coast toward the Bay of Gods Mercy was interpreted from the imagery of April 28; a condition which has been described (Dunbar and Greenaway, 1956, p418; Bowley, 1969, pl3). Another feature was the persistence of an ice boundary in Roes Welcome Sound throughout the period April 28 to July 12 (Figure 20a). In 1961, aerial reconnaissance indicated open water throughout Roes Welcome Sound on July 13 (Archibald et al., 1962, their Figure 8) and the relatively high (25 kg cal  $cm^{-2}$ ) seasonal heat storage to about the latitude of Wager Bay (Barber, 1967) indicated that open water had occurred there early in the season and persisted up to the time of the temperature observations. Clearly, the heat storage in the area in

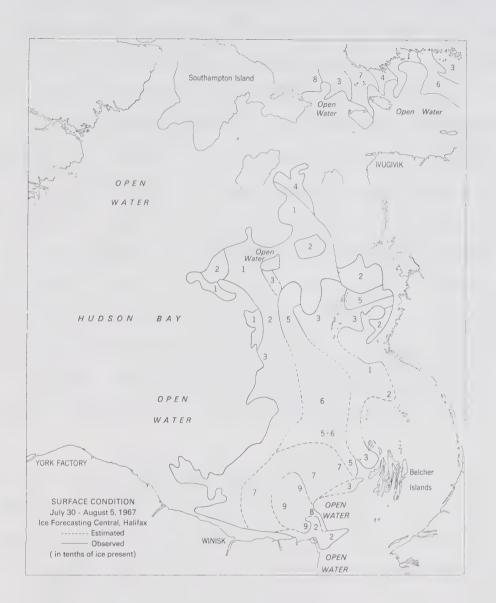


Figure 17. The distribution of open water in Hudson Bay about the end of July, 1967 from Ice Forecasting Central, Department of Transport, Halifax.

1967, if observed, would have proved much less than in 1961. Similarly the slow development of open water in the east suggested that the absorbed incoming shortwave radiation, and hence the surface temperature and seasonal heat storage of the water, would be much reduced. This prompted a request to the Department of Transport for bathythermograph observations in the northeast. These were obtained by CCGS "Labrador" (Anon., 1968) and revealed that the seasonal heat storage was low (Figure 18), certainly much lower than occurred at about the same time in 1961 (Barber, 1967). The subsequent airborne survey in late August (Wendland and Bryson, 1967, their Figure 3) indicated a relatively low surface temperature, which confirmed that in the eastern portion of Hudson Bay the peak of the seasonal heat storage was less than average. Hence their conclusion concerning the "signature of last ice".

On the other hand surface temperature data observed between the Bay of Gods Mercy and Churchill about mid-July in 1967 (Figure 19) indicated a level of temperature close to that observed about the end of July in 1961 (Barber and Glennie, 1964, their Figure 17). If it is assumed that the development of the surface layer was similar in the two years, then it is a tentative conclusion that the peak value of the seasonal heat storage there was at least as great in 1967 as in 1961.

More recently, the 1969 season in terms of development of open water also appears to have been

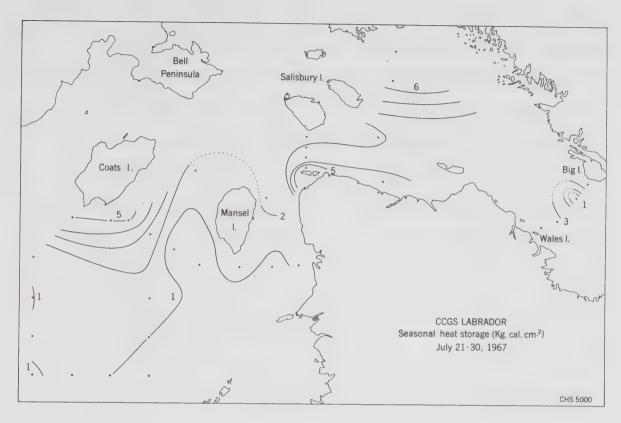


Figure 18. The distribution of the seasonal heat storage in late July, 1967 as interpreted from bathythermograms obtained in the C.C.G.S. "Labrador".

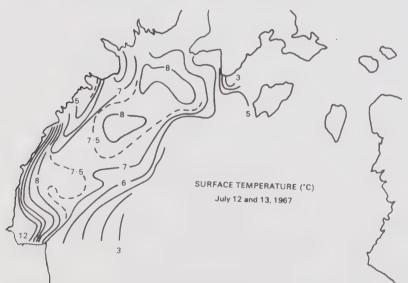


Figure 19. The distribution of sea surface temperature as observed July 12 and 13, 1967 using an airborne radiation thermometer (from Wendland and Bryson, 1967, their Figure 1).

anomalous for, "in southern Hudson Bay final clearing after mid-September was the latest on record" (Anon., 1971e, pl); the persistence was related to the occurrence of "some heavier than usual ice" (p42). At the same time the existence of, "Broad areas of open to close pack ice," west of Belcher Islands to August 20 (p42) was also considered unusual as was the, "Break-up and disintegration of ice which was earlier than normal in northern Hudson Bay..." (pl). The distribution of ice in May and June strongly suggested the influence of advection due to wind and examination of the data indicated that winds from the west through to north occurred 75 percent of the time at Chesterfield, i.e. about 25 percent greater than normal. This pattern of wind no doubt contributed to the earlier open water in the north and to the observed accumulation of ice toward the south and east. It is suggested that the accumulation here was such as to maintain a high average albedo well into the normal melting season and was a factor in the persistence of the ice cover there.

The deeper temperature data of 1967 (Figure 20a), while limited, suggest the existence of a slighly warmer water, to -0.9°C, in a tongue-like distribution southwest of Coats Island similar to that observed in September 1962 (Figure 20b). Such warming is thought to be due to downward mixing of heat absorbed at the surface in that season, but was not considered to occur as early as observed in 1967, i.e. as early as the end of July. It seems possible that the extensive open water in 1967 in the northwest

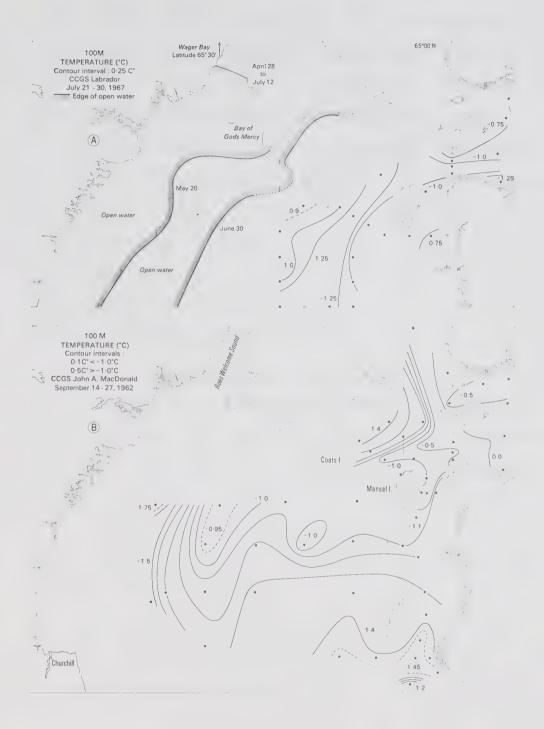


Figure 20. The temperature distribution at 100m depth in 1967 and 1962. (a) From the 1967 "Labrador" bathythermograms. The edges of open water at the dates indicated are an interpretation by the author of satellite (APT) imagery. (b) From the 1962 "John A. Macdonald" serial data (Anon., 1966).

may have contributed to this result which, due to the similarity in open water, probably occurred again in the 1969 season.

There is one major caveat here concerning the 1967 data and that is they are the result of but one instrument, i.e. one bathythermograph, so that both the values and pattern (Figure 20a) are suspect. The chances are, however, that the bathythermograph data are appropriately used here relative to the assessment of heat storage. It is realized that this may not be the case for all data observed within the system, for certain of the apparent differences are disconcerting. The most recent example is that of Pelletier et al. (1968) who, even though aware of earlier interpretations, could report for Hudson Bay, temperatures to -2.0°C (their Figure 4 and p565) and salinities in excess of 34.0 0/00 (their Figure 6\* and p565) without particular comment, either to the precision of the observations or their relevance. extreme values were reported on earlier occasions, but at a time when data were relatively few. For example Campbell (1964, p49) reported temperatures to -1.97°C and salinities to 34.07 % oo at unusually shallow depth in Foxe Basin.

<sup>\*</sup>Their figure 6 indicates a value of 35.56 O/oo which may be an error in the preparation of the illustration.

The existence of such extreme value generally leads to conjecture about the process which forms the water; a process which may be a major feature of the oceanography of the region. A consideration of all the data led to the definition of the probable winter surface salinity (Barber, 1967) and to a simple model (p54) of the Hudson Bay system. The model also served as a review of earlier concepts and emphasized that these were hampered through lack of winter data. Hachey (1931b) realized this and developed the view that, "the waters of Hudson Bay differ markedly from the waters of Hudson Strait and the waters of the open ocean". Of course he was not aware that conditions in Hudson Strait exhibit marked seasonal variation. He was led (Hachey, 1954) therefore to consider the evolution of Hudson Bay water alone and offered two conjectures. In one was visualized removal of the observed summer low salinity surface layer prior to freeze-up, so that water cooled at the surface would sink and become the deep water. In a variation, he visualized that water initially at the surface in winter in Hudson Strait moved into Hudson Bay replacing a less-dense seaward flowing water. In the other he suggested a movement "of deep, cold, saline water" into Hudson Bay from Foxe Basin; the origin of such water remained a question. Campbell (1964) expressed his accord with the latter hypothesis, describing observations made in 1955 and 1956 in support, and developed a theory wherein the cold high-salinity water in Foxe Basin was directly related to the production of winter ice in

the extensive inter-tidal zone (tide flats) of Foxe Basin. Apparently he considered the "exposure" of the tide flats the significant feature in the process and visualized the production of a "concentrated brine solution". No primary evidence which might lend credence either to the existence of such a brine solution, or to the main hypothesis, was presented. However, the described data are important as they demonstrate the existence below sill depth in Foxe Channel in both years of a high salinity water (about 33.7 O/oo) close to the freezing point. This conditioning must have occurred at the surface in association with cooling during the winter months. The water so conditioned would sink and move away, if not depth limited, from the source area eastward into Hudson Strait, and into Hudson Bay, enhanced perhaps through increased heat losses in areas of open water, including tide flats. The extent of open water in the system at this time is not known; however, amounts likely significant in the formation of the water are believed to occur in northern Foxe Basin, throughout Hudson Strait, Foxe Channel, and Frozen Strait to Roes Welcome Sound.

Looking beyond the system, it is suggested that the process would be favoured by an extension of the Baffin Current into Hudson Strait along the north shore and into the area of Foxe Channel. The evidence for such a current during the winter is limited, but it is the hypothesis that it exists and in association with a process similar to that described by Mosby (1934) for the formation of

a cold, saline water on the Antarctic continental shelf, particularly in the Weddell Sea. In this, he elaborated the conclusion of Brennecke (1921), that the deep and bottom water of the Antarctic is derived in part from an extremely cold and moderately saline water formed in winter through cooling at the surface and freezing on the shallow shelf. Mosby emphasized the importance of horizontal movement whereby the surface water sinks and moves away from the source area, and off the shelf. The similarity of the situation in the Antarctic with that observed in the Foxe Basin - Hudson Bay system is heightened for according to Mosby, a portion of the winter-formed water is prevented from mixing into the Antarctic deep through the existence of a limiting depth. Thus, it retains its main characteristics, i.e. like the deep water in Foxe Channel a temperature very close to the surface freezing point. It was these data, specifically those at "Deutschland" station 125, that led Brennecke to the hypothesis concerning the formation of the winter shelf water, and its subsequent contribution to the deep water. Deacon (1937) and Fofonoff (1956) have contributed further to the understanding of the process.

In the situation for Hudson Bay it is thought that the observed cold and saline water at depth in Foxe Channel is the result of the cumulative effects of cooling and ice formation at the surface of the current throughout its movement along the east Baffin Island coast, into Hudson Strait, and into Foxe Channel and Foxe Basin. It

is not recognized as a distinct water outside the system in the Labrador Sea because it is too dense to remain at the surface and is not sufficiently dense to participate in the formation of a distinct bottom water. It must then contribute to an intermediate water there.

An assessment of this contribution will not be attempted here, although it is of major interest, for it seems that much more needs to be known of the influence of Hudson Strait, where intense mixing due to tides leads to a further modification of water characteristics. These subsequently became recognizable within Hudson Bay through an inward movement associated with the estuarine circulation, which in effect results in considerable recirculation, i.e. the coupling of the system to the Atlantic through the estuarine circulation is limited by mixing in Hudson Strait.

Assessments of plankton data by Grainger (1961) and Bursa (1961) showed that both Atlantic and Arctic plankton types were found in Hudson Bay, and Grainger (1962) made a similar observation with regard to Foxe Basin. In a study of the distribution of three species of copepod, Grainger (1963) showed that while two Arctic species were widespread in Hudson Bay, the Atlantic was not. Huntsman (1954) in a discussion of the production of life in Hudson Bay outlined a number of reasons for the apparent low productivity as compared to areas at similar or higher latitude and, while appearing to emphasize the "lack of heat," concluded that as yet the data are

too few "on which to base proper judgement". Other factors which might influence the "apparent low production" (Grainger, 1968, p357) include "the effects of long periods of ice cover" (p358) and consequences related thereto. Dunbar (1970) on the other hand apparently considers nutrient levels to be important. He suggested that the oceanographic regime of Hudson Bay would be altered if the "supply were cut off" of Arctic water moving eastward through Fury and Hecla Strait. He considered the significant alteration would occur in the stability which would become less and as a consequence of mixing processes, more of a nutrient of the deeper water would become available to the surface layer, where productivity would increase. Of a number of questions which might be raised, that concerning the influence of Hudson Strait seems the most important. It is visualized above that mixing processes in the strait determine to considerable degree the distributions within Hudson Bay, such that a change in a surface water characteristic, e.g. to a higher salinity, would be reflected in a change in the same direction in the deeper water. Thus, a decrease in the contribution to the system of low-salinity surface water brought about by damming Fury and Hecla Strait would lead to an increase in salinity throughout, perhaps with little change in stability.

It seems that Dunbar assumed that the nutrient level in the deeper water is high, but this may not be. Should the uncoupling and recirculation suggested here

be significant it is possible that the level of nutrient, or of a nutrient, may be generally low throughout the system (it may be that the limited occurrence of Atlantic copepod there may also be a partial result of this uncoupling).

Nevertheless, it would be of interest to demonstrate that Dunbar's secondary consideration is indeed possible, i.e. that productivity would increase in a region of annual ice cover were the level of nutrients increased. As noted, data are not available for Hudson Bay, but the evidence for some other areas with ice cover indicate that a depletion of at least one nutrient occurs by mid-summer (Apollonio, MS undated; McLaren, 1969). If this is so, it seems that an experiment (McLaren, 1969) in which nutrient is applied to a relatively isolated (uncoupled) body of water in order to avoid the development there of a period of nutrient depletion might provide a useful result. Omarolluk Sound is suggested as a suitable site should it be determined that the nutrient depletion occurs there.

A further consideration here relative to Hudson
Bay is that of the dissolved oxygen. It seems that it
is characteristic of the water in Hudson Strait to be near
saturation levels, from which it follows that the moderate
depletion of oxygen observed in the deeper water of Hudson
Bay occurs entirely within the bay. If the consumption
there were known, it would be possible to estimate an age
for the oxygen depleted deep water. This is apparently

not known, but is likely about the 0.21 ml/l/year suggested for the North Atlantic by Riley (1951). Assuming about 80 percent is retained after leaving the surface and subsequent oxidation of surface nutrients (Redfield et al., 1963), then the age is in the range 5 to 7 years.

2.4 Longer term change within the system

Bailey and Hachey (1951) recognized that the general level of salinity observed by Hachey (1931b) in 1930 in Hudson Bay was low. They compared the data to observations made in 1948 and suggested that the observed difference was due to an increased Atlantic influence. It was suggested (Barber, 1967, p55) that it is not possible through study of other data to reject their hypothesis. In the latter work and in Bailey and Hachey (1951) it appears to have been assumed that the tabulated depth and salinity values (Hachey, 1931b, p96) are without more than the usual error.

It is not difficult to accept the assumption as regards the salinity data for, although they are generally low, it is possible to see in the observed distributions similarities with more recent data. For example, the distribution in Hudson Bay at 50m (Hachey, 1931b, his Figure 4) can be interpreted so that the pattern is compatible with a 50m distribution shown here (Figure 4) based on much more data. However, there are two salinity values in Hudson Strait, each at 200m depth at stations 57 and 48 (Hachey, 1931b, p96), which on the basis of all the data are 0.5 % oo lower than is to be expected at the

depth. At nearby station 58 the salinity values within the surface layer, 32.5 °/oo, are appropriately high. Thus, the data in Hudson Strait and in Hudson Bay suggest that if a persistent error exists it is not in the salinity and could be in the depth.

Depth data presented in the report of the 1930 observations (Hachey, 1931b) are of two kinds. One is the depth to the bottom at each station (p95) which when plotted appear to fit current information. The other data are the tabulated value of the serial samples (p96). feature of the tabulation is that each is at a "nominal" depth. This could, of course, be easily achieved, particularly in Hudson Bay where water movements are generally small. In Hudson Strait, however, strong water movements are known to exist so that extraordinary consideration would have had to be given in order to achieve a desired sampling depth, particularly in the circumstance that a cast comprised a lowering of only one reversing bottle (Hachey, 1931b, p95). It is suggested that this did not occur and consequently the depth data are liable to more than the usual error. Consideration of the validity of the hypothesis of Bailey and Hachey (1951) must therefore include an evaluation of the precision of the depth data.

It does appear therefore that there is little good physical evidence on which to base hypothesis for recent changes of marine climate, i.e. beginning about 1930. It seems that such hypotheses were of much interest in about the late 1940's and early 1950's, e.g. Dunbar's

(1951) extensive work, so that even limited data in rather complicated fiords were used to suggest the existence of change (Nutt and Coachman, 1952); more recently such data were said to "document" a change in the water of the fiord (Coachman, 1969, p215).

## 3. Discussion

The region is ice and snow covered for much of the year so the extent that the relative proportion of land and water areas are altered by the project would not have a major influence during the winter because the ice and snow largely serve to uncouple the underlying ground or water and the atmosphere. This uncoupling and the fact that the winter climate is due to air mass movement of global scale suggests that if the project is to have an influence on the water it would occur at other seasons. We know that the world ocean is largely buffered against change and that the James Bay Project would not have significant impact were the project not located in a system so uncoupled from the world ocean. The limitation to the coupling is due largely to processes within the system including tidal mixing in Hudson Strait. In James Bay the coupling to Hudson Bay is also limited, but through a weak convective circulation. It seems that the project could influence this convective circulation.

It is characteristic that the freshwater from runoff entering James Bay would lead to the observed layer

of low salinity water in the surface and to a distinctive pattern of circulation. The freshwater moves seaward in this surface layer, entraining salt from below and eventually leaves the bay as salt water. The outward movement of salt is balanced, during some time interval, through a sub-surface inflow of water of relatively high salinity so that:

$$T_i S_i = T_0 S_0;$$

As the freshwater moves seaward at a rate equal to the supply then:

$$T_0 = T_i + R.$$

If these be applied in the section across the entrance where  $S_i$  and  $S_o$  may be 29 % of and 22 % of respectively, then the inflow is three times the runoff; thus the net transport is a significant part of the total circulation.

Present evidence indicates that the strength of the outflow is strongly time-dependent, presumably in association with seasonal changes in the volume of runoff. The outflow would probably be a maximum at some time after the peak in runoff in early summer. Earlier, in late winter, a minimum in the runoff is indicated about which time a minimum in the outflow would occur. As well, the runoff would occur to a region covered with ice and may not entrain the deeper seawater to the same extent as when not ice covered. The energy for mixing is believed input to the system from tides, wind and surface processes of heating and cooling. An ice cover would reduce the influence of wind considerably as well as of the other

surface processes, so that mixing would be less under an ice cover and the volume of the inflow would be even closer to that of the runoff than it is in the late summer.

The significance of the foregoing is that the runoff does not provide a strong coupling to the water of Hudson Bay, i.e. the estuarine circulation is relatively weak, so that changes in the pattern of runoff may not change the coupling significantly. Should changes in the extent of the ice cover occur, say to decreasing cover, then an increase in the coupling could be expected and a greater forced circulation would result. However, the return flow comprises a relatively cold water, -1.4° to -1.0°C, so that an increased coupling during the summer (period of heating) would tend to a colder surface layer. Conversely, uncoupling would tend toward a warmer surface layer.

For example, if James Bay were completely uncoupled from Hudson Bay through a physical barrier across the entrance, except that an outflow occurred to balance the inflow as in a lake, then a direct influence of Hudson Bay water would not occur. Eventually the bay would likely contain freshwater only, which would be warmer than now, probably close of 4°C in the deeper water, and would likely undergo a wider range of temperature in the surface as a greater storage of seasonal heat would take place. An ice cover would still form and would likely be similar to that which occurs now. The sum of the radiative and turbulent flux terms in the annual heat budget would be

close to zero, as it is now believed to be (assuming the average temperature of the inflow is the same as the outflow), but the average temperature would be higher throughout.

If the coupling were somehow steadily increased the water in the bay would tend toward that of Hudson Bay, in particular toward that comprising the inflow, i.e. toward a water of salinity 29 0/oo and temperature about -1.4° to -1.0°C, but eventually without an ice cover. Of course this situation would not likely be achieved, for processes within James Bay would, at some stage, begin to influence distributions within Hudson Bay, i.e. the radiative and turbulent flux terms of the heat budget would show a large deficit which could only be balanced by advection of water from Hudson Bay. It seems, therefore, than an increased coupling to Hudson Bay would tend to a decrease of water temperature, while a decrease in the coupling would have the opposite result. An increase in the runoff to James Bay, for example by diversion from Grande Rivière de la Baleine, would increase the coupling, while the smoothing of the runoff so that a greater portion occurs under an ice cover would tend to reduce the coupling. It does not seem possible now to provide a quantitative estimate of the results which may be anticipated, except that the influence of the smoothed runoff would be greater. This would decrease the coupling in summer so that the water in summer would tend to be warmer, or at least would tend to store more heat. Most of this heat would be given

off during the period of net heat loss prior to ice formation. During the latter part of this period rather massive losses occur, perhaps as large as 500 - 600 g cal cm<sup>-2</sup> day<sup>-1</sup>, so that in order for the time of formation of first ice to be influenced, the stored heat would have to be appropriately large. An estimate of the difference has not been possible. Neither has it been possible to determine the changes which might be brought about in the average pattern of ice cover.

In this simple analysis it is assumed that changes in the surface condition in Hudson Bay, mainly that of the ice and snow cover, will not occur. There is, however, the possibility, described in earlier section (2.2.1), that the present distribution of ice in early summer there may reflect the inferred strong time-dependence in the outflow at the surface from James Bay. Some conclusions might be possible were we to achieve an understanding of the factors which determine the present average situation.

winds, which are generally from the northwest and north early in the breakup, have generally become light and variable by mid-July and continue so into August. Thus, after creating an accumulation of ice toward the south of Hudson Bay by about the end of July, the wind has little or no influence so that no subsequent ice movement occurs due to wind. Furthermore, there is evidence that during this period of stagnation the frequency of both cloud cover and fog is increased so that less insolation reaches the ice surface and melting is retarded;

complete clearance of the ice may not occur till about the end of August on the average.

In conclusion, it seems likely that advection is not a significant factor in the total heat balance of the system and although a real understanding of the relative influence of inflows and outflows and of processes is not achieved, the system within Hudson Strait appears to be closed rather than open. This tentative result emphasizes the importance of the radiative and turbulent flux terms, i.e. the mainly climatic factors, which in turn reflect the influence of a process global in scale, at least in winter through to early summer, which determines the nature of that important variable the surface condition. at this time either ice or snow. Within the system the water of James Bay is not strongly coupled to that of Hudson Bay but as similar climatic conditions prevail so the surface condition is one of ice or snow for a significant period. It does not seem likely that the project could measureably influence the coupling of the system to the world ocean. It could, however, influence the coupling of James Bay to Hudson Bay to the extent that a greater storage of heat may occur within James Bay during summer after breakup and influence the distribution of ice in the approach to the bay during breakup. Thus the impact of the James Bay Project on the system will be relatively small and it may be necessary, in order to assess the impact, to acquire a particularly extensive body of data; a significant portion of which should be obtained during the period of ice cover.

At other times it may be possible to utilize remote sensing techniques to considerable advantage; although the fact that the peak of the present runoff occurs when there is still considerable ice cover emphasizes the difficulty of the field problem. Of particular value now would be information on the water structure over a period of a year at one postion within the bay, although data in sections seaward of Rupert Bay, Fort-George and Poste de la Baleine and eventually at Inoucdjouac would be preferred. Part of the field programme could be based on relatively simple oceanographic instruments with a frequency of observations about once a week and perhaps sustained by residents there, at least during the period of ice cover, as occurred at Tuktoyaktuk (Kelly, 1967, p8).

## 4. Acknowledgements

The Satellite Data Laboratory of the Meteorological Branch (now the Atmospheric Environment Service) provided reproductions of data received from the satellite ESSA 2 and Ice Forecasting Central of the same Branch provided advice and forecasts relating to probable ice movement. The 1967 observations in CCGS "Labrador" were carried out by ship's personnel with equipment from the Bedford Institute. Mrs. P. Milner prepared the material for publication and Mrs. K. Ramlo and Mrs. P. McAinsh typed the report.

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shown in Figure 21 and the date of sampling

## 7. Appendix

Figures 21, 22, 23 and 24 based mainly on the "Calanus"

1959 data (Grainger, 1960), but including some of the

"Theta" data (Anon., 1964a).



Figure 22. Distribution at various depths of salinity (o/oo) and temperature (°C) as interpreted from the "Calanus" data of 1959. The positions of the stations are shown in Figure 21 and the date of sampling in Table 2 where it may be seen that the interval between occupations was relatively long.

(a) Surface. (b) 10m. (c) 25m.
(d) 30m. (e) 50m. (f) Deepest.



Figure 22(a)

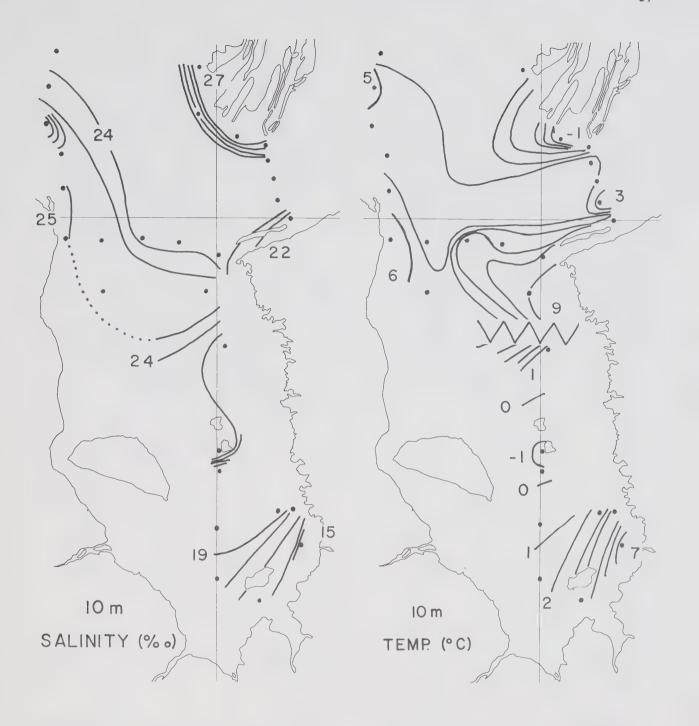


Figure 22(b)

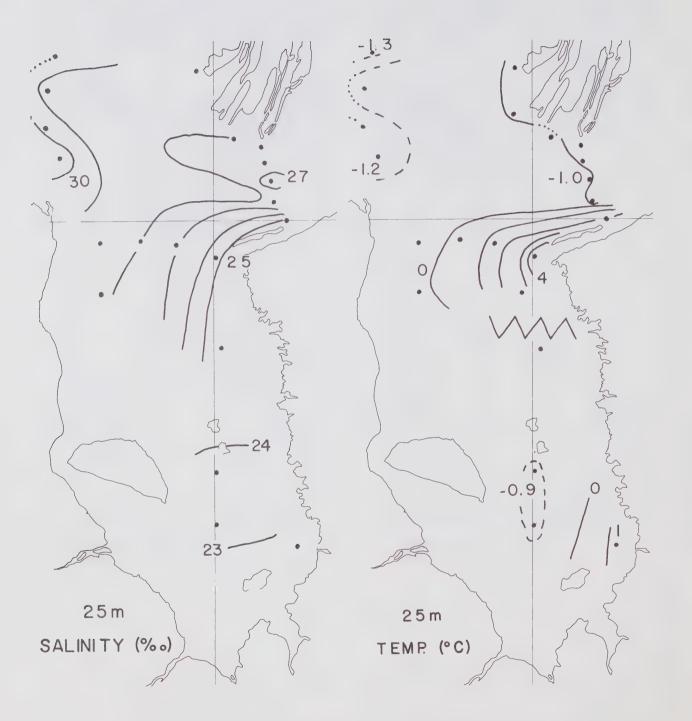


Figure 22(c)

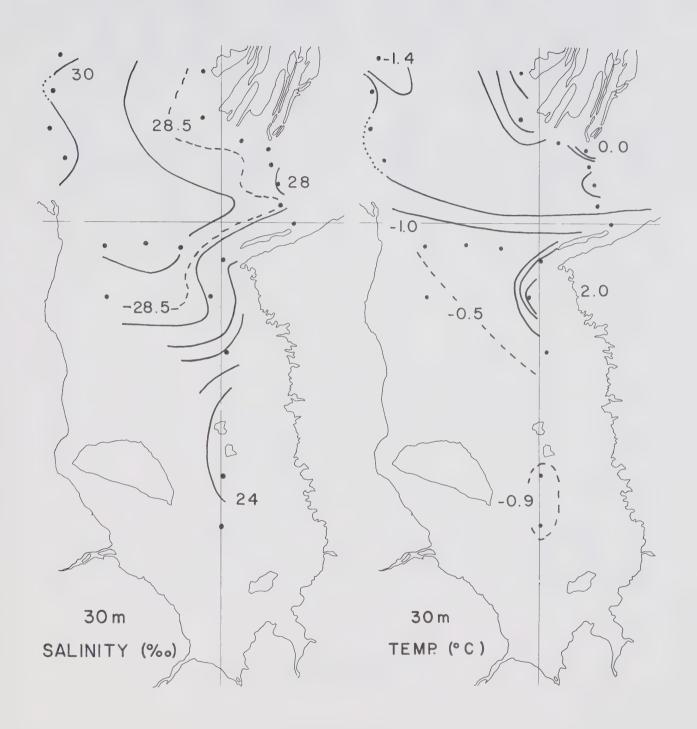


Figure 22(d)

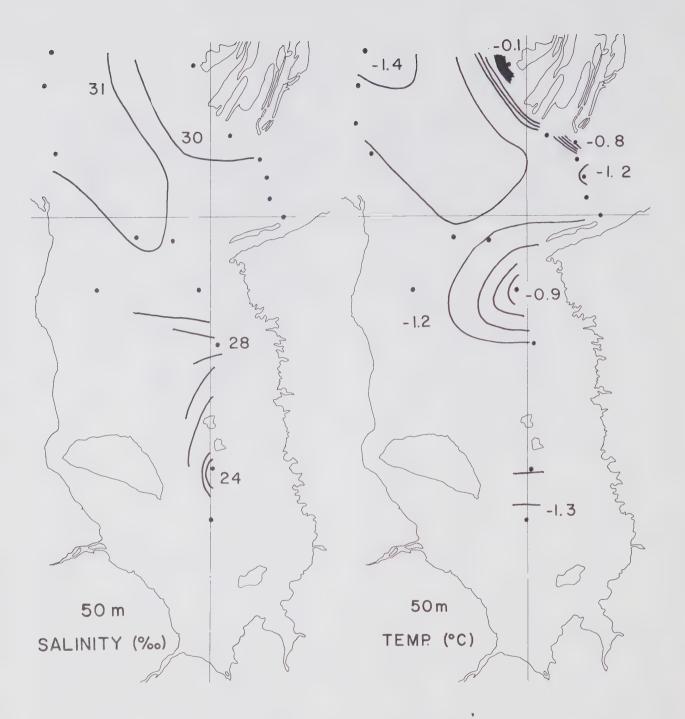


Figure 22(e)

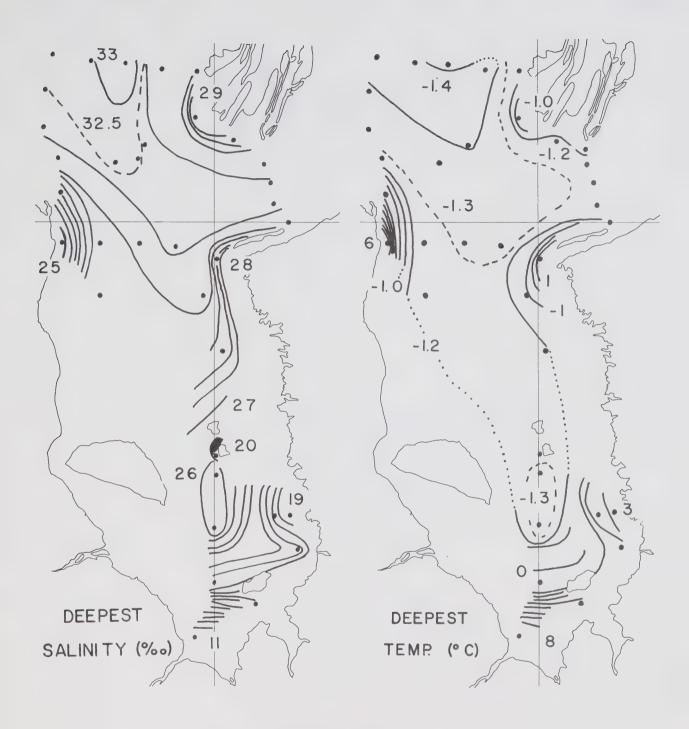


Figure 22(f)

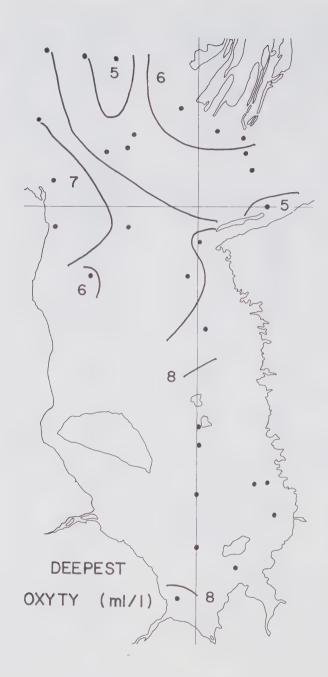


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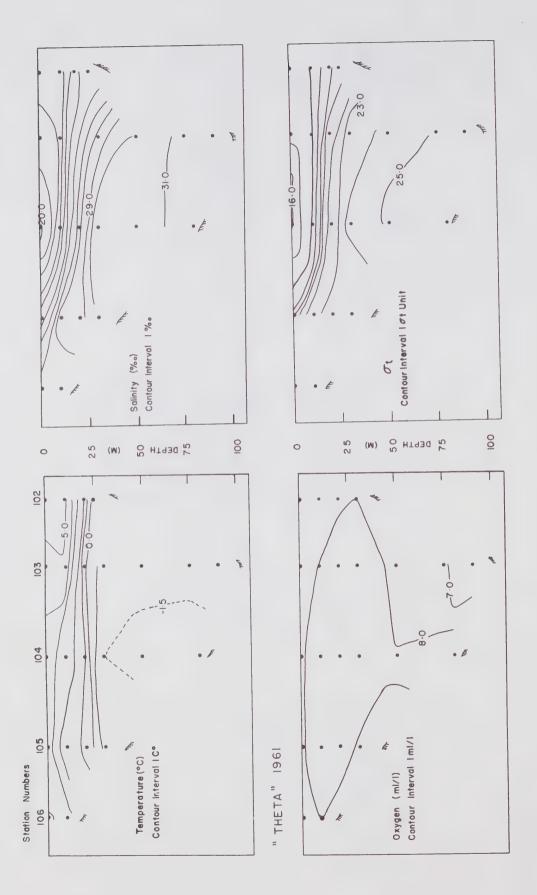


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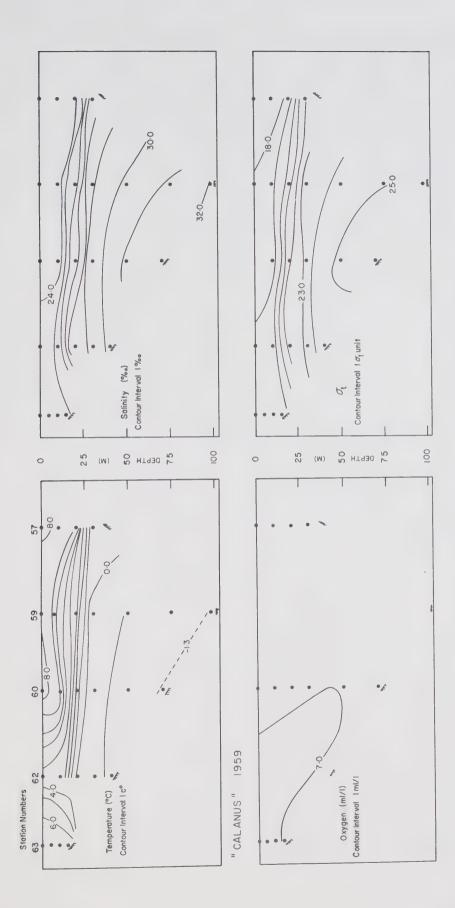


Figure 24(b)

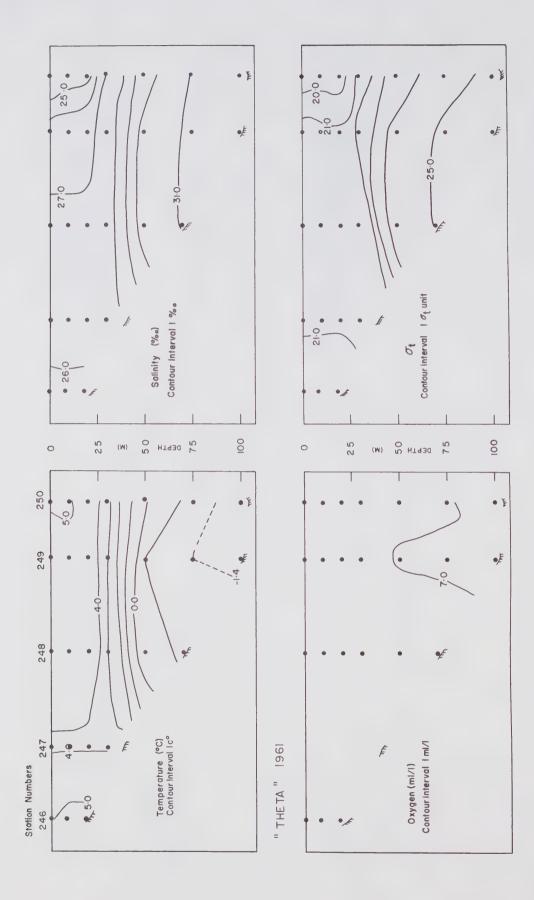
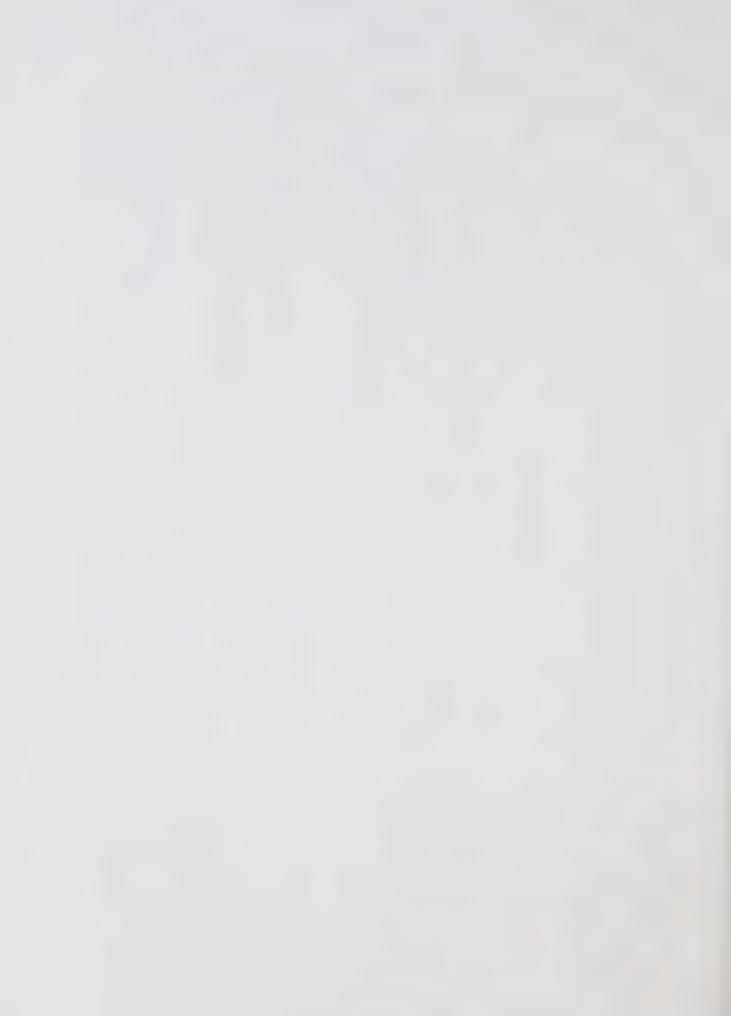


Figure 24(c)

The tides in James Bay

by

Gabriel Godin



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## 1. The data available

James Bay (frontispiece) because of its subarctic climate and its inhospitable shores has not attracted a large permanent population around its rim. Scientific information about its geography, bathymetry and oceanographic characteristics has been acquired in a succession of hesitant steps. Its extensive flats, its hundreds of islands and reefs make most of it virtually inaccessible to any hydrographic vessel. This severe limitation on ship traffic inhibits any motivation for acquiring more extensive soundings with the help of launches, an expensive and time-consuming project. As a consequence the bathymetric information about the western and eastern sides of James Bay consists mostly of blanks. The extent of soundings in the central portion of the bay however, suffice to delineate a relatively deep channel of depths exceeding 25 m, extending from its mouth down to the west side of Charlton Island, while the remainder of the bay remains very shallow.

Temporary tide gauges have been installed at odd stations at irregular time intervals and records of short duration have been accumulated. The majority of stations were established on islands such as Bear, Strutton, Charlton and Stag Islands, as well as inside sand and sediment choked rivers such as La Grande-Rivière, Eastmain, Moose and Albany. None of these observations can be considered of good quality with the exception of those carried out in the Moose River (Langford, 1963). As a matter of fact, the careful observations in the Moose River, which monitored currents and water levels from well outside the

estuary up to the Government Wharf in Moosonee at the foot of rapids, help assess the representativeness of the observations gathered at the other stations. For instance the tide takes 2 1/2 hours to travel the 13 nautical miles from Sand Head to Moosonee; therefore, one must expect a lag of the same order of magnitude between the tide recorded at other trading posts located at the head of rivers and the tide present in their respective estuaries, which makes such observations quite unrepresentative of what is happening in James Bay proper. The observations gathered at the island stations should represent the actual tide in James Bay more adequately. Even at that, there exists a peculiar phase difference between the M2 tide observed at Charlton Island and Strutton Island which might be accounted for by the very sheltered position of Charlton Depot.

## 2. Cotidal charts using the gauge observations and a one dimensional model

With the help of these rather sparse observations and the use of a simplified one dimensional model of James Bay, it is still possible to gain a fair idea of the progression of the tide inside the bay and of its change in amplitude.

The tidal wave which progresses along the southern rim of Hudson Bay is strongly refracted around Cape Henrietta Maria and enters James Bay as a damped progressive wave. Its advance appears to be considerably retarded in the vicinity of Akimiski Island and it eventually reaches Hannah Bay and Rupert Bay seven hours after it has rounded Cape Henrietta Maria. Its amplitude is larger on the western and southern portions of the bay due to the presence of a degenerate node of the semidiurnal tide on the

eastern side of the bay. The mean amplitude of the tide varies from over 90 cm in the western and southern portions to less than 40 cm in the eastern portions near Akimiski Island. Table 1 lists the amplitude and phase of the major constituents of the tide observed at various gauge stations established in James Bay; we include Poste de la Baleine and Winisk to relate these observations to Hudson Bay proper. M2 is the major lunar semidiurnal constituent; it represents the mean tide. S2, the solar semidiurnal constituent, adds or subtracts its contribution to that of M2, creating spring and neap tides. N2 is of lunar origin and reflects the variable distance of the moon from the earth. We have kept it in parentheses because most likely it has not been properly separated from the other semidiurnal constituents. K1 and O1 are diurnal constituents and are related to the declination of the orbits of the moon and of the earth. Their amplitude and phase are irregular and Ol appears to be twice as small as K1 throughout the bay. Normally O1 should be about two thirds of Kl. If there were no friction in James Bay, both of these constituents should have a node just at the mouth of the bay, the node of K1 being positioned a little further inside the bay. In practice the node is degenerate and somehow O<sub>1</sub> enters the bay considerably weakened compared to K<sub>1</sub>. All in all, the diurnal constituents are very much smaller than the semidiurnal constituents thus indicating that the tide throughout James Bay is truly semidiurnal at all times.

Amplitude and phase of various tidal constituents observed at thirteen stations around James Bay. Table 1

		M2		S2		(N <sub>2</sub> .)		K1		01
Station	amp	phase	amp	phase	amb	phase	amp	phase	amp	phase
Poste de la Baleine	63	232	15	302	6	184	4	25	0	334
Pte. Louis XIV	64	222	20	281	20	191	9	27	П	311
Fort George	99	238	15	329	10	211	∞	116	Ŋ	51
Eastmain	34	49	9	134	ſΩ	349	10	145	7	110
Strutton	52	14	IJ	124	ıΩ	355	10	141	9	06
Charlton	63	48	16	128	12	n	12	128	m	78
Stag	16	8 0	20	164	16	38	12	145	σ	110
Sand Head	94	44	22	126	18	357	16	128	7	63
Ship Sands	71	71	16	155	12	31	15	152	m	64
Moosonee	62	111	14	200	11	89	12	171	2	09
Fort Albany	<b>∞</b>	40	22	120	18	0	15	139	4	06
Bear Island	97	219	19	302	19	182	∞	62	m	349
Winisk	109	က	29	158	19	45	ſΩ	322	m	280

We may use the data presented in Table 1 to draw cotidal charts of the constituents over James Bay. In these charts the lines of equal phase may be considered as delineating the front of the tidal wave as it progresses up James Bay. A phase difference of 60° between two such lines represents about two hours for a semidiurnal tide and four hours for a diurnal tide. The lines of equal amplitudes denoted by dotted lines define areas within which the amplitude exceeds or is less than the value indicated on the rim.

Such isopleths may be drawn in an almost infinite number of ways because of the scarcity and problematic value of the actual observations. It seems wise at this stage to solve the equations of hydrodynamics over a simplified one dimensional model of James Bay. In this way we may predict the mean vertical tide and tidal currents at various sections of the bay assuming that the tidal wave is perfectly reflected at the head. The results of such calculations may be combined with the coastal observations to draw the most plausible system of cotidal and coamplitude lines, besides yielding values of the tidal currents which have not been observed in the bay except around the Moose River. The equations of hydrodynamics in one dimension are

$$\frac{\partial}{\partial x}(Au) + B\frac{\partial Z}{\partial t} = 0 \tag{1}$$

$$\frac{1}{g} \frac{\partial \mathbf{u}}{\partial t} + \frac{|\mathbf{u}|}{C^{2}H} \mathbf{u} + \frac{\partial \mathbf{z}}{\partial \mathbf{x}} = 0$$
 (2)

where x is the distance along the bay moving northward,

A is the area of a vertical section of the bay drawn across its width,

B is the width of such a section,

H is the mean depth of such a section,

Z is the amplitude of the tide and

u is the tidal current,

while C stands for the Chézy coefficient which measures the intensity of friction. Since the friction all over the bay must be rather high we take the value of C as  $C = 45 \text{ m}^{\frac{1}{2}}/\text{sec}$ .

Because of the friction term  $(|u|/C^2H)u$ , the phase difference between Z and u varies from section to section and for a sinusoidal oscillation of frequency  $\sigma$  (a given tidal constituent), we may write a solution to (1) and (2) in the form

$$Z=Z_1(x) \cos \sigma t + Z_2(x) \sin \sigma t$$
 (3)

$$u=u_1(x) \sin \sigma t - u_2(x) \cos \sigma t,$$
 (4)

so that we may write for the x dependence of Z and u:

$$\frac{\partial}{\partial x}(Au_1) = B\sigma Z_1 \tag{5}$$

$$\frac{\partial}{\partial x}(Au_2) = B\sigma Z_2 \tag{6}$$

$$\frac{\partial Z_1}{\partial x} = -\left[\frac{\sigma}{g} u_1 - \frac{|u|}{c^{2H}} u_2\right] \tag{7}$$

$$\frac{\partial Z_2}{\partial x} = -\left[\frac{|u|}{C^{2H}} u_1 + \frac{\sigma}{g} u_2\right] \tag{8}$$

u<sub>1</sub>, u<sub>2</sub>, Z<sub>1</sub>, Z<sub>2</sub> are inextricably linked because of the friction.

(5) to (8) become useful to us in difference form:

$$(Au_1)_{j+2} = (Au_1)_{j} + \Delta x \sigma B_{j+1}(Z_1)_{j+1}$$
 (9)

$$(Au_2)_{j+2} = (Au_2)_j + \Delta x \sigma B_{j+1}(Z_2)_{j+1}$$
 (10)

$$(z_1)_{j+1} = (z_1)_{j-1} - \Delta x \left[ (\sigma/g) (u_1)_{j} - (|u|_{j}/c^2 H_{j}) (u_2)_{j} \right]$$
 (11)

$$(z_2)_{j+1} = (z_2)_{j-1} - \Delta x \left[ (|u|_j/c^2H_j)(u_1)_j + (\sigma/g)(u_2)_j \right]$$
 (12)

By dividing the bay into segments of length  $\Delta x$ , (9) to (12) allow us to evaluate Z and u at section  $2\Delta x$  apart from the values set at the boundaries. We have in fact, subdivided James Bay into 23 segments of length  $\Delta x$ =10 nautical miles extending from 51°10'N to 55°00'N. The integration has to be performed for each constituent. We restrict ourselves to M<sub>2</sub> and K<sub>1</sub> only, since calculating S<sub>2</sub>, N<sub>2</sub> and O<sub>1</sub> would simply be repetitious. We use for M<sub>2</sub> ( $\sigma$ =28.98°/hour) as boundary conditions

$$u=0$$
 at  $j=0$  (51°10'N) (13)

 $Z=91(\cos \sigma t -40^{\circ})$  cm at j=1 (51°20'N),

which is equivalent to taking

$$Z_1=70$$
 cm and  $Z_2=59$  cm at  $j=1$  (14)

For  $K_1$  ( $\sigma=15.04^{\circ}/hour$ ) we take

$$u=0$$
 at  $j=0$  (15)

 $Z=15\cos(\sigma t-125^{\circ})$  cm at j=1

or 
$$Z_1=-9$$
 cm  $Z_2=12$  cm at  $j=1$  (16)

Table 2 lists values of B, A and H which have been derived from our schematization of the bay into the 23 subsections; the actual profiles are presented in the appendix. Table 3 contains the results of the integration of (9) to (12) using (13) to (16). These values cannot describe any two dimensional features of the tidal motion and represent averages over a whole section. They are plotted at the points indicated by x and 0 in the cotidal charts and they supply considerable assistance in the drawing of plausible and consistent coamplitudes and cotidal lines. We draw similar charts for S<sub>2</sub>, N<sub>2</sub> and O<sub>1</sub> using the features that had been delineated for M<sub>2</sub> and K<sub>1</sub> using the one dimensional model. All these charts are shown in Figures 1 to 5.

A side product of the numerical calculations is that they yield values of the mean current at the various sections; these have been noted on the charts of  $M_2$  and  $K_1$ . Actual

Table 2 A one dimensional schematization of James Bay indicating the section j, (equation 13) and the width B, the area A and the depth H (equations 1 and 2).

Section j	Width B (km)	Vertical Area A (xl0 <sup>6</sup> m <sup>2</sup> )	Depth H (m)
0			
1 :	38		
2		.68	11
3	109		
4		2.07	16
5	118		
6		2.91	19
7	203		
8		3.68	18
9	126		
10		3.09	24
11	125		
12		3.50	44
13	195	`	
14		3.42	33
15	200		
16		5.80	29
17	175		
18		7.21	37
19	194		
20		9.31	51
21	170		
22		11.44	69
23	245		

Values of M2 and K1 deduced from the one dimensional model of James Bay. Table 3

6	amp phase cm deg	15 125	15 124	14 123	13 22	_			7 112	4 100	3 87	2 62	2 29	2 359
., K1	phase	ر ت	n n	34	ന	33	33	32	31	29	7.0	را ر د بر	2 4	1
	amp cm/sec 0	c	N	m	ហ	9	∞	Φ	0	9	L	) 4	, ,	
7	phase	40	31	22	ហ	223	7 U	000	258	238	226	216	207	199
	amp	91	78	64	49	000	) \ \	n F	19	75	78	92	71	65
M2	phase	l r	310	304	299	291	285	274	252	224		O CO	0 7	, H
	amp cm/sec		76	31	37	45	61	56	υΩ	, c	) (C	) c	o 00	0
·n			N W	4 r	9 7	- 00 0	010	12	13	115	17	0 6 6	210	23

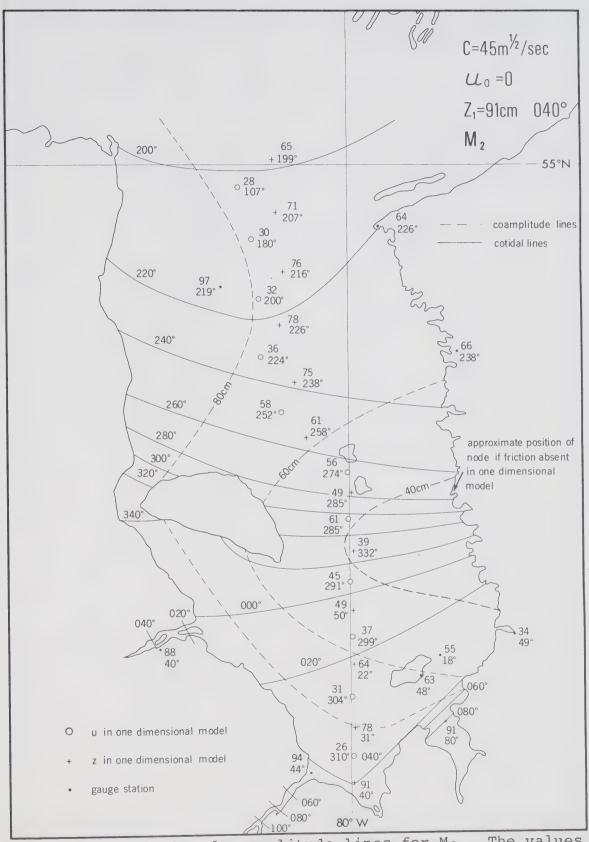


Figure 1 Cotidal and coamplitude lines for M2. The values observed are squared. x and 0 indicate the values of Z and n deduced in the one dimensional model.

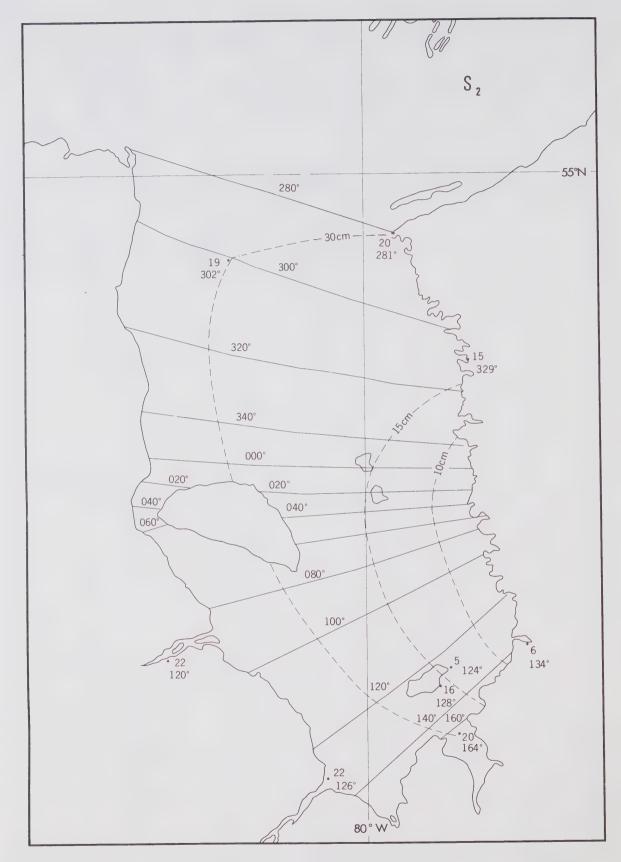


Figure 2 Cotidal and coamplitude lines for  $S_2$ .

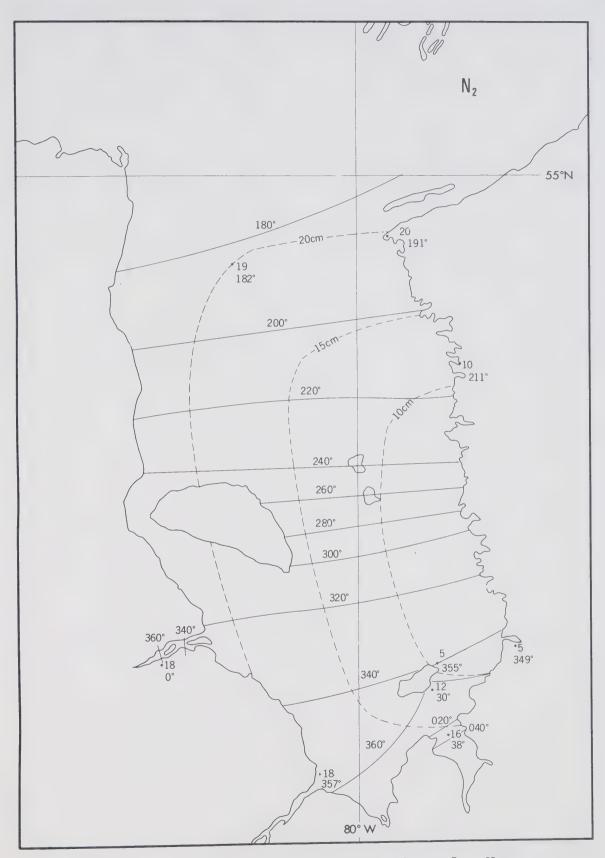


Figure 3 Cotidal and coamplitude lines for  $N_2$ .

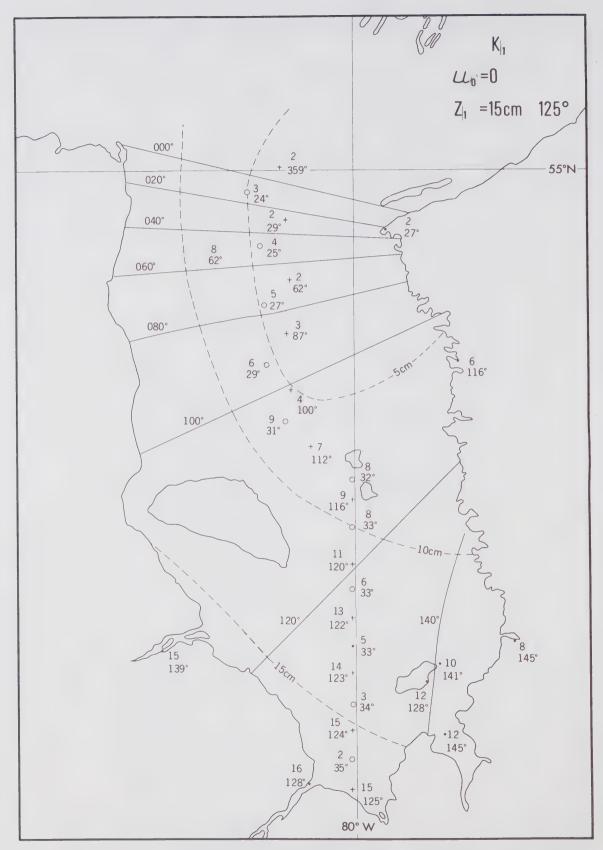


Figure 4 Cotidal and coamplitude lines for  $K_1$ .

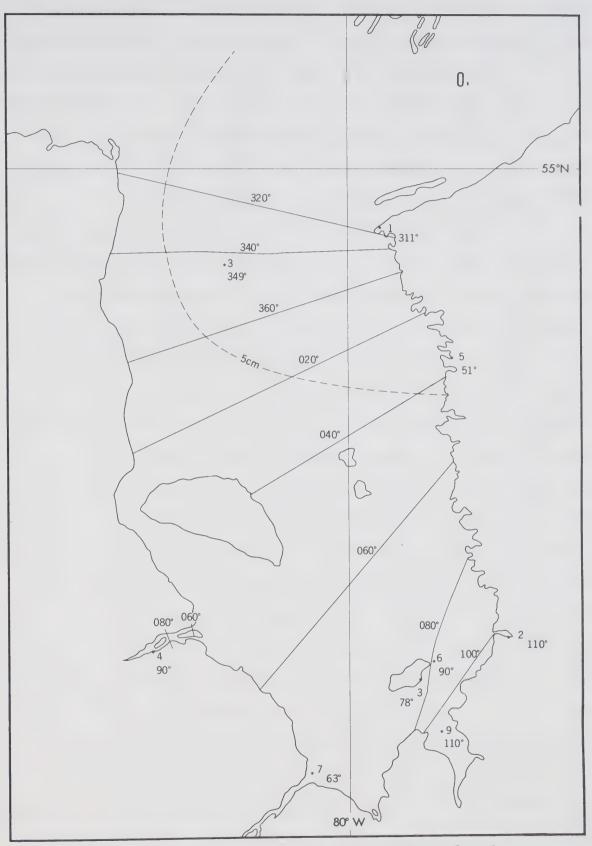
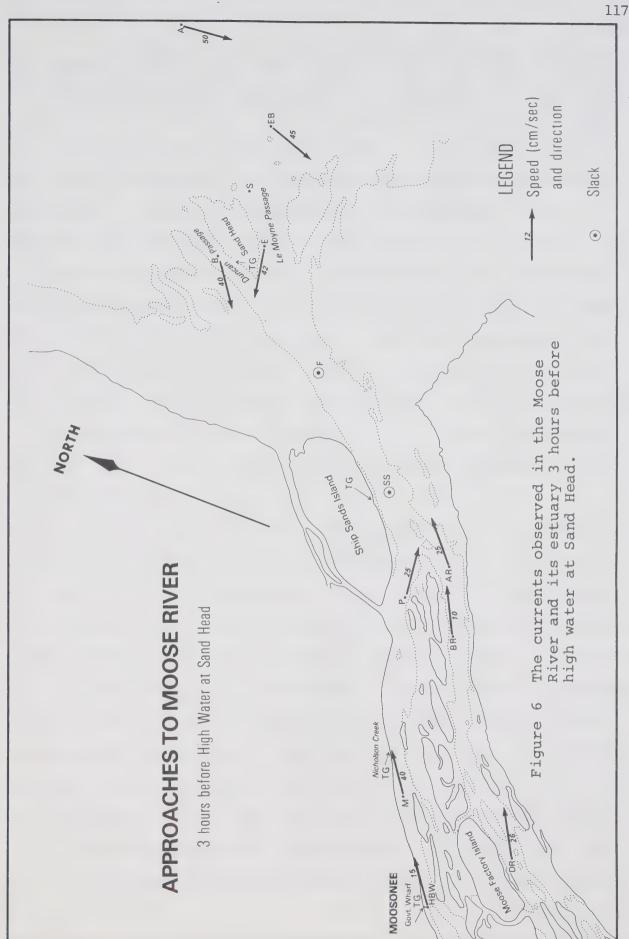


Figure 5 Cotidal and coamplitude lines for  $O_1$ .

observations of the currents in James Bay are not available but some currents were measured in the estuary of the Moose River.

Our solution to (1) and (2) depends on our assumption of perfect reflection at the head of the bay; in practice the reflection is rather diffuse and we cannot expect that the current observed in that area should lead the vertical tide by 90° of phase as our calculations suggest.

In the vicinity of Sand Head in the estuary of the Moose River, the vertical tide was observed on the tidal flat while the current was monitored in two neighbouring channels. At station B (Figure 6) the current was exactly 90° of phase ahead of the vertical tide as predicted by the model, but at Station E the current was more like 45° of phase ahead; the current at E is more representative of the current in the bay proper because the channel there is deeper and the flow is much less influenced by friction as it is at Station B. The observed currents are also larger than those predicted by the model, but it must be kept in mind that the model predicts the average current over a whole section while the current observed is present at a given point which happens to be more shallow than the remainder of the section. Moving away from the boundary the currents predicted from the model should be more representative and we notice that they reach their maximum of 61 cm/sec off Akimiski Island for M2 and 9 cm/sec for K1. Taking the contribution of S2 and N2 into account, the currents caused by the semidiurnal tide can therefore reach a magnitude of 80 cm/sec at spring tide when the moon is in perigee. This gives an



indication that the tidal currents over the body of James Bay in the vicinity of Akimiski Island can reach quite a respectable intensity.

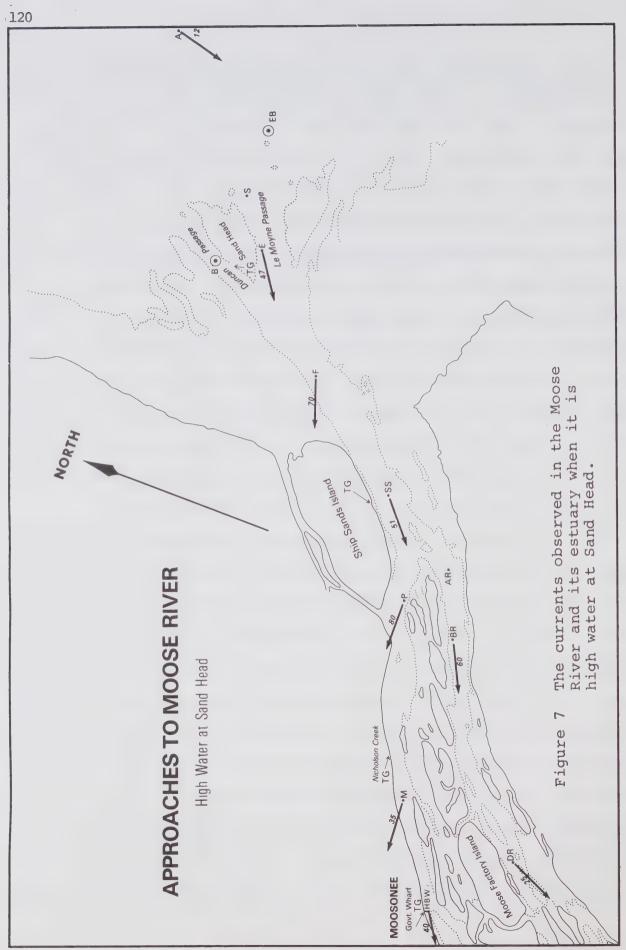
3. A detailed view of the flow caused by the tide in Moose River

Up to now we have struggled with a handful of data and
a coarse model in order to obtain some idea of the tidal motion
in James Bay. It comes as a relief to this bleak situation to
have in our possession the results of the detailed survey carried
out in the Moose River and its estuary by the Canadian Hydrographic Service (Langford, 1963). In this fashion we can
actually witness the flow of water in and out of a tributary of
James Bay which most likely is typical of what is happening
inside the other rivers emptying into James Bay.

The bed of the Moose River is quite undefined, shallow and cursed with numerous islands and drying flats of mud, sand and boulders. Sand flats nearly block its estuary and only the pressure of its impounded water manages to keep a gully open towards the sea (the Le Moyne Passage). The shores of the Moose River consist in many places of soft materials which are undermined by the ice and rushing waters during the spring freshet; later in the season, these mined areas collapse. Rapids effectively block the river upstream. The Moose River appears as rather typical of the other rivers around the bay and a study of the intricacies of the water movements inside it during a tidal cycle will help give an idea of the actual complexity of the patterns of flow in the vicinity of any of the other tributaries.

In Figures 6 to 9, the solid line delineates the shores and the islands and the dotted line the flats that dry at lower low water. The tide gauge stations indicated by "T.G." were set up at four locations. The first station at Sand Head measures with little distortion the tide that comes in from the main body of the bay. The distance between the first and last gauge amounts to approximately 13 nautical miles. Current meter stations were established at points indicated by the origin of the arrows; they are labelled by letters. We have shown 11 such stations. A circled point at one station indicates that at the moment of observation, no detectable current was noticed (slack water, turn). The arrow shows the orientation and the velocity in cm/sec is indicated. The origin of time is chosen as the moment when it is high water at Sand Head, the station located furthest out in the estuary.

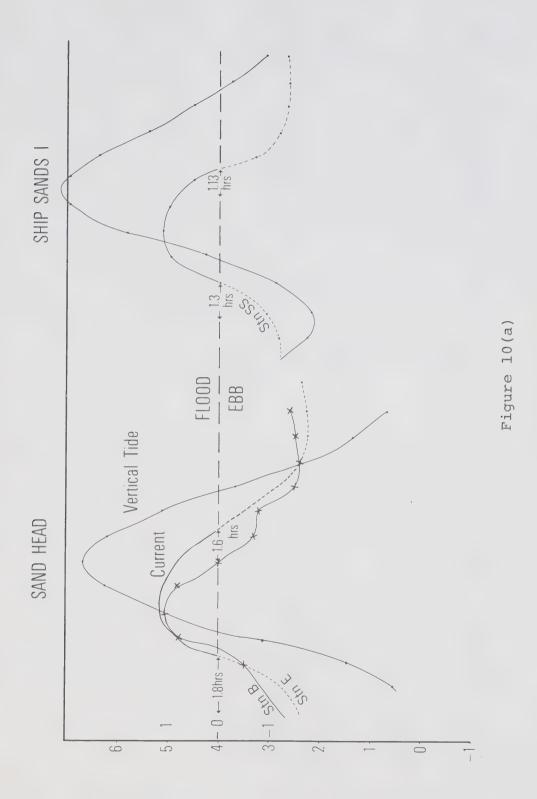
In Figure 6, three hours before high water at Sand Head, which corresponds to nearly 1/4 period of the semidiurnal tide (and therefore approximately to mean water), flood has been established in the estuary while the water is still ebbing in the upper regions of the river; as a consequence we find a region of no motion in the vicinity of Ship Sands Island. This front moves gradually upstream as the tide rises and by the time it is high water at Sand Head (Figure 7) the river is in flood as far as Moosonee. We may notice that the directions of flow in individual channels depends very sensitively on the bottom configuration and that it may change rapidly at times. In the estuary, slack water is already reached in Duncan Passage and in the shallower portions

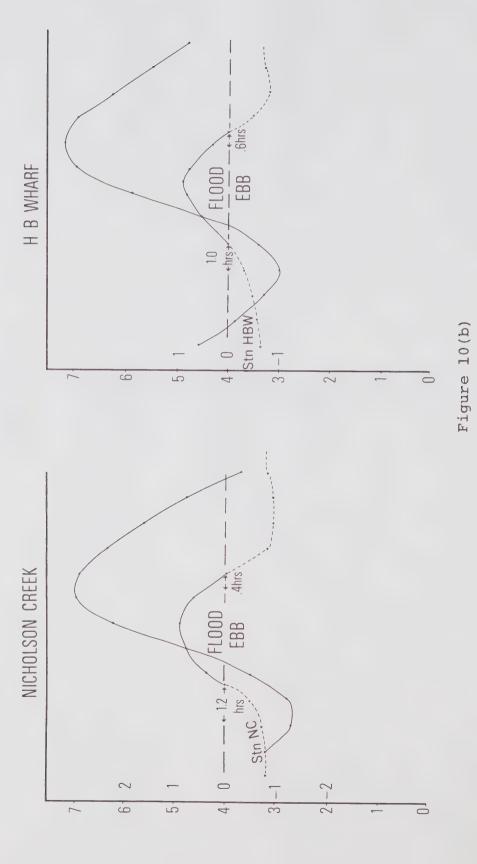


of the estuary while it is still flood over the deeper portions. Slack water occurs earlier in shallow areas because of the friction and it will take a further 1 1/2 hours before the current turns in the deeper portions of the estuary. currents turn clockwise there, while in Le Moyne Passage it slackens suddenly and turns to flood 1 1/2 hours after high This zone of slack water progresses upstream as the water level starts lowering, reaching the northern portions of the river first and the southern portions later. Three hours after high water (Figure 8) the river is in ebb. The water level is decreasing, high water having occurred at Ship Sands one hour after Sand Head while the water was still in flood. High water occurred at Nicholson Creek two hours after Sand Head and then it was slack water in its vicinity. Finally high water at Moosonee occurred 2 1/2 hours after Sand Head while the ebb started to take place in its vicinity. Six hours after high water at Sand Head (Figure 9) it is low water at Sand Head and the entire river continues ebbing.

Figure 10 shows the relationship between the water level at the four gauge stations and the current in their immediate vicinity. The time interval between high water and slack is noted since it gives some idea of the phase difference between the current and the vertical tide. Because of the net outflow from the river, the actual time difference between flood and high water is actually longer than the time noted on the graphs.

Figure 10 The relationship between the vertical tide and the current observed in the vicinity of the observing station.





## 4. Irregularities in the tidal regime

Two factors may affect the regularity of the tides: the spring freshet and weather disturbances. During April-May ice breaks up in the rivers surrounding James Bay and their discharge increases abruptly. An increase in discharge inhibits the penetration of the tide into the river and a wall of fresh water may extend a fair distance into James Bay where the tide will be reduced in range and where the mixing of salt and fresh water will take place at an intensified rate. This damping of the tide by a strong current is indicated by the solution of the equations of hydrodynamics (1) and (2) if we represent the current by

$$u = u_0 + u_1 \cos (\sigma t - \alpha), \qquad (17)$$

where  $u_0$  represents the steady current due to the discharge of fresh water and the second term represents the oscillatory term contributed by the tide. Assuming that the tidal current is never strong enough to cause any flood current,

$$|u|u = u^2 \simeq u_0^2 + \frac{1}{2}u_1^2 + 2u_0u_1\cos(\sigma t - \alpha)$$
 (18)

for which (1) and (2) have a solution of the form

u,h 
$$\sim \exp i\sigma t \exp \pm \frac{i\sigma}{\sqrt{gD}}$$
  $\sqrt{1 - \frac{2igu_0}{C^2\sigma D}} \times$  (19)

There is a damping term in x of the form

$$\exp - \left[ \frac{\sqrt{2u_0\sigma}}{CD} \sin \left( \frac{1}{2} \arctan \frac{2u_0g}{C2\sigma D} \right) \right] x \tag{20}$$

which would vanish if uo were zero.

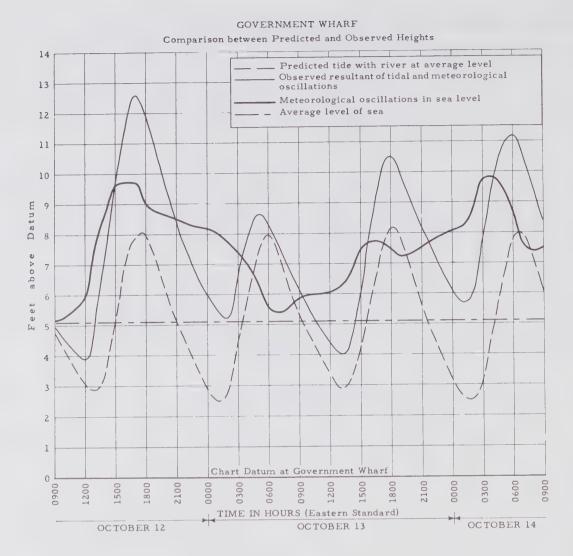
This situation must occur with various degrees of intensity depending on the strength of the freshet. At that time the water will be fresh in the rivers, the overall salinity of James Bay in the vicinity of their estuaries will be significantly reduced and processes of mixing will be intensified but will take place further away from the shores. The stronger currents will carry sediments into the bay at an accelerated rate, particles of a larger size will be dragged along and most probably during this time the major portion of the nutrients fed into James Bay over the course of one year will be carried into it. Observations on the Moose River support this conjecture.

Weather disturbances must also have a marked effect on the motion of water on account of the general shallowness of the area and the presence of modes of resonance in the bay of a period of only a few hours. James Bay might respond dramatically to the passage of some of the cold fronts which are frequent and most severe during early summer and the fall and to the passage of depressions. Rupert Bay, Hannah Bay and the portion of James Bay south of Akimiski Island have minimodes of their own which could be excited by specific patterns of wind and pressure. During the course of a weather surge, the level of water will undergo irregular oscillation and transient currents will be

created. Masses of salt water could invade the rivers. The non-linear interaction of the surge with the tide may enhance or decrease the water level depending on fortuitous combinations of pressure gradients and wind stress. Surges are most likely to occur during the storm seasons, namely between September to December and April to June. A surge, a mild one, was observed on the Moose River while tidal observations were being carried out; it occurred between October 12 and October 14, 1963 and had a height of 120 cm with a period of 30 hours. The accidental detection of such a surge indicates that these phenomena are probably quite frequent in the James Bay area (Figure 11).

# 5. Changes caused by the regulation of the Nottaway, Broadback and Rupert Rivers

Once the NBR project is completed, the tide in Rupert Bay will not be inhibited during the spring freshet as it normally is at that time of the year; this implies that a more regular tidal regime will tend to prevail in and around Rupert Bay throughout the year. The zone of mixing of fresh water and salt water which must wander appreciably inside Rupert Bay at the spring freshet will no longer suffer such fluctuations and will be restricted to an area defined by the intensity of spring and neap tides and the mean value of the regulated outflow from the NBR complex. Consequently, there will be fewer localized short-range fluctuations in density, salt and sediments. The net amount of nutrients brought into Rupert Bay will be reduced since the larger currents necessary for their intensified transport will no longer have any occasion to prevail: the migratory birds



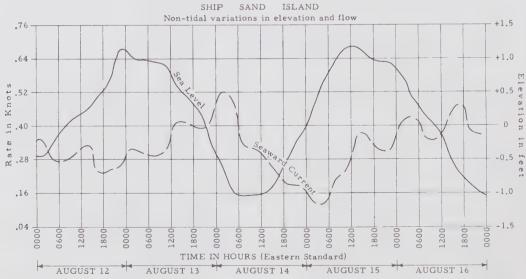


Figure 11 A surge observed in the Moose River during October, 1963.

which feed in the bay for a short time might choose to move to better feeding grounds while the local sea life which is already at a low level will be further reduced.

The probability of occurrence of surges will in no way be affected by the more regular outflow from the NBR complex.

The primary factors controlling surges are wind strength, pressure patterns and pressure fluctuations; changes in the density of the air and of the water affect them too but imperceptibly compared to the other factors.

6. Other sources of energy in Hudson Bay and Hudson Strait

Besides the obvious sources of hydroelectric power which it is now intended to tap around James Bay, another potential source of energy could be exploited in the New Quebec with a minimum of ecological upheaval: the very large tide in Ungava Bay. It is a little known fact that the tidal range in the vicinity of Leaf Basin in Ungava Bay is as large as in Minas Basin in the Bay of Fundy and that at some times of the year the tide in Ungava Bay is larger than in the Bay of Fundy. In Ungava Bay the tidal range increases southwestward and reaches its maximum in the vicinity of Leaf Basin and the Koksoak River. Not only is the tide very large but numerous fjords and basins exist in that area which would be most suitable as reservoirs for tidal energy. Not only do they have an appreciable storage capacity but the width of their mouth is relatively small so that construction of a dam across them would not be at all prohibitive. As well, the depths in the entrances are smaller than those which would have to be faced in the Bay of Fundy.

The exploitation of the tidal energy in these basins would not preclude in any way the conventional exploitation of the upper waters of the rivers emptying into some of them for the production of hydroelectric power. As a matter of fact the two modes of exploitation would rather complement each other since the construction of hydroelectric plants upstream would reduce the flow of sediments into the tidal basin and this would prolong the useful life of the tidal plant. Such an arrangement is not possible in James Bay because of the reduced range of the tide there.

The tidal power which can be extracted from some basins in Ungava Bay may be estimated using the formula (Godin, 1969)

$$P = \frac{\rho g S_0 [(2h_0)^2 + (1/3)e^2]}{T}$$
 (21)

where  $\rho$  = density of the water,  $S_0$  = the area of the basin,  $h_0$  = the amplitude of  $M_2$ , 1/2e = the amplitude of  $N_2$  +  $S_2$  +  $K_2$  and T = the tidal period, 12.4 hours. The equation (21) is an upper limit of the value of the power; in practice a tidal plant can extract at most 30 percent of this power. Table 4 lists the potential power output of some basins in Ungava Bay.

Table 4

Potential tidal power for some sites in Ungava Bay.

Basin	Surface area	<sup>M</sup> 2	S2 <sup>+N</sup> 2 <sup>+K</sup> 2	30 percent of potential tidal power
	10 <sup>7</sup> m <sup>2</sup>	m	m	Megawatts
Payne River Leaf Basin Koksoak River George River Abloviak Fjord	11.4 45.2 5.7 8.0 4.4	4.33 4.09 3.37	2.20 2.28 2.12 1.68 1.73	617 2450 274 258 57

The total assessed power amounts to 3650 MW and compares with the whole power potential of the Grande-Rivière. In addition, the tidal power output is highly reliable on a yearly basis (Bernshtein, 1965) while fluctuating from day-to-day and its yearly regularity could be used to compensate the yearly variabilit of the output of the rivers exploited around James Bay. This would involve integrating the network of the Ungava Bay project with the NBR, Grande-Rivière and Eastmain River projects; this would be most natural to do in any case, simply to take advantage of the networks already established.

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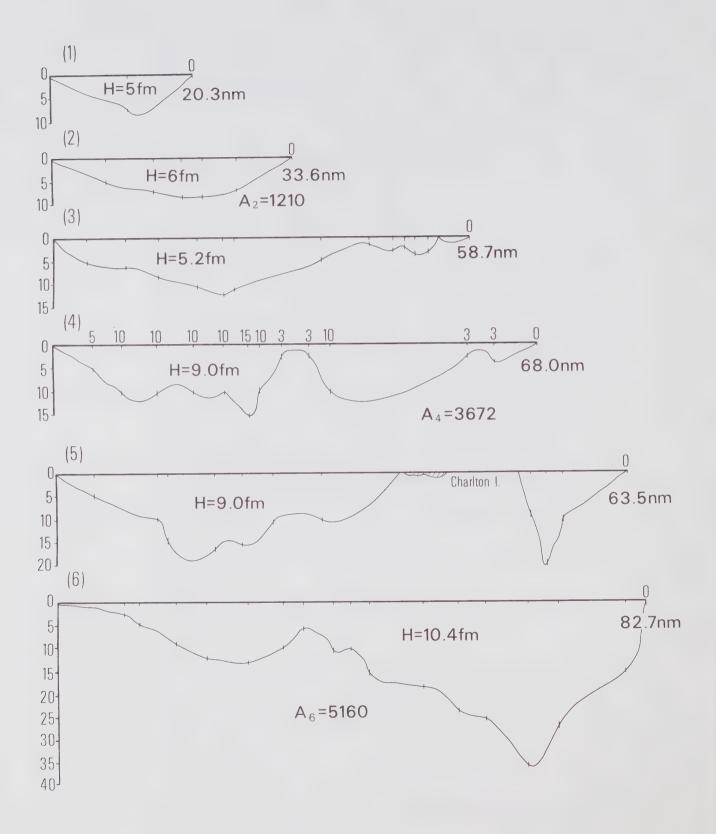
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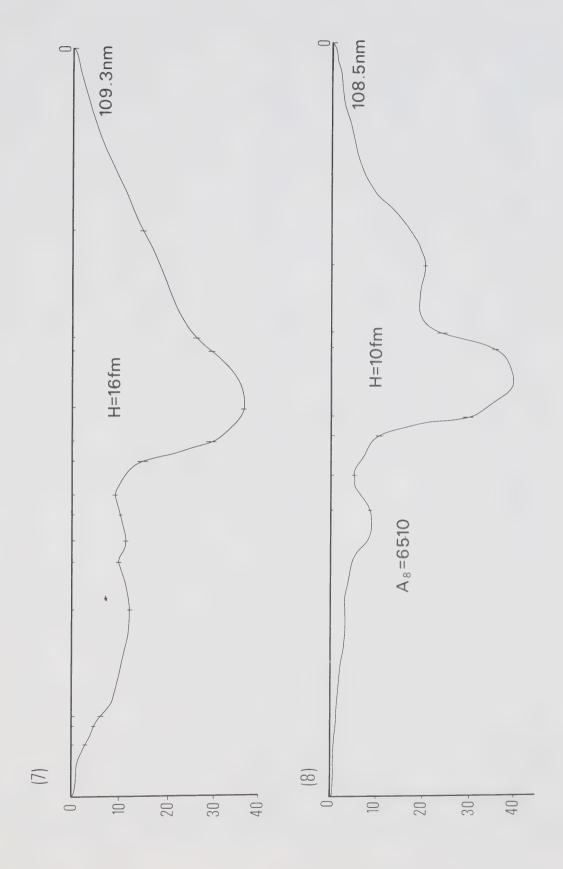
## 8. List of figures and tables

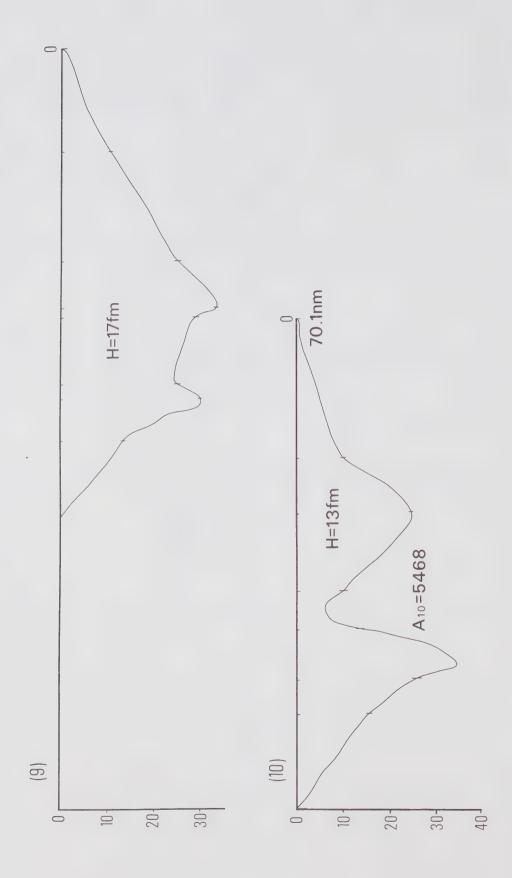
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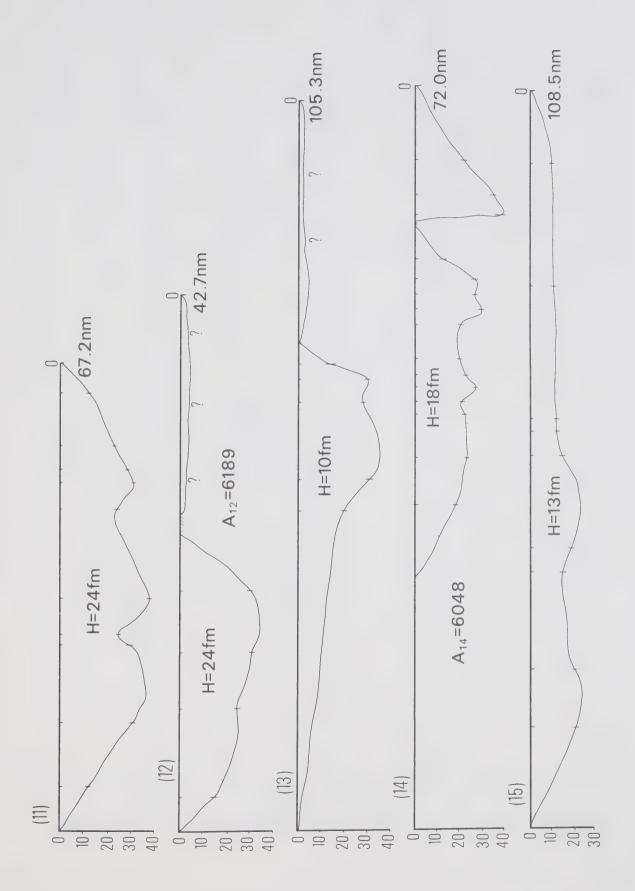
## 9. Appendix

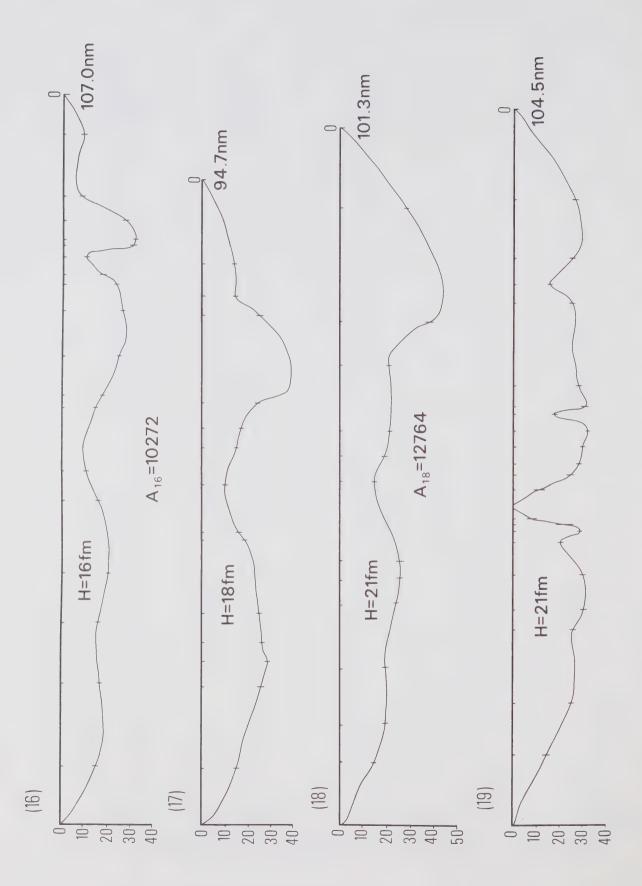
Pictorial presentation of the profiles derived from the schematization of the bay (Table 2).

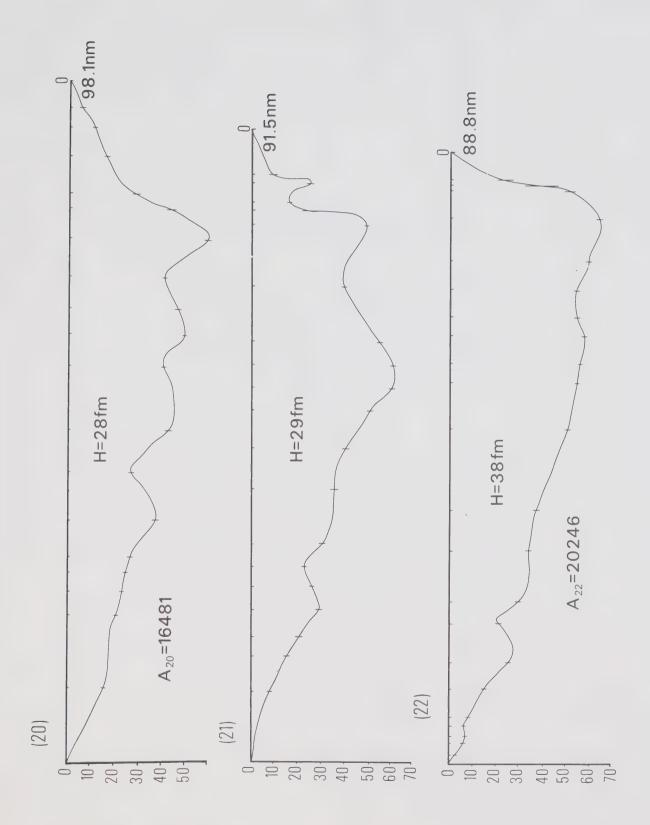


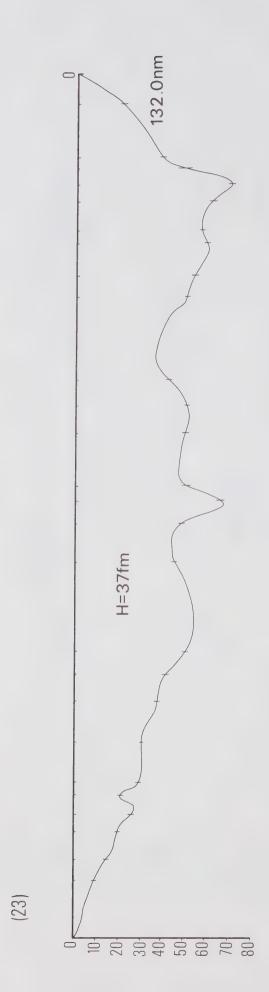












Circulation in James Bay

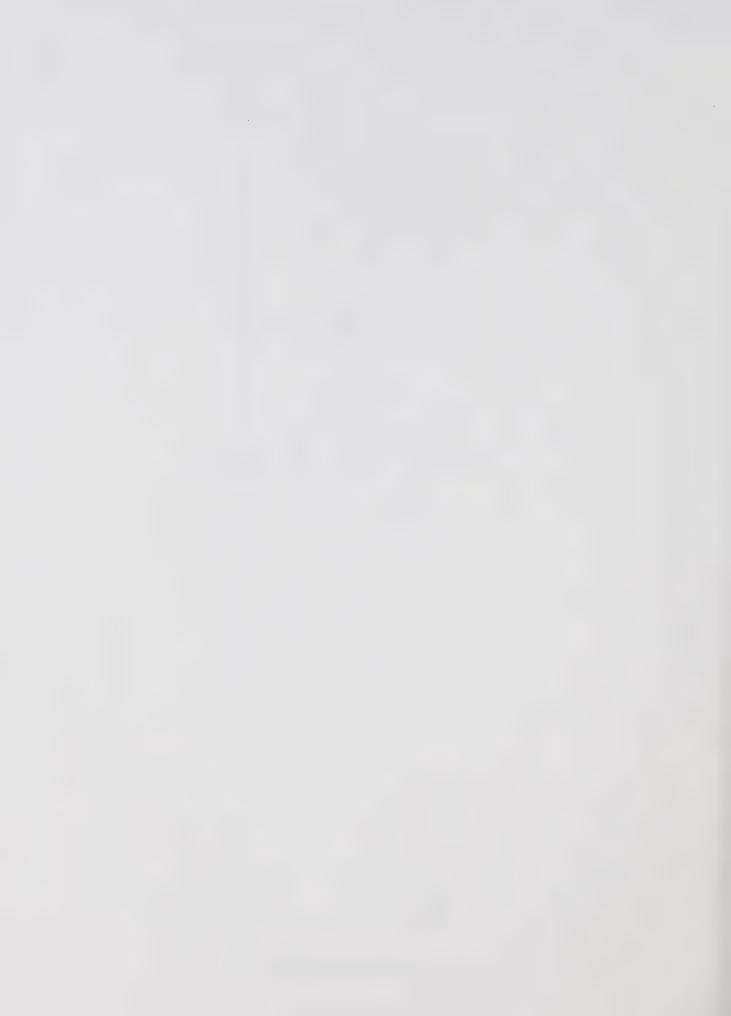
by

T.S. Murty

Marine Sciences Branch

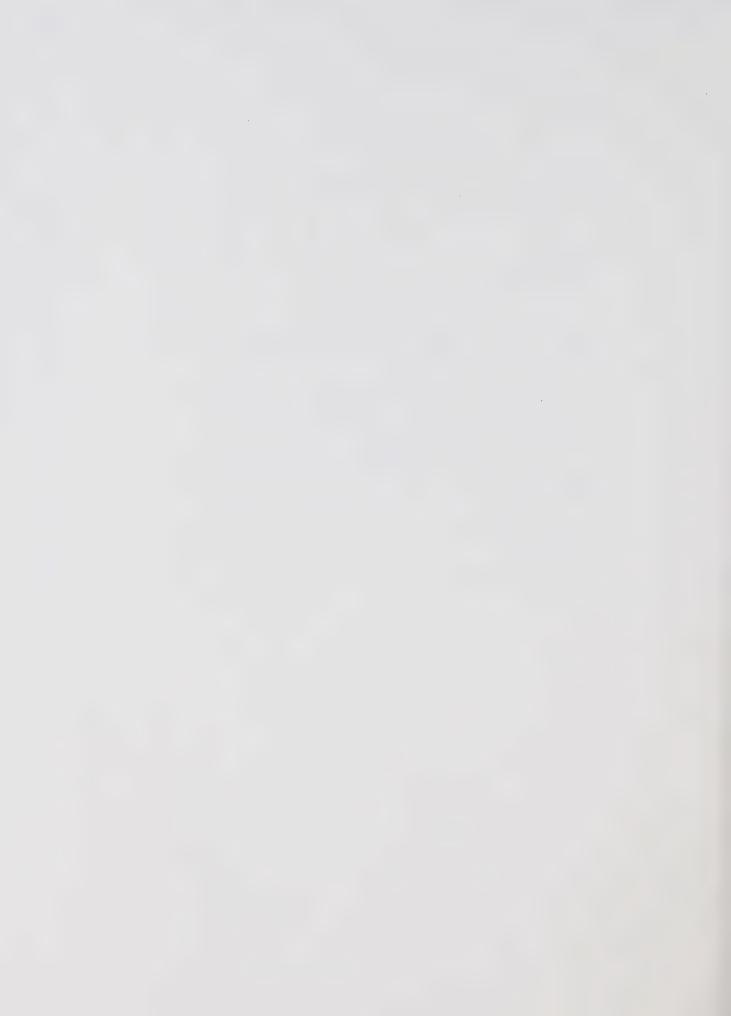
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#### 0. Abstract

This is a preliminary study on some physical oceanographic problems of James Bay. The interest in this study arose in connection with the proposed hydroelectric power project on some rivers draining into James Bay. The normal modes of response of James Bay are calculated using a one dimensional topographic model. Using the method of characteristics it has been shown that at the southern shores of James Bay the storm surge amplitude could reach up to 19 feet. It was shown that the thermohaline circulation in James Bay could be as intense as the wind generated circulation and this may explain some of the anomalous ice drift patterns. Although no barotropic coastal jets appear to be possible in James Bay, baroclinic coastal jets with a width of 7 kilometers (km) can occur with the core of the jet being at about 7 km seaward from the position where the thermocline intersects the The circulation in some rivers draining into James Bay is examined theoretically using the concepts of the three modes of estuarine circulation, namely the river-discharge mode, wind-stress mode and gravitational-convection mode. The question of atmospheric water balance over these river basins is briefly examined and some possible consequences of the man-made changes have been speculated.

#### 1. Introduction

This is a preliminary study of some physical oceanographic problems in James Bay. Because of the deadline imposed on the time available for this study no rigorous numerical modelling has been attempted. The present interest in this study arose from the recent announcement by the government of Quebec to divert the waters of certain rivers flowing into James Bay for hydroelectric power purposes.

Barber (1971) summarized some of the physical oceanographic problems in James Bay and its connected waters, namely Hudson Bay. He also pointed out the considerations that led to this present study. Godin (1971) studied the propagation of the tide into James Bay using a one dimensional topographic model. The topographic data used in the normal mode study has been kindly supplied to me by G. Godin and the input data on the estimates for the circulation problem has been supplied by F.G. Barber.

#### 2. The normal modes of James Bay

A knowledge of the normal modes of response of James Bay is useful in understanding the phenomena of tides, storm surges and circulation and for this reason these have been computed using a one dimensional topographic model. The linearized momentum and continuity equations for flow in a water body of varying cross-section are given by the so-called channel equations (Rao, 1968),

$$\frac{\partial \hat{M}}{\partial t} = -gA \frac{\partial \hat{\eta}}{\partial x} \tag{1}$$

$$\frac{\partial \hat{\eta}}{\partial t} = -\frac{1}{B} \frac{\partial \hat{M}}{\partial x} \tag{2}$$

where the x axis is locally tangent to the principal axis of the channel, t is time and g is gravity. The other parameters have the following meanings:

 $\hat{M}(x,t)$  = volume transport through a vertical section,

 $\hat{\eta}(x,t)$  = water level deviation from the undisturbed level,

A(x) = area of cross-section and

B(x) = surface width of the section.

In the above equations, friction and terms due to the earth's rotation are ignored. Later we will estimate the possible effects of these terms on the frequencies of the normal modes. For free oscillations of the systems assume,

$$\hat{M}(x,t) = M(x) \cdot \sin(\sigma t) \text{ and}$$
(3)

$$\hat{\eta}(x,t) = \eta(x) \cdot \cos(\sigma t)$$

Here M(x) and  $\eta(x)$  are the space dependent normal mode functions of James Bay and  $\sigma$  is the frequency. Substitution of (3) into (1) and (2) gives:

$$\sigma M = -gA \frac{d\eta}{dx} \tag{4}$$

$$\sigma \eta = \frac{1}{B} \frac{dM}{dx} \tag{5}$$

The boundary conditions are the following:

M is arbitrary and 
$$\eta=0$$
 at the mouth of James Bay, (6)

Equations (4) to (7) form an eigenvalue problem of the frequencies  $\sigma$ .

Define, following Rao (1968)

$$C_{i} \equiv \frac{\Delta x}{gA_{i}}$$
 and  $D_{i} \equiv -B_{i}\Delta x$  (8)

Then equations (4) and (5) become:

$$\eta_{i+1} = \eta_{i-1} - \sigma C_i M_i, i=2(2)22$$
 (9)

$$M_{i+1} = M_{i-1} - \sigma D_{i} \eta_{i}, i=3(2)23$$
 (10)

These finite difference forms are written with respect to a staggered grid in which M and  $\eta$  are computed at alternate points. The James Bay topography used here is the same as the one used by Godin (1971) in his calculation of the propagation of the tide. The grid distance  $\Delta x$  is 20 nautical miles, this being the distance either between two successive M sections or two successive  $\eta$ 

sections. In our model the grid numbering is such that, section one is the mouth and section twenty-four is the head of James Bay, opposite to the scheme of Godin (1971).

The equations (9) and (10) were solved by prescribing  $\eta_1=0$  and  $M_2=10^8$  cm<sup>3</sup> sec<sup>-1</sup> and alternately solving these equations. Since these are linear equations, any arbitrary value can be assigned to  $M_2$ . A trial value is provided for  $\sigma$  for the first mode and if  $M_24$  is not very close to zero, the value of  $\sigma$  is changed slightly and the computation is repeated. This is continued until that value of  $\sigma$  is found which makes  $M_{24}$  equal to zero for practical purposes. This value of  $\sigma$  gives the frequency of the fundamental mode. The same procedure can be used to calculate the frequencies of the higher modes.

The trial values for  $\boldsymbol{\sigma}$  for each mode can be calculated from the Merian formula:

$$T_{m} = \frac{4L}{m\sqrt{gH}}$$
 (11)

where  $T_m$  is the period of the mth longitudinal mode, L is the total length of James Bay and H is the average depth of James Bay. The values used were L=230 nautical miles and H=32 metres. Table 1 shows the periods of the normal modes of James Bay calculated both from the Merian formula and from the topographic model.

Table 1 Periods of the normal modes of James Bay.

Modal Number	Period Calculated from Merian Formula Hours Min		Period Calculated from Topographic Model Hours Min	
1	26	45	22	42
2	13	23	8	54
3	8	55	6	00
4	6	41	4	24
5	5	21	3	48
6	4	28	3	06

Figures 1(a) to 1(f) show the structures of these six longitudinal modes in terms of the modal functions M(x) and  $\eta(x)$ . In each, the abscissa denotes the grid point number with 1 denoting the mouth and 24 the head. The ordinate scale on the left side shows  $\eta$  and on the right side shows M. Figure 2 shows the positions of the nodes in James Bay computed from the topographic model.

In the above calculation the effect of the earth's rotation is ignored, which is a serious drawback considering the size of James Bay. The inclusion of rotation not only changes the frequency of the mode but also destroys the standing nature of the oscillation by introducing amphidromic systems into the modal structures. James Bay may be treated as a rectangular bay with length about three times the average width for purposes of estimating the effect of rotation on the frequency of a given mode. Rao (1968) mentions that Van Dantzig and Lauwerier (1960) gave the following formula for a rectangular bay with length twice the width,

$$\sigma = \sigma_0 + 0.504 \frac{f^2L}{2\pi C} + O(f^4)$$
 (12)

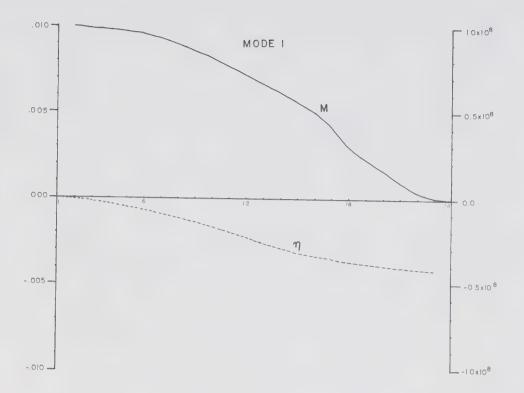
where  $\sigma_{\text{O}}$  is the lowest nonrotating frequency, given by,

$$\sigma_0 = \frac{\pi C}{2L}$$
 and

$$C = \sqrt{gH}$$
 (13)

and where the other symbols have the same meaning as before.

Figure 1 Structure of the first six longitudinal modes of James Bay. The abscissa shows the grid number (l is the mouth and 24 is the head). The ordinate scales on the left and right sides are for the water level \( \eta \) and the volume transport M respectively. Since this is a linear problem the actual units are arbitrary.



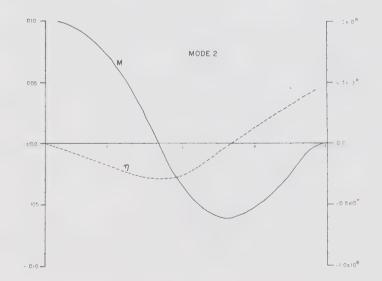
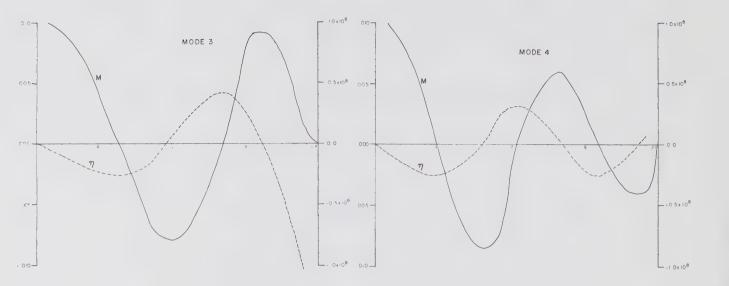


Figure 1



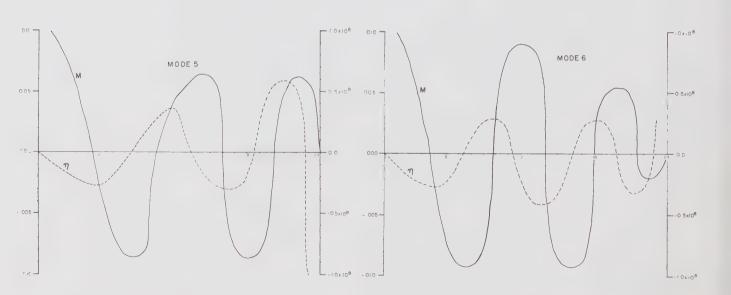


Figure 1 (con't)

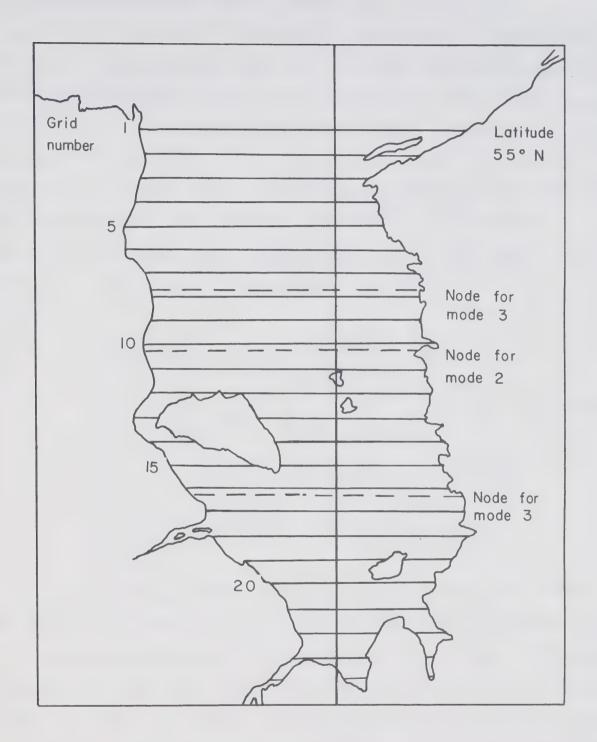


Figure 2 Nodal positions in James Bay calculated from the topographic model.

The equation gives a result that is exactly opposite to what happens in a completely closed basin. That is, in the case of a rectangular bay, rotation increases the frequency of the lowest mode. Since in our case (Table 1) the period of the fundamental mode is 22 hours 42 minutes then,

$$\sigma_0 = \frac{2\pi}{22.7 \times 3600} \text{ sec}^{-1} \tag{14}$$

Since 
$$\frac{L}{C} = \frac{\pi}{2\sigma_0} = \frac{22.7 \times 3600}{4}$$

$$= 20,430 \text{ sec}$$

From (12), (13)

$$\frac{\sigma L}{C} = \sigma_0 \frac{L}{C} + 0.504 \frac{f^2}{2\pi} \left(\frac{L}{C}\right)^2$$
$$= \frac{\pi}{2} + 0.335$$

The nondimensional period,  $T = 2\pi/(\sigma L/C)$  is therefore 3.3, while in the nonrotating case it is 4. Thus rotation can reduce the period by 17.5 percent for a rectangular bay with length and fundamental mode period of James Bay and with width equal to half the length of James Bay. Since, in reality, James Bay is about three times as long as it is wide, the reduction of the period due to rotation would be less than 17.5 percent.

#### 3. Storm surge estimation for James Bay

Thompson (1968) summarized the climate of Hudson Bay. He stated that the upper air circulation over Hudson Bay is the persistent counter-clockwise air flow around a low pressure vortex over northern Baffin Island in winter. This gives rise to a general transport of cold air in a northwest to southeast direction in winter. Thompson further stated that since many of the storm centers lack sufficient moisture to cause heavy precipitation, their principal effect is in inducing strong north winds over Hudson Bay as they travel across it. He further stated:

"In contrast to the broad expanses of Arctic tundra that surround Hudson Bay the sub-Arctic lands bordering James Bay are partially forested and thus protected from strong winds. As a result two of the most distinguishing features of Hudson Bay winter climate - wind chill and blowing snow - are not nearly as evident near James Bay. With this important exception and the fact that James Bay is several degrees of latitude farther south, the factors that influence the climate of James Bay are essentially the same as those outlined for Hudson Bay".

Archibald (1969) studied the storm tracks over Hudson Bay and eastern Canada. As can be seen from his diagrams the storms move from different directions over James Bay. Since my present interest is in the maximum possible amplitudes of storm surges in James Bay, I will concern myself mainly with storms moving from north to south over James Bay thereby piling up water on the southern shore and causing large surges. Since the magnitude of the surge is strongly dependent upon the nature of

the coast of James Bay I will briefly summarize some relevant facts from the article by Robinson (1968).

It can be seen from Figure 1 of Robinson's article that the eastern half of James Bay falls into the classification of "East Coast Uplands", while the western half falls into the category "South Coast Lowlands". Regarding the south coast lowland, I quote from Robinson (1968):

"Along the whole coast there is a flat strip five to ten miles wide, with the widest parts generally being to the north. The coastal zone is treeless, but grass or marshes are common. Storm beaches, a few feet high, are the only topographic features. Tidal flats may be exposed for one to six miles; even at high tide shallow water extends far offshore. ...Deeper water may be found at the drowned river mouths, but shifting sand bars and minor deltaic deposits are navigation hazards".

Regarding the east coast upland I again quote Robinson (1968):

"The east coast and adjoining interior have a combination of landform features which are different from those west of Hudson and James bays, but in some characteristics they have regional similarities. The central and northern sections of the east coast are high and rugged, with rocky hills and drift-covered uplands inland, as is characteristic of the Northwest Hills; the southern section along James Bay is a poorly-drained lowland, although not as wide as that on the west side."

Because of the time limitation of this study, no numerical modelling with topography taken into account has been attempted. Instead, use was made of the analytical solutions developed by Rao (1969) for a rectangular bay of uniform depth, making use of the method of characteristics. In this study

James Bay will be treated as a rectangular bay of length 230 nautical miles, of width up to 100 nautical miles and of average depth 32 meters, connected to Hudson Bay at the mouth. James Bay can be considered narrow because its length is about three times its average width and a one dimensional model is used here and the effect of earth's rotation is suppressed. Since we are interested in the transient response of the bay to a time and space-dependent wind field, the bottom friction is not included. It can be shown by order of magnitude considerations that the nonlinear term in the momentum equation can be omitted. Hydrostatic assumption has been made for the pressure field. For convenience, the relevant equations and the solutions will be summarized here.

Let  $\bar{x}$ ,  $\bar{z}$  represent a cartesian coordinate system such that the  $\bar{x}$ -axis is oriented along the length of James Bay with  $\bar{x}$ =0 at the mouth and  $\bar{x}$ = $\bar{L}$  at the southern shore. The  $\bar{z}$ -axis is positive upwards and is measured from the undisturbed position of the water surface. Then the bottom is at  $\bar{z}$ =- $\bar{h}$ . Here, superior bar denotes dimensional quantities. The vertically integrated forms of the momentum and continuity equations are:

$$\frac{\partial \overline{M}}{\partial \overline{t}} + \overline{g} \overline{h} \frac{\partial \overline{\eta}}{\partial \overline{x}} = \overline{R}$$
 (15)

$$\frac{\partial \bar{\eta}}{\partial \bar{\tau}} + \frac{\partial \bar{M}}{\partial \bar{x}} = 0 \tag{16}$$

where

$$\bar{M} = {}^{\circ}\int_{\bar{h}} \bar{U} d\bar{z}$$
 (17)

is the volume transport through a vertical section,  $\bar{U}$  being the horizontal velocity in the  $\bar{x}$  direction and  $\bar{\eta}$  is the deviation of the water level from its mean position,  $\bar{g}$  is gravity,  $\bar{t}$  is time and  $\bar{R}$  is the force due to wind stress given by:

$$\bar{R} \equiv \frac{\bar{\tau}}{\bar{\rho}}$$
 (18)

The boundary conditions are the following:

$$\bar{M}=0$$
 at the head of James Bay, i.e. at  $\bar{x}=\bar{L}$ , (19)

$$\bar{\eta}=0$$
 at the mouth, i.e. at  $\bar{x}=0$ . (20)

Initially, 
$$\bar{M}=0$$
,  $\bar{\eta}=0$  at  $\bar{t}=0$  for all  $\bar{x}$ . (21)

Rao (1969) introduced the following nondimensionalization:

$$x \equiv \frac{\bar{x}}{\bar{L}}$$
  $\eta \equiv \frac{e^2}{\bar{R}_0 \bar{L}} \bar{\eta}$   $R \equiv \frac{\bar{R}}{\bar{R}_0}$  (22)

$$t \equiv \bar{t} \frac{\bar{c}}{\bar{L}} \qquad \qquad M \equiv \frac{\bar{c}}{\bar{R}_0 \bar{L}} \bar{M}$$

where 
$$\bar{c} = \sqrt{\bar{g} \bar{h}}$$
 (23)

and  $\bar{R}_0$  is a scale value of the wind stress. Using this scheme, equations (15) and (16) can be nondimensionalized. From addition and subtraction of these we get:

$$\frac{d}{dt} (M \pm \eta) = R \text{ for } \frac{dx}{dt} = \pm 1$$
 (24)

This states that the quantity M  $\pm$   $\eta$  is conserved along the positive and negative characteristics  $dx/dt = \pm 1$ , respectively. Since the bay is assumed to have a uniform depth  $\bar{h}$ , the positive and negative characteristics are straight lines given by  $x = \pm t + const$ .

I will consider two different types of stress bands and assume that the wind stress force R is in the form of a step function and moves with a constant speed  $\bar{V}$  starting from x=0 at t=0. The nondimensional form of this translational speed is

$$V \equiv \frac{\overline{V}}{C} \tag{25}$$

The value of R at any given point in the bay is given by

$$R = \begin{cases} 1 & \text{for } t \ge \frac{x}{V} \\ 0 & \text{for } t \le \frac{x}{V} \end{cases}$$
 (26)

This stress band crosses the head of the bay (x=1) at t=1/B and in the case of the semi-infinite band, for t>1/V the entire bay is under the influence of the wind stress. In the case of the

finite stress band we assume a pulse of a square wave shape with zero forcing ahead and behind the pulse and the width of the pulse being V(t-t') = VT = X where T is the time taken by the pulse to travel past a given point in the bay. The forcing function for the finite stress band can be obtained by superimposing on (26) another semi-infinite band R'(x',t') where

$$R'(x',t') = \begin{cases} 0 & \text{for } t' \leq \frac{x}{V} \\ -1 & \text{for } t' \geq \frac{x}{V} \end{cases}$$
 (27)

and t'  $\leq$  t. Hence, the total forcing function for the case of finite band width is

$$R+R' = \begin{cases} 0 & \text{for } t \leq \frac{x}{V} \\ 1 & \text{for } t \geq \frac{x}{V} \geq t' \end{cases}$$

$$0 & \text{for } t' \geq \frac{x}{V} \end{cases}$$

$$(28)$$

Since our equations are linear, the solutions due to forcing from R' can be obtained from the solutions due to forcing from R by changing the sign of  $\eta$  and replacing t with t' = t - T.

Rao showed that the maximum possible elevation at the head is given by:

$$\bar{\eta}_{\text{max}} = \frac{2 \bar{L}}{\bar{g} \bar{h}} \bar{R}_0 \tag{29}$$

He used the following representation for  $\bar{R}_{\text{O}}$ 

$$\bar{R}_0 = 4 \times 10^{-6} |W|W$$
 (30)

where W is the wind speed (ft  $\sec^{-1}$ ) along the axis of the bay. If we take W=75 ft  $\sec^{-1}$ , corresponding to 51 mph we get

$$\bar{\eta}_{\text{max}} = \frac{2 \times 230 \times 6080}{32 \times 32 \times 3.281} \times 4 \times 10^{-6} \times 75^{2}$$

= 18.8 feet

For higher winds speeds the amplitude of the surge is even higher. The reason for the possibility of such large storm surges in James Bay is the long fetch and the relatively shallow water in the bay.

Thus the storm surges appear to attain very large amplitudes at the southern end of James Bay when intense storms pass over it. This result coupled with the fact that the shores are very flat could give rise to very serious storm surge inundation.

# 4. Circulation in James Bay

The circulation in James Bay cannot be studied completely independent of the circulation in Hudson Bay to which it is connected at its widest portion. As is expected for any large water body in the northern hemisphere, the circulation in Hudson Bay is dominated by a large counterclockwise cell (Murty and Yuen, 1970). Figure 3 shows the circulation in September

(left side) and in May (right side) calculated by these authors using wind stress computed through the balance equation. Figure 4 shows the corresponding water level deviations. These calculations were made using a steady state topographic model and pressure data averaged over several years has been used in the wind stress computation. It can be seen from Figures 3 and 4 that while the circulation itself is extremely weak and not too well defined in the southern portion of Hudson Bay and James Bay, the water level gradients in the equilibrium state are strong in James Bay (relative to Hudson Bay).

The Ice Central Office at Halifax which is responsible for predicting the movement of ice has faced some difficulty in their prediction because of anomalous drift patterns,

"...The summer of 1968 provided another such situation. The general water circulation pattern is believed to be a cyclonic drift around the Bay. This would normally be expected to carry ice from the Cape Churchill area southeastward toward James Bay. A west to northwest wind accompanying such drift should move the ice southeastward rather quickly...the combination of wind and water current drifts is almost negligible.

It must be assumed, therefore, that during the melting period thermohaline processes occur which tend to disrupt the normally weak circulation in the southern bay. The addition of fresh water through runoff from the rivers along the west and south coasts also tends to disrupt the salinity balance of the southern section of the Bay. However, the failure of ice to drift southeastward under the influence of persistently favourable winds leads again to the idea of a countercurrent in the southern part of the Bay, especially during the period when melting ice is present.

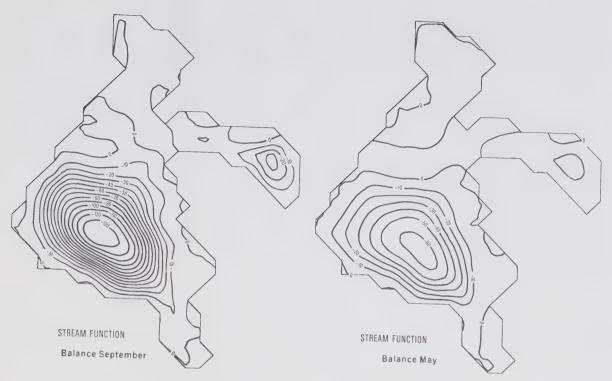


Figure 3 Stream function (for the volume transport) in units of  $10^{12}~\rm cm^3~\rm sec^{-1}$  for September (left) and May (right) using data averaged over 30 years.

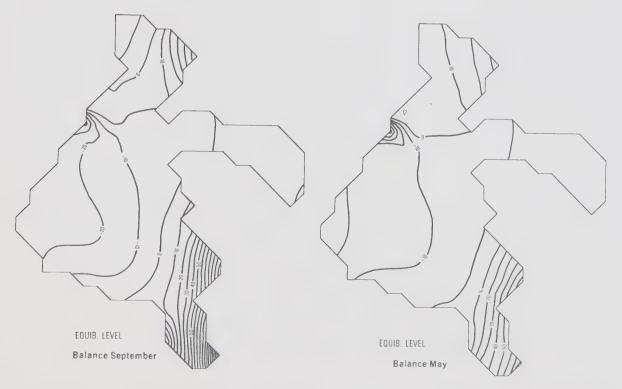


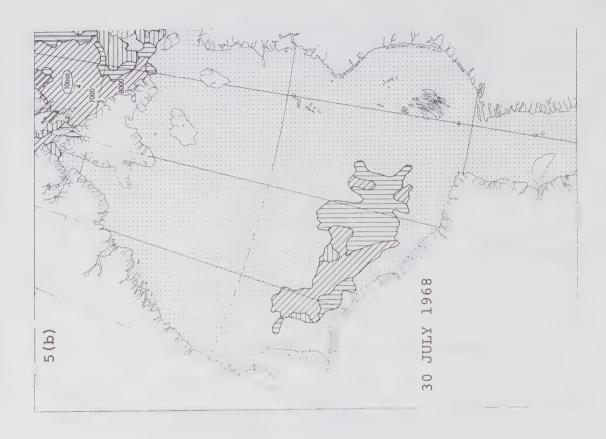
Figure 4 Water level deviation in units of  $10^{-2}$  cm for September (left) and May (right) corresponding to Figure 3.

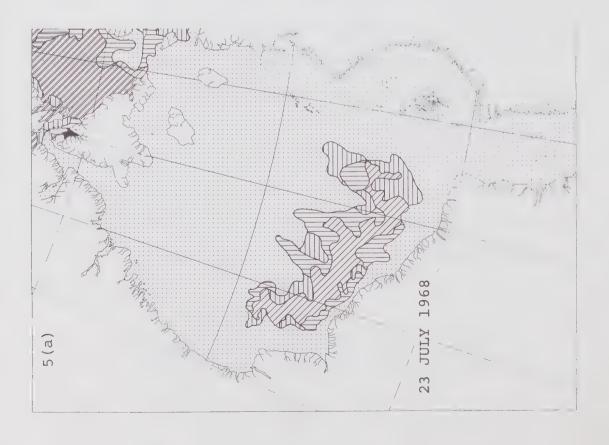
It is obvious that a serious study of thermohaline relationships to water circulation in this area is necessary before useful forecasts can be provided for southern Hudson Bay". (Anon, 1970).

Figures 5(a) to 5(f) show a sequence of diagrams in which the ice conditions in Hudson Bay are shown on different days. These figures are respectively the figures 12, 14, 16, 18, 20 and 22 in the Ice Central Report for 1968. It can be seen from these that the ice simply melted away and did not move with the water current pattern.

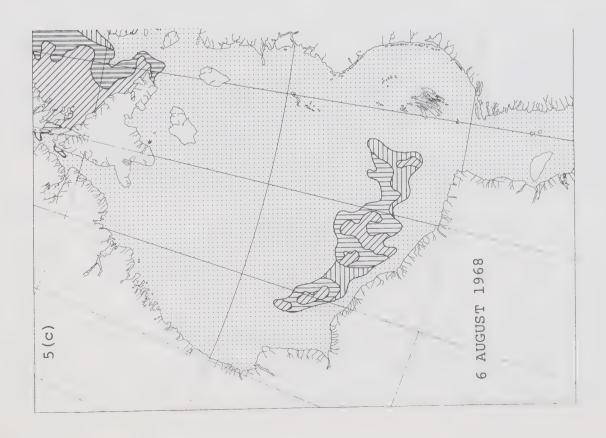
It is clear that an estimation of the intensity of the thermohaline circulation is essential, especially to answer the question whether this could be strong enough to offset the wind-generated circulation. Another consideration is that because James Bay really is part of the complex consisting of Hudson Bay, Hudson Strait and Foxe Basin, this system should be studied together and any estimates made should involve the length and time scales representative of this system as a whole. The north-south extent of this system is large enough for  $\beta$ , the variation of the coriolis parameter with latitude to play a significant role. Because of the time limitations on this study no numerical model has been developed. Instead, estimates have been made using an analytical study made by Gates (1968).

"The action of surface wind stress on an underlying ocean has been extensively studied since the pioneering work of Ekman (1905). The characteristic depth Figure 5 Ice conditions in Hudson Bay for the winter period of 1968-69 on six different dates.

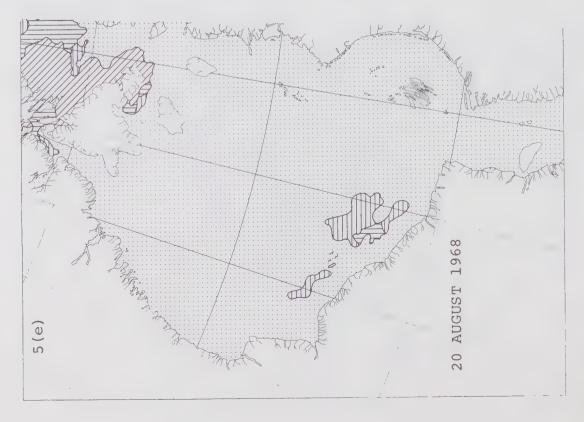












of penetration of wind-induced motion (Ekman depth) depends upon both the vertical eddy diffusion and the coriolis parameter, and the net transport in deep waters is normal to the surface wind stress with the horizontal velocity displaying the well known spiral structure. The modifications of these features in shallow water and near coast lines were also explored by Ekman by the introduction of a bottom boundary layer whose transport complements that of the surface laver. addition to describing the horizontal transport in terms of the wind stress, the Ekman theory provides an estimate of the net vertical flux required beneath the surface boundary layer by mass continuity. This is the so-called Ekman vertical velocity."

Gates (1968) studied at first, a homogeneous ocean of uniform depth on a  $\beta$ -plane and then considered stratification. His assumptions include steady-state, neglect of inertial terms and horizontal friction. The homogeneous ocean is divided into an interior region in between two boundary layers, one at the top and one at the bottom. It should be cautioned that because of the shallow nature of the Hudson Bay System it is debatable whether there really is an interior region where the vertical eddy viscosity is negligible.

The characteristic thickness E of the Ekman boundary layer is:

$$E = \sqrt{\frac{2v_z}{f}}$$
 (31)

where f is the coriolis parameter and  $\nu_Z$  is the vertical eddy viscosity. Gates defined E slightly differently from the conventional definition of Ekman depth, the difference being,

here E is thickness at which the speed is reduced by the factor  $e^{-1}$  and this is  $1/\pi$  times the conventional definition. Since f increases with increasing latitude, E decreases with increasing latitude provided  $\nu_Z$  does not change. Another assumption involved in Gates' study is that the thicknesses of the boundary layers are small compared to the total depth H (assumed uniform).

The equilibrium pressure field P is determined from the following relation:

$$E\nabla^2 P + \frac{\beta}{f}$$
 (2H-E)  $\frac{\partial P}{\partial x} - \frac{\beta E}{f} \frac{\partial P}{\partial y}$ 

$$= 2 \operatorname{curl}_{z} \tau + \frac{\beta}{f} \tau_{sx}$$
 (32)

where  $\tau$  is the vector wind stress with x and y components  $\tau_{SX}$  and  $\tau_{SY}$ . It can be seen that, when  $\beta=0$  this reduces to

$$\nabla^2 P = \frac{2}{E} \operatorname{curl}_{Z} \tau \tag{33}$$

which was originally derived by Ekman (1923). The vertical velocity at the bottom of the surface boundary layer is

$$W_1 = \frac{1}{\rho_0 f_0} \operatorname{curl}_z \tau \tag{34}$$

which agrees with the Ekman vertical velocity. Here,  $\rho_0$  is the density of the homogeneous fluid and  $f_0$  is the coriolis parameter at the standard latitude such that

$$f = f_0 + \beta y \tag{35}$$

The vertical velocity in the interior region is given by

$$W_{I} = W_{K} + \left(\frac{\beta}{f^{2}\rho_{0}} \frac{\partial P}{\partial x}\right) z \tag{36}$$

where  $W_K = W_K(x,y)$  is an arbitrary function to be determined knowing the wind stress. In the absence of  $\beta$ , this is uniform in depth.

Next, Gates introduced a surface temperature distribution and considered its effects on the Ekman vertical velocity. This temperature distribution is assumed to be maintained by surface heat exchanges and Gates assumed that the horizontal temperature gradient decreases linearly with depth, from a maximum at the surface, z=H, to zero at the bottom, z=0. Thus if T is the temperature field, then

$$\frac{\partial \mathbf{T}}{\partial \mathbf{x}} = \frac{\mathbf{z}}{\mathbf{H}} \left( \frac{\partial \mathbf{T}}{\partial \mathbf{x}} \right)_{\mathbf{0}} \tag{37}$$

$$\frac{\partial \mathbf{T}}{\partial \mathbf{y}} = \frac{\mathbf{z}}{\mathbf{H}} \left( \frac{\partial \mathbf{T}}{\partial \mathbf{y}} \right)_{\mathcal{O}}$$

where the subscript o denotes the value at the surface. Let the equation of state be

$$\rho = \rho_0 (1 - \alpha T) \tag{38}$$

where  $\alpha$  is the coefficient of thermal expansion of water. That

is, here we ignored the effect of salinity variations on the density. However, qualitatively speaking it will add one more term in equations (39) and (40). In (38), T is the temperature relative to a reference value corresponding to  $\rho=\rho_0$ .

Assuming that  $\eta << H$  (where  $\eta$  is the perturbation of the free surface) and  $\rho \approx \rho_0$  when  $\eta$  and  $\rho$  are undifferentiated we can write the following expressions for the interior geostrophic velocity components

$$U_{I}(z) = -\frac{g}{f} \frac{\partial \eta}{\partial y} + \frac{\alpha g(H^{2} - z^{2})}{2fH} \left(\frac{\partial T}{\partial y}\right)_{0}$$
 (39)

$$V_{I}(z) = \frac{g}{f} \frac{\partial \eta}{\partial x} - \frac{\alpha g (H^{2} - z^{2})}{2fH} \left( \frac{\partial T}{\partial x} \right)_{0}$$
 (40)

Thus the presence of a temperature gradient permits the reversal of the current with depth.

The surface boundary conditions are now

$$\frac{\partial \mathbf{U}}{\partial \mathbf{z}} = \frac{\tau_{\mathbf{SX}}}{\rho_{0} \nu_{\mathbf{z}}} + \frac{\alpha \mathbf{g}}{\mathbf{g}} \left( \frac{\partial \mathbf{T}}{\partial \mathbf{y}} \right)_{0} \tag{41}$$

$$\frac{\partial V}{\partial z} = \frac{\tau_{SY}}{\rho_{O} v_{z}} - \frac{\alpha g}{f} \left( \frac{\partial T}{\partial x} \right)_{O} \tag{42}$$

both at z=H. The vertical velocity in the interior is given by

$$W_{I}(z) = \frac{\beta g}{f^{2}} \frac{\partial \eta}{\partial x} z - \frac{\beta \alpha g}{2f^{2}H} \left( \frac{\partial T}{\partial x} \right)_{O} \left( H^{2}z - \frac{z^{3}}{3} \right) + W_{K}$$
 (43)

Thus, in the non-homogeneous case the vertical velocity varies as the cube of the depth in the interior while in the homogeneous case it varies linearly. The equation to determine  $\eta$  for the non-homogeneous case is given by

$$\text{EV}^2 \eta \ + \ \frac{\beta}{\text{f}} \ \text{(2H-E)} \ \frac{\partial \eta}{\partial x} \ - \ \frac{\beta \text{E}}{\text{f}} \ \frac{\partial \eta}{\partial y} \ = \ \frac{2}{\text{g}\rho_0} \ \text{curl}_z \ \tau \ + \ \frac{\beta \tau_{\text{SX}}}{\text{gf}\rho_0} \ + \ \alpha \text{E} \left(\frac{\text{H}}{2} \ - \ \text{E}\right) \ \nabla^2 T_0$$

$$+ \frac{\alpha \beta H}{f} \left( \frac{2H}{3} - \frac{E}{2} \right) \left( \frac{\partial T}{\partial x} \right)_{0} + \frac{\alpha \beta E}{f} \left( \frac{3E}{2} - \frac{H}{2} \right) \left( \frac{\partial T}{\partial y} \right)_{0}$$
(44)

NOTE: equation (32) gives a corresponding equation for the homogeneous case.

It can be seen that when  $\beta=0$  this reduces to

$$E\nabla^{2}\eta = \frac{2}{g\rho_{0}} \operatorname{curl}_{z} \tau + \alpha E \left(\frac{H}{2} - E\right) \nabla^{2}T_{0}$$
(45)

Thus the effect of a local excess of temperature (for example,  $\nabla^2 T_0 < 0$ ) is analogous to that of the curl of the anticyclonic wind stress (curl<sub>z</sub>  $\tau$  < 0).

Without giving any more details we will simply state that starting with (45) Gates showed that the thermally forced contribution could be comparable to the wind-induced portion.

## 5. Possibility of coastal jets in James Bay

The dimensions of James Bay (length 230 nautical miles and average width of 100 nautical miles) are comparable to the Great Lakes. James Bay is very shallow (with an average depth of 32 meters) like Lake Erie. The main difference between James Bay and the Great Lakes, besides the saline water in the former and fresh in the latter, is that James Bay has a very wide mouth and

is connected to Hudson Bay while the Great Lakes are essentially closed systems with only some interconnecting rivers. Csanady (1967, 1968a, 1968b) studied the coastal jets in the Great Lakes and his concepts (Csanady, 1971) could be applied here to study the possibility of coastal jets in James Bay:

"Theoretical studies of some simple model Great Lakes (Csanady, 1967, 1968a,b) have suggested that shoreward wind drift may set up concentrated boundary currents of essentially zero potential vorticity in the Great Lakes by the mechanism of vortex stretching. Along an infinite vertical shore, as shown by Charney (1955) in an extension of Rossby's (1938) work on geostrophic adjustment, some coastal jets may be generated in this manner, although these would possess practical significance only in the baroclinic modes. In a closed basin of constant depth, the implusive application of wind stress leads, by the same mechanism, to the generation of slow baroclinic Kelvin waves (Csanady, 1968b), the profile of which, perpendicular to shore, is very similar to the steady-state coastal jets.

The application of these previous results to the real Great Lakes or other similar bodies of water is not immediate for at least two reasons: 1) depth variations in the shore zone are certain to modify the vortex stretching mechanism; and 2) a linear, frictionless theory (from which the above results were obtained) can hardly be expected to describe water movements accurately, particularly in the shallow coastal zones."

Csanady (1971) at first considered barotropic boundary currents near a sloping shore. He postulates that a boundary current is generated by the arrival of water in the shore zone, either driven by wind or due to the passage of a long slow wave. Csanady shows that no narrow boundary currents are possible in the barotropic case. Then Csanady considered a two-layer system

and showed that the analysis for the barotropic case with a modified scale factor can be used for the baroclinic case. Let  $\rho$  and  $\rho'$  be the densities of the top and bottom layers and h and h' be the depths. Then  $\epsilon$  =  $(\rho'$  - $\rho$  )/ $\rho'$  is the fractional density defect. The top layer depth h is assumed to be uniform while the depth h' of the bottom layer is variable and is assumed to be

$$h' = \begin{cases} sy \text{ for } y \leq y_0 \\ h'_0 \text{ for } y \geq y_0 \end{cases}$$

$$(46)$$

where y denotes a coordinate normal to the shore and it is assumed that conditions are uniform along the shore (for example,  $\partial/\partial x = 0$ ). Here, s is the beach slope.

Csanady showed that for this problem there are two important parameters, K and L where

$$L \equiv \left(\frac{g\epsilon h}{f^2}\right)^{1/2} \tag{47}$$

and

$$K = \frac{f}{2s} \left( \frac{h}{g\varepsilon} \right) 1/2 \tag{48}$$

Next he defines

$$S = \frac{sg\varepsilon}{f} \tag{49}$$

Here S is the slope length scale for effective gravity gs in the baroclinic case. Thus for the baroclinic case, there are two length scales S and L. A key parameter of this problem is  $\kappa$  defined in (48) which is also the same as

$$\kappa = \frac{L}{2S} \tag{50}$$

For James Bay, we can take

$$s = 7 \times 10^{-4}$$
 $\epsilon = 10^{-2}$ 
 $h = 5 \text{ meters}$ 
 $f \sim 10^{-4} \text{ sec}^{-1}$ 
(51)

From (48) we get

$$K = 1/2$$

$$L = S = 7 \text{ km}$$
(52)

Thus, the width scale of the boundary current L for James Bay is 7 km. Also the peak current occurs at a distance of order L from the point where the thermocline intersects the bottom.

#### 6. Some miscellaneous topics

Here I will discuss the estuarine circulation in the rivers that drain into James Bay. Because of strong tidal action the saline water intrudes into the rivers and this might give rise to the type of estuarine circulation discussed by Hansen

and Rattray (1965). They derived a set of coupled partial differential equations to study the circulation and salt flux processes for estuaries in which turbulent mixing results primarily from tidal currents and they separated the circulation into modes analogous to the barotropic, baroclinic and Ekman modes of oceanic circulation. These authors state that their solutions (although derived for the Delaware river estuary) hold well for sea straits and narrows having strong tidal currents and well-defined water and density budgets.

The present analysis may be of some use in understanding the dispersion processes and circulation in those regions of the rivers that are affected by strong tidal action from James Bay.

Hansen and Rattray divided the estuary into the three regimes, namely the outer, central and inner regimes, the outer being the closest to the connecting sea. For the circulation in the central regime they derived the following relation using similarity theory.

$$\phi(\eta) = \frac{1}{2} \underbrace{(2 - 3\eta + \eta^3)}_{I} - \frac{T}{4} \underbrace{(\eta - 2\eta^2 + \eta^3)}_{II}$$

$$-\frac{vR_{a}}{48}\underbrace{(\eta - 3\eta^{3} + 2\eta^{4})}_{TTT}$$
 (53)

where  $\phi$  is a stream function,  $\eta$  is a non-dimensional vertical coordinate, T is a non-dimensional wind-stress,  $\nu$  is a constant and  $R_a$  is an estuarine Rayleigh number. Thus, equation (53) expresses the circulation in the central regime as the sum of

three modes: the river discharge mode (term I), the wind-stress mode (term II) and the gravitational-convection mode (term III). The important result is that only the river-discharge mode allows a net transport of water. In deriving equation (53) it has been assumed that, in the central regime the width and depth of the river as well as the river discharge are constant. This assumption would be more questionable both in the outer and inner regimes. In the next section (Section 7) we will speculate on the relative influences of these three modes in some of the rivers and how the situation may change due to the proposed hydroelectric power project.

Next we consider the question of atmospheric water balance. A study incorporating this along with the hydrology of the various river basins before and after the man-made changes will help to evaluate the effects of the man-made changes. At present, although in principle, the study could be carried out for the present situation, because of the necessity to handle a large amount of data and because of the time limitations on this study, this part of the study is not carried out. However, I will outline a procedure described by Rasmussen (1970) and speculate on this in the next section. Rasmussen has developed the following model to study the atmospheric water balance and hydrology of the upper Colorado River Basin. Since any climatic effects are intimately connected with the atmospheric water balance, the following procedure developed by Rasmussen to study the hydrometeorology of a given river basin using data on precipitation and evaporation will be especially useful. I quote from Rasmussen:

"Traditionally, studies of the hydrologic balance of river basins have been approached from the viewpoint of the terrestrial part of the hydrologic cycle. The factors determining the runoff from an area are precipitation, evaporation, change in ground-water storage, and underground seepage from the basin. Such an approach to the study of hydrologic problems is often plaqued by measurement deficiencies ... If a large mountainous region is studied. the measurement problem is maximized because not only is the density of observations low, but the observations are biased toward the lower elevations where precipitation is lower.

Alternately, the atmospheric part of the hydrologic cycle may be studied to evaluate the net deposition of water over an area. A budget parallel to that of the terrestrial part of the hydrologic cycle must be observed. The atmospheric water balance may be expressed as the evaporation minus precipitation occurring at the ground over an area balanced by the net transfer of water mass through the atmospheric volume over the area and by the change in storage of water mass within the atmospheric volume. In theory then, given a continuous distribution in time and space of the atmospheric water mass, an accounting can be done to determine as a residual the quantity evaporation minus precipitation. In practice, however, the distribution of water in the atmosphere is not continuously known; rather only water in the vapour state is sampled and only at time intervals of 12 hours and over distances of hundreds of kilometers. The problem is to approximate the water balance from this imperfect sampling procedure, realizing that the computation is only meaningful over sufficiently large areas and for sufficiently large weather systems."

In a coordinate system with pressure as the vertical coordinate, the time rate of change of water and water vapour can be written as:

$$\frac{d}{dt} (q+r) = \frac{\partial}{\partial t} (q+r) + \nabla_2 \nabla_2 (q+r) + \omega \frac{\partial}{\partial p} (q+r)$$
 (54)

where q is the specific humidity, r is the ratio of the mass of water (either in liquid or ice form) to the mass of air,  $\nabla_2$  is the velocity vector on a pressure surface and  $\omega = dp/dt$  is the vertical velocity in the isobaric coordinate system. Equation (54) can be integrated from the earth's surface to some pressure level at which the amount of water in any form is negligible. Then, making use of the continuity equation and the divergence theorem due to Gauss one obtains the so-called atmospheric water balance equation:

$$E - P = \frac{1}{g} \int_{dp} \int_{dA} \frac{\partial}{\partial t} (q+r) dA dp$$

$$+ \frac{1}{g} \int_{dp} \oint_{dt} (c_n q + c_n r) d7 dp$$
(55)

where P and E are the rates of precipitation and evaporation at the earth's surface, g is gravity, dA is an area increment on a pressure surface, dl is a line increment on the vertical boundary and  $c_n$  is the component of the wind vector  $v_2$  normal to the walls of the volume such that it is positive upward. If the integrals in (55) can be properly evaluated through observational data, the exchange of water and water vapour at the earth's surface, given by E - P can be determined as a residual.

One should be able to obtain the same exchange of water at the earth's surface from the hydrological balance provided we

confine ourselves to surface waters. The hydrologic balance equation for a river basin is

$$P - E = R_O + \Delta W + L \tag{56}$$

where  $R_0$  is the runoff from the entire basin,  $\Delta W$  is the total change in the surface and subsurface water storage and L is the depletion from the basin due to use within the basin itself or due to man-made diversion from the basin.

Thus the results from (55) and (56) should agree and this provides a check on the validity of the computation.

Equation (56) is straightforward enough for computation but (55) has to be replaced by the following finite-difference form that is convenient for computation

$$P - E = -\frac{1}{g} \frac{\Delta}{\Delta t} \sum_{j=1}^{N} \sum_{i=1}^{M} q_{ij} \Delta^{A}_{ij} \Delta^{P}_{ij}$$

$$-\frac{1}{g} \sum_{j=1}^{N} \sum_{i=1}^{M} c_{nij} q_{ij} \Delta^{I}_{ij} \Delta^{P}_{y}$$
(57)

Here i denotes the grid point number (with maximum M) and j denotes the pressure level with N being the top most one.

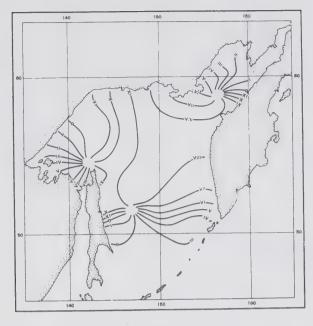
7. Speculation on the possible effects due to man-made changes

It has been shown in Section 4 that the thermohaline processes could give rise to circulation as strong as that due to wind-generation in James Bay and the southern part of Hudson Bay. The discharge from the rivers flowing into James Bay will be

effected when the project proposed by the Quebec government is carried out and it appears that towards the end of the winter season, the thermohaline processes could become stronger and thus further impede the movement of ice. Though this is highly speculative it appears as though the movement of ice in the southern part of Hudson Bay may be even more restricted towards the end of the winter season.

Another speculation could be made on the ice pressure. It appears from Russian work (Kagan, 1967a, 1967b) that because of tidal movement, ice could concentrate in the regions near amphidromic points. To illustrate this we reproduced Figures 2 from Kagan's reports of 1967a and 1967b to produce our Figure 6. The left side shows the cotidal lines for the M2 tide in the Okhotsk Sea while the right side shows the lines of convergence and divergence in the ice pattern. Figure 7 shows the semidiurnal tide in the James Bay-Hudson Bay-Hudson Strait complex (Dohler, 1964). This figure shows that there are no amphidromic points in James Bay while there are two in Hudson Bay. Hence it is possible that while strong and well defined zones of convergence and divergence may not occur in James Bay, these could occur in Hudson Bay and thus create ice pressure situation. However, the proposed hydroelectric power project may not effect the tidal effects on ice.

Although the total water discharged from the rivers into James Bay annually will be essentially unaffected by the hydroelectric power project, it is very likely that the amount of spring discharge into James Bay will be reduced. Because of the



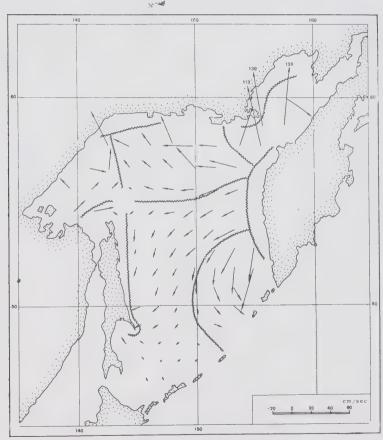
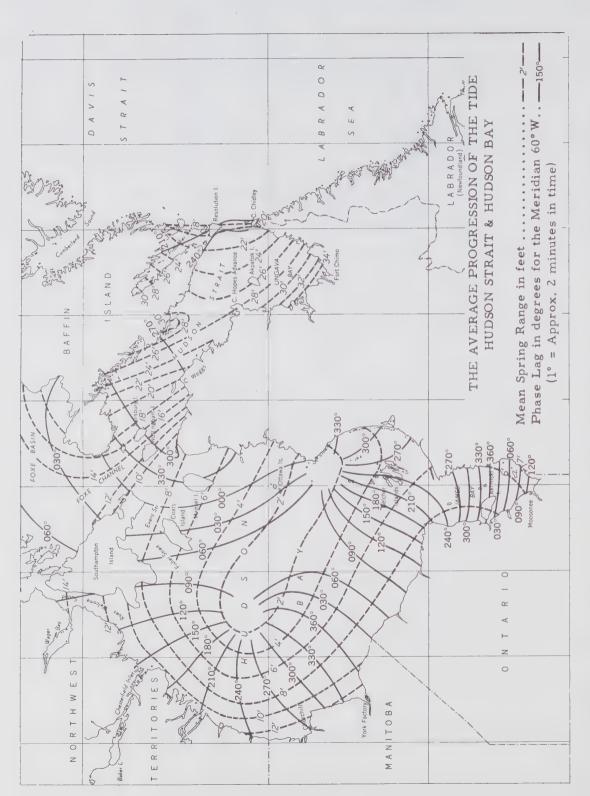


Figure 6 Cotidal lines for the M<sub>2</sub> tide (left side) and convergence-divergence pattern in the ice (right side) in the Okhotsk Sea.



in Hudson Bay-Hudson Cotidal lines for the semidiurnal tide Strait-James Bay complex. Figure

reduction in the freshwater content at the surface, the stability will be reduced and this may give rise to increased upwelling.

The storm surge activity is mainly due to wind stress effects and the effects of baroclinicity usually are negligible. Thus we do not expect any significant changes in the storm surge amplitudes after the construction of the hydroelectric power project.

It appears that the construction of the project may have noticeable effects on the baroclinic coastal jet in James Bay. Because of the reduced freshwater content, the density difference between the bottom and top layers may be reduced. Suppose  $\epsilon$  is of the order of 5 x  $10^{-3}$ , then the coastal jet will be somewhat weakened in the sense that it may have only a width of about 3.5 km.

Next we consider the effect of the project on the estuarine circulation in the central regime of the rivers. As was explained in Section 6 through equation (53), the circulation could be visualized as the combination of a river discharge mode, wind stress mode, and a gravitational convection mode. It is unlikely that the wind stress mode will be affected by the project. However, it is very likely that the dominant term will be the river discharge mode. Naturally the relative influence of this term varies not only with season but also from one river to the other. The gravitational convection mode may play an important role in determining the salinity pattern, noting that the strong tidal action will make some saline water intrude into the rivers.

One of the intractable problems a priori is the effect of the project on the local climate. Our approach to this is to

calculate the atmospheric water balance and compare this with the hydrologic water balance. In the equation for the hydrologic water balance given by (56) allowance is made for the change in water content due to man-made projects (L). However, in the atmospheric water balance equation given by (55) any consequence of human interference has to be felt only indirectly. Thus the only meaningful way to attempt to answer this is to calculate the atmospheric water balance over the river basin system using data before and after the power project construction.

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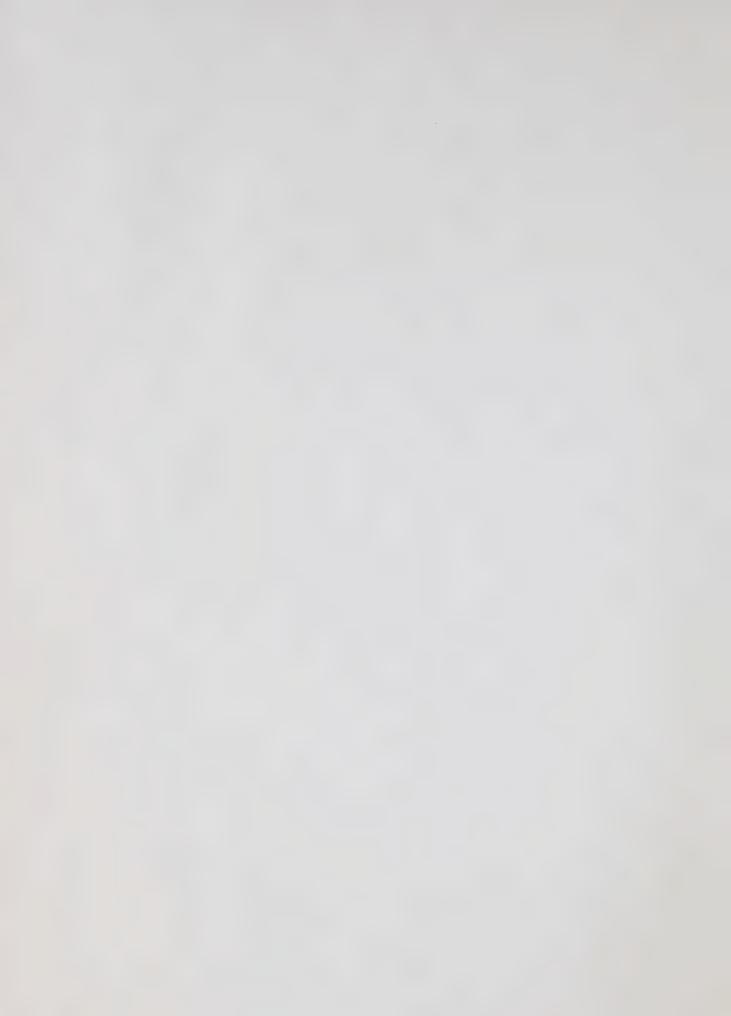
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